

Magnetization Dynamics II: Magnonics: Trends and Challenges

Burkard Hillebrands

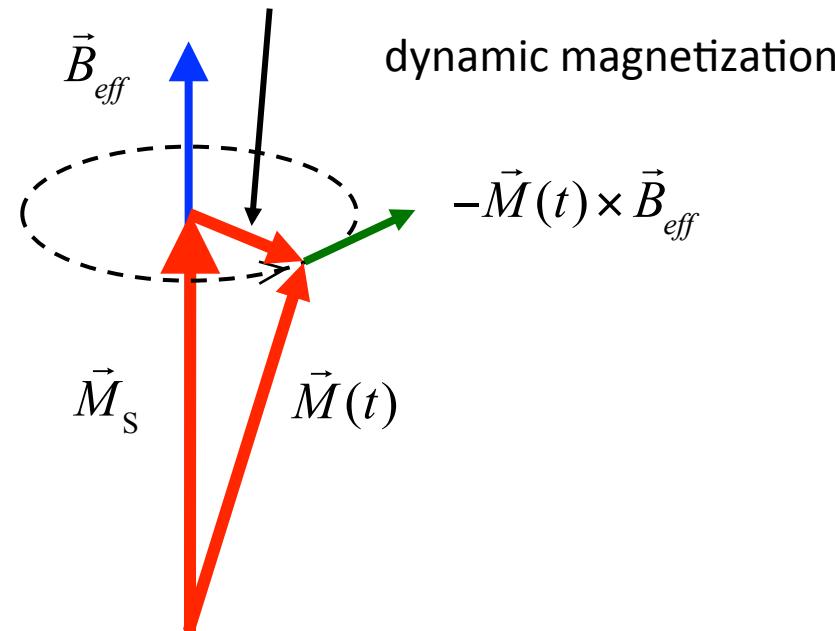
Fachbereich Physik and Landesforschungszentrum OPTIMAS,
Technische Universität Kaiserslautern, Germany



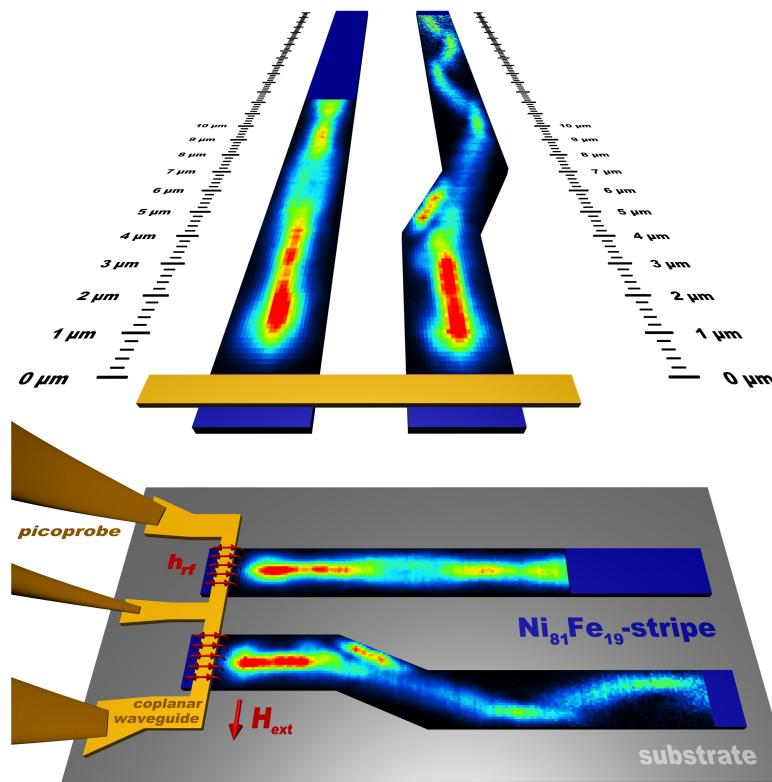
Landau-Lifshitz torque equation

$$\frac{1}{|\gamma|} \frac{d\vec{M}(t)}{dt} = -\vec{M}(t) \times \vec{B}_{eff}(t) + \frac{\alpha}{M_s} \vec{M}(t) \times \frac{d\vec{M}(t)}{dt}$$

$$\vec{m}(\vec{r}, t) = \vec{m}_0(\vec{r}) \times e^{i(\vec{k}\vec{r} - \omega t)}$$



Spin wave bus

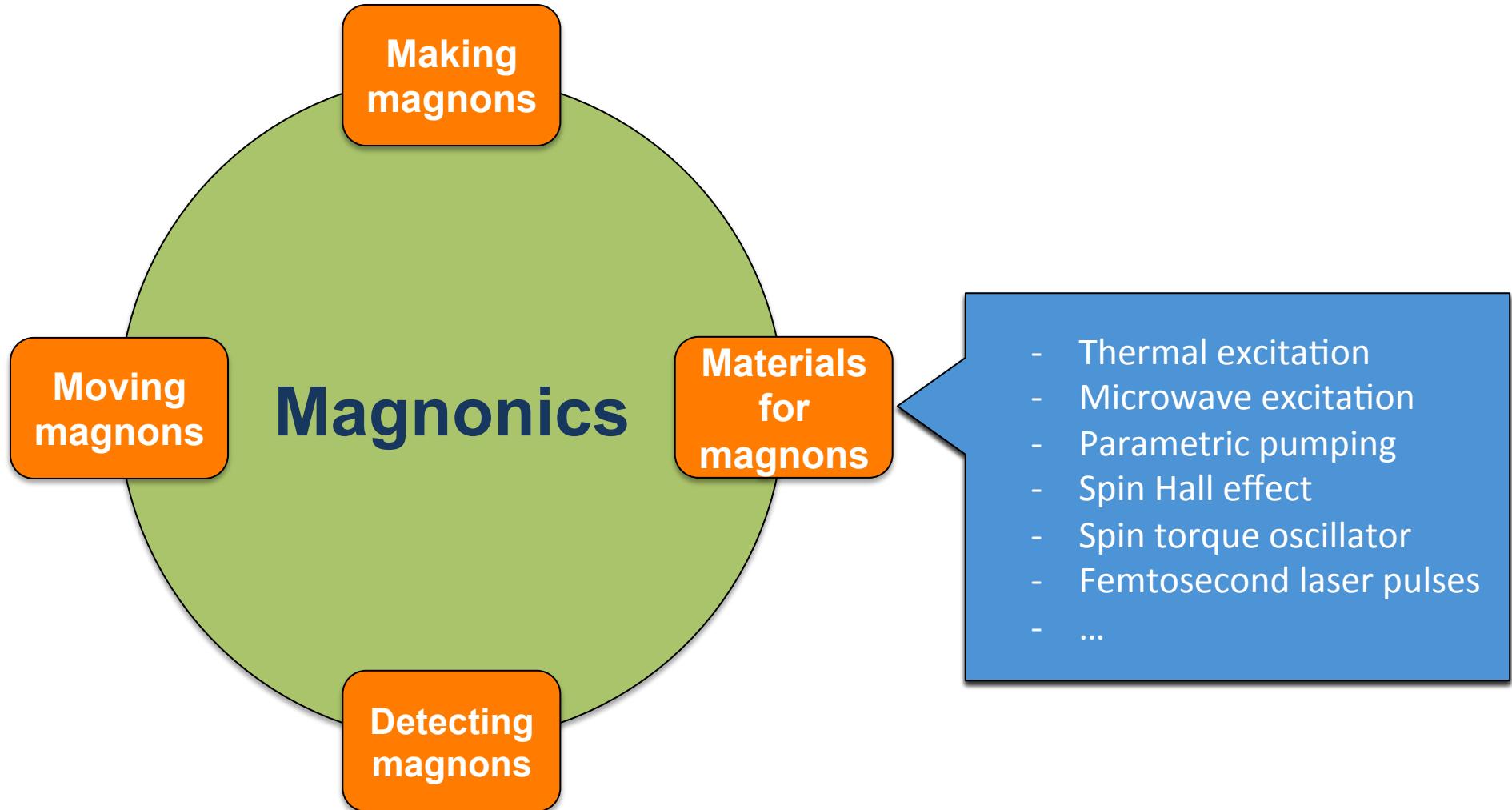


Travelling magnons allow one to:

- **transfer** spin information over centimeter distances
- **process** the information (using wave nature of magnon)
- **operate** in **insulator**-based technology

Fundamental properties:

- Minimal wavelength is down to **several nm**
- Frequency is in GHz and up to the **THz** range
- **Energy:** $E_{\text{magnon}} \ll k_B T$
- **Lifetime:** up to several 100 ns



Kaiserslautern PI Team



A. Chumak



V. Vasyuchka



A. Serga



B. Leven

Main External Collaborators

Y. Ando, E. Saitoh (Tohoku University, Sendai, Japan)

G.A. Melkov (National Taras Shevchenko University of Kiev, Ukraine)

A.N. Slavin (University of Rochester, Michigan, U.S.A.)

A. Karenowska (Oxford University, U.K.)

M. Kostylev (University of Western Australia)



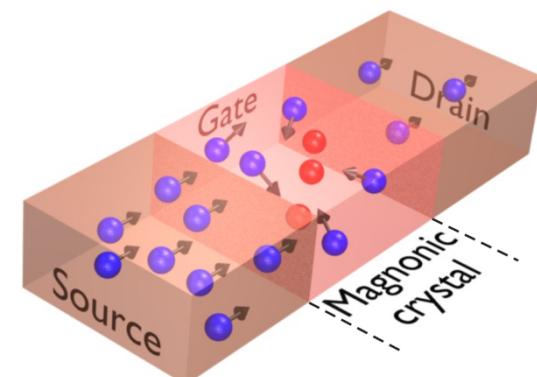
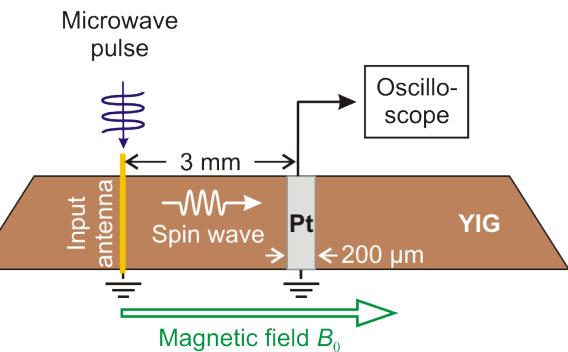
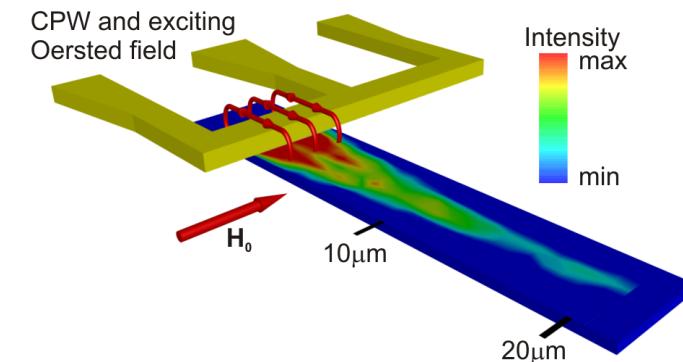
D. A. Bozhko, Dr. M. Agrawal, T. Fischer, S. Klingler, Dr. B. Leven, S. Keller, P. Clausen, L. Mihalceanu, Dr. A. A. Serga, L. Gareis, T. Langner, J. Greser, Dr. P. Pirro, Dr. A. Conca Parra, Jun.-Prof. Dr. E. Th. Papaioannou, Dr. A. Ruiz Calaforra, Dr. T. Brächer, F. Heussner, T. Meyer, Dr. V. I. Vasyuchka, V. Lauer, Prof. Dr. B. Hillebrands + (Dr. A. V. Chumak).

I. New materials for magnonics

II. Novel means for magnon detection

III. Data processing using magnons

IV. Magnonic supercurrents



I. New materials for magnonics

Main requirement: small damping parameter

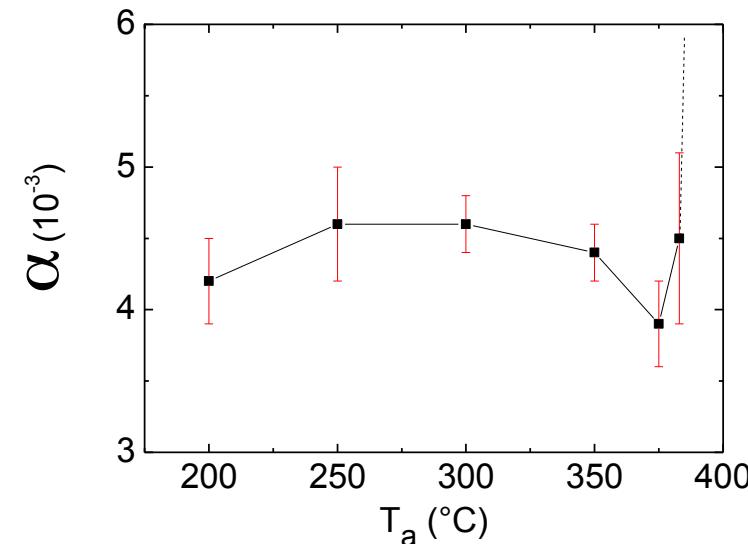


Commonly used material: Permalloy (Py, NiFe): $\alpha = 8 \times 10^{-3}$

I. New materials for magnonics

Main requirement: small damping parameter

- Commonly used material: Permalloy (Py, NiFe): $\alpha = 8 \times 10^{-3}$
- CoFeB: low damping (both am. and crystalline phase): $\alpha = 4 \times 10^{-3}$



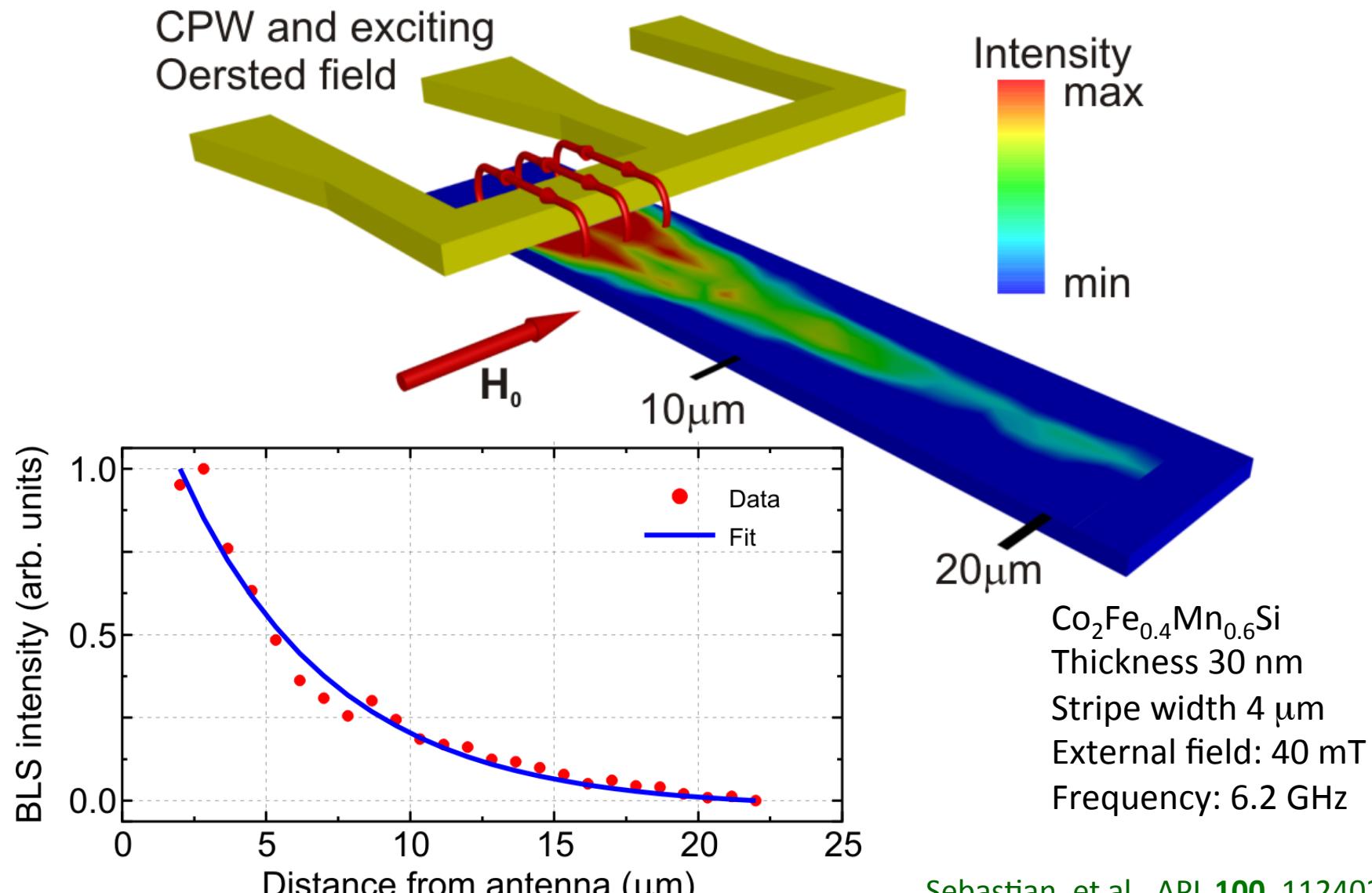
Liu, et al. JAP **110**, 033910 (2011)
Conca, et al. APL **104**, 182407 (2014)

I. New materials for magnonics

Main requirement: small damping parameter

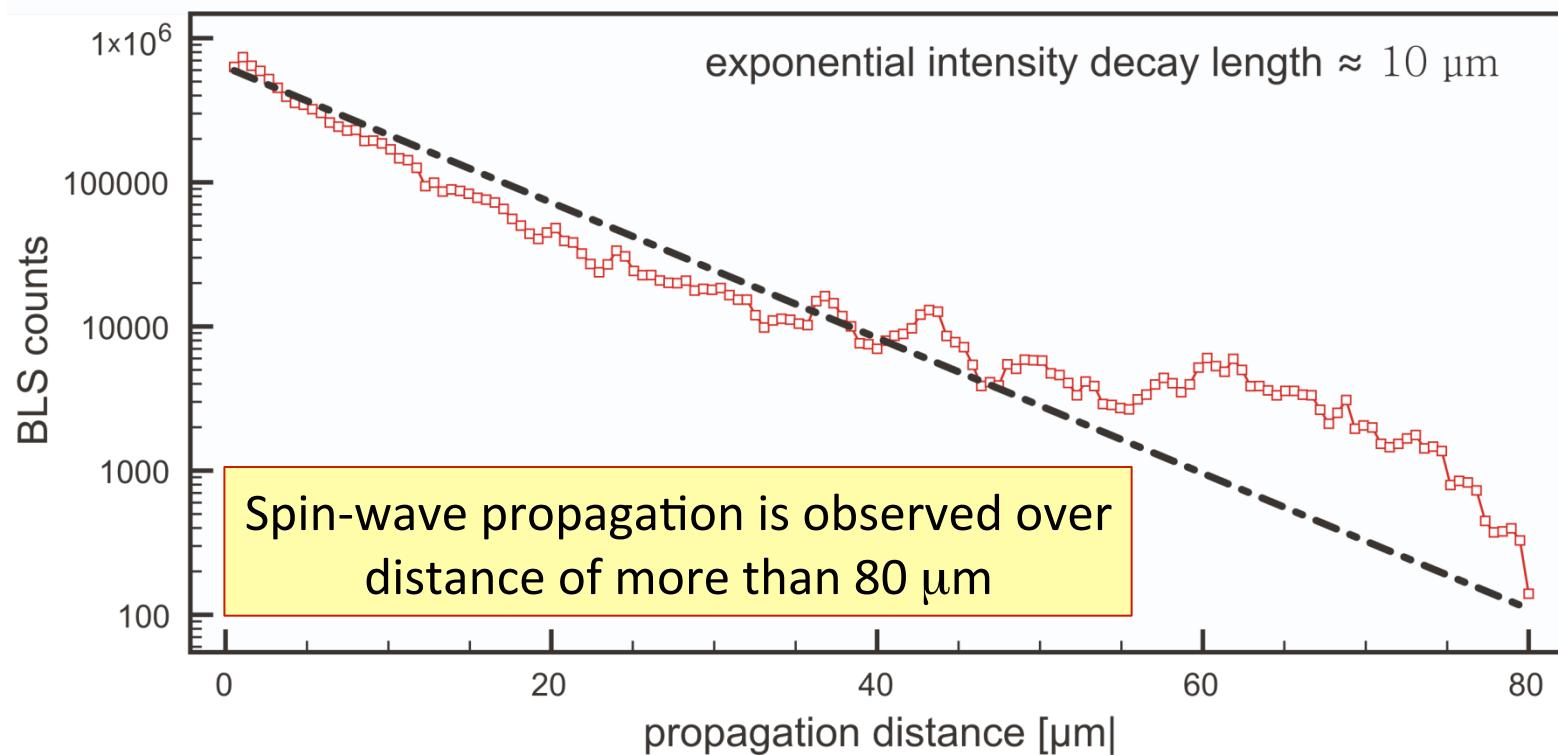
- Commonly used material: Permalloy (Py, NiFe): $\alpha = 8 \times 10^{-3}$
 - CoFeB: low damping (both am. and crystalline phase): $\alpha = 4 \times 10^{-3}$
- Novel Heusler compounds: $\alpha = 3 \times 10^{-3}$

Spin-wave propagation in $\text{Co}_2\text{Fe}_{0.4}\text{Mn}_{0.6}\text{Si}$ Heusler waveguides



Sebastian, et al., APL **100**, 112402 (2012)

Spin-wave propagation in $\text{Co}_2\text{Fe}_{0.4}\text{Mn}_{0.6}\text{Si}$ Heusler waveguides



$$I(x) = I_0 \exp\left(-\frac{2x}{\delta}\right) + b$$

$$\alpha = \frac{1}{\tau \gamma \mu_0 (H_{eff} + M_s)} \quad \tau = \frac{\delta}{v_G}$$

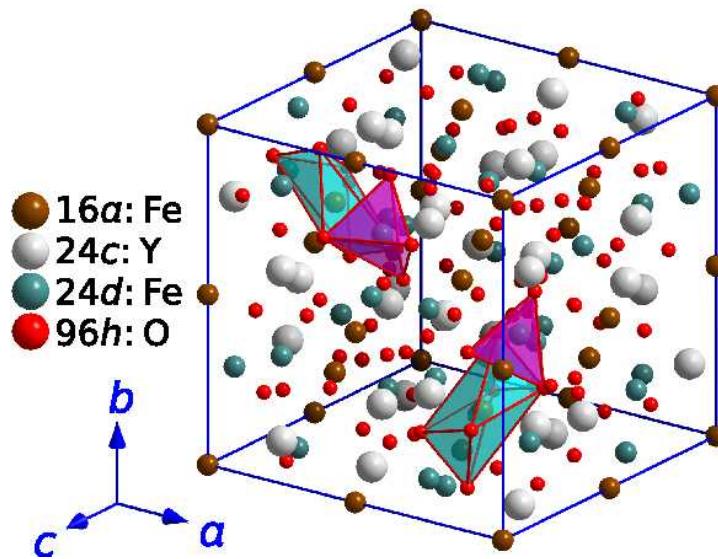
Decay length: $10.6 \mu\text{m}$
 Damping α : 4.7×10^{-3}

Sebastian, et al., APL **100**, 112402 (2012)
 Sebastian, et al., PRL **110**, 067201 (2013)

I. New materials for magnonics

- Commonly used material: Permalloy (Py, NiFe): $\alpha = 8 \times 10^{-3}$
 - CoFeB: low damping (both am. and crystalline phase): $\alpha = 4 \times 10^{-3}$
 - Novel Heusler compounds: $\alpha = 3 \times 10^{-3}$
-  Micro-structured Yttrium Iron Garnet $\text{Y}_3\text{Fe}_5\text{O}_{12}$ (YIG): $\alpha = 4 \times 10^{-5}$

Yttrium Iron Garnet $\text{Y}_3\text{Fe}_5\text{O}_{12}$ (YIG)



YIG:

- magnetic insulator
- smallest spin-wave damping

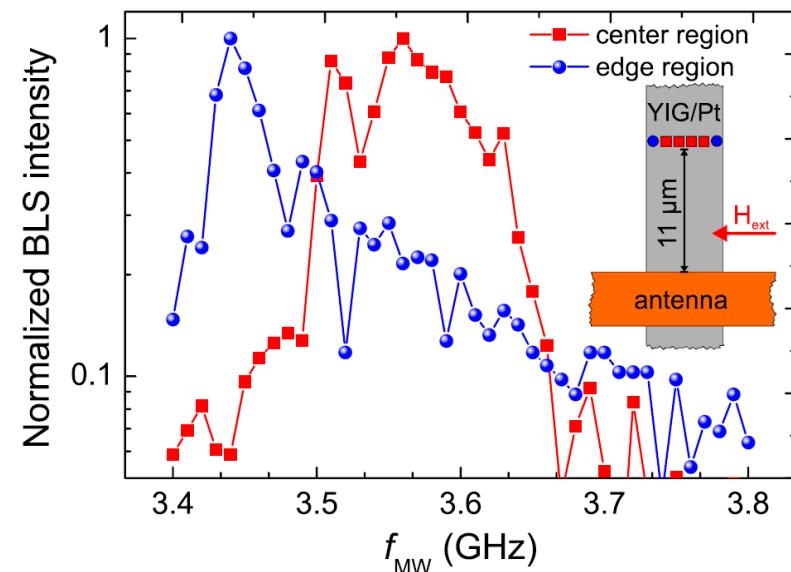
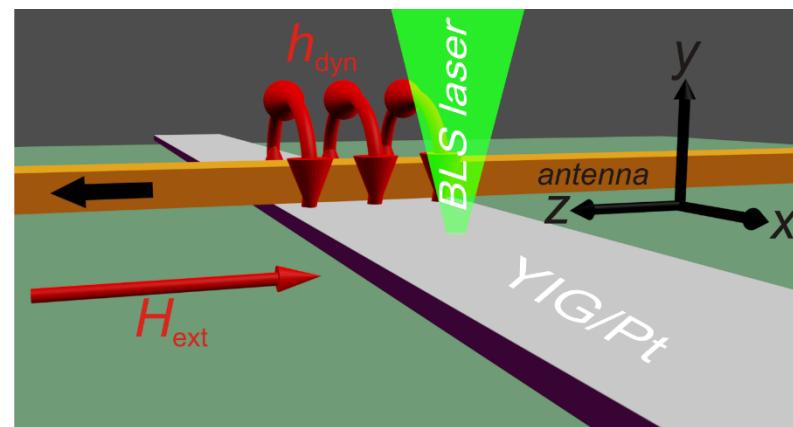
Preparation via:

- liquid phase epitaxy
- sputtering
- pulsed laser deposition

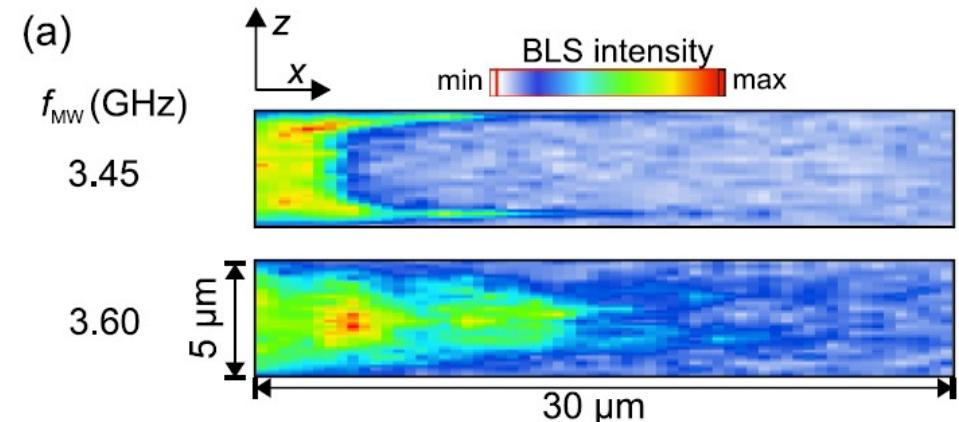
A. Kreisel, Europhysics News (2006)

Sample fabrication

Micro-focused Brillouin Light Scattering setup was used for magnon detection



BLS intensity map: YIG thickness: 100 nm
made by liquid phase epitaxy



For standard YIG quality:
free path will be up to **1 mm**

P. Pirro, et al., APL **104**, 012402 (2014)

I. New materials for magnonics

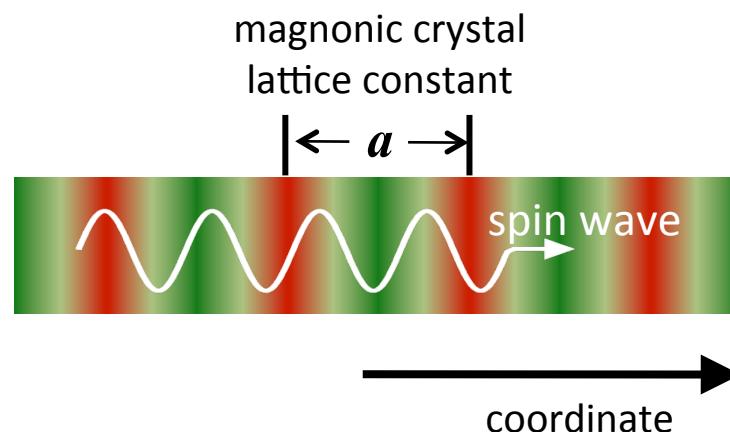
- Commonly used material: Permalloy (Py, NiFe): $\alpha = 8 \times 10^{-3}$
 - CoFeB: low damping (both am. and crystalline phase): $\alpha = 4 \times 10^{-3}$
 - Novel Heusler compounds: $\alpha = 3 \times 10^{-3}$
 - Micro-structured Yttrium Iron Garnet $\text{Y}_3\text{Fe}_5\text{O}_{12}$ (YIG): $\alpha = 4 \times 10^{-5}$
- Magnonic crystals: artificial magnetic materials (static and dynamic)

What is a “magnonic crystal”?

Magnonic crystal – magnetic meta-material:

- ❖ artificial medium with periodic lateral variation in magnetic properties

One-dimensional magnonic crystal:

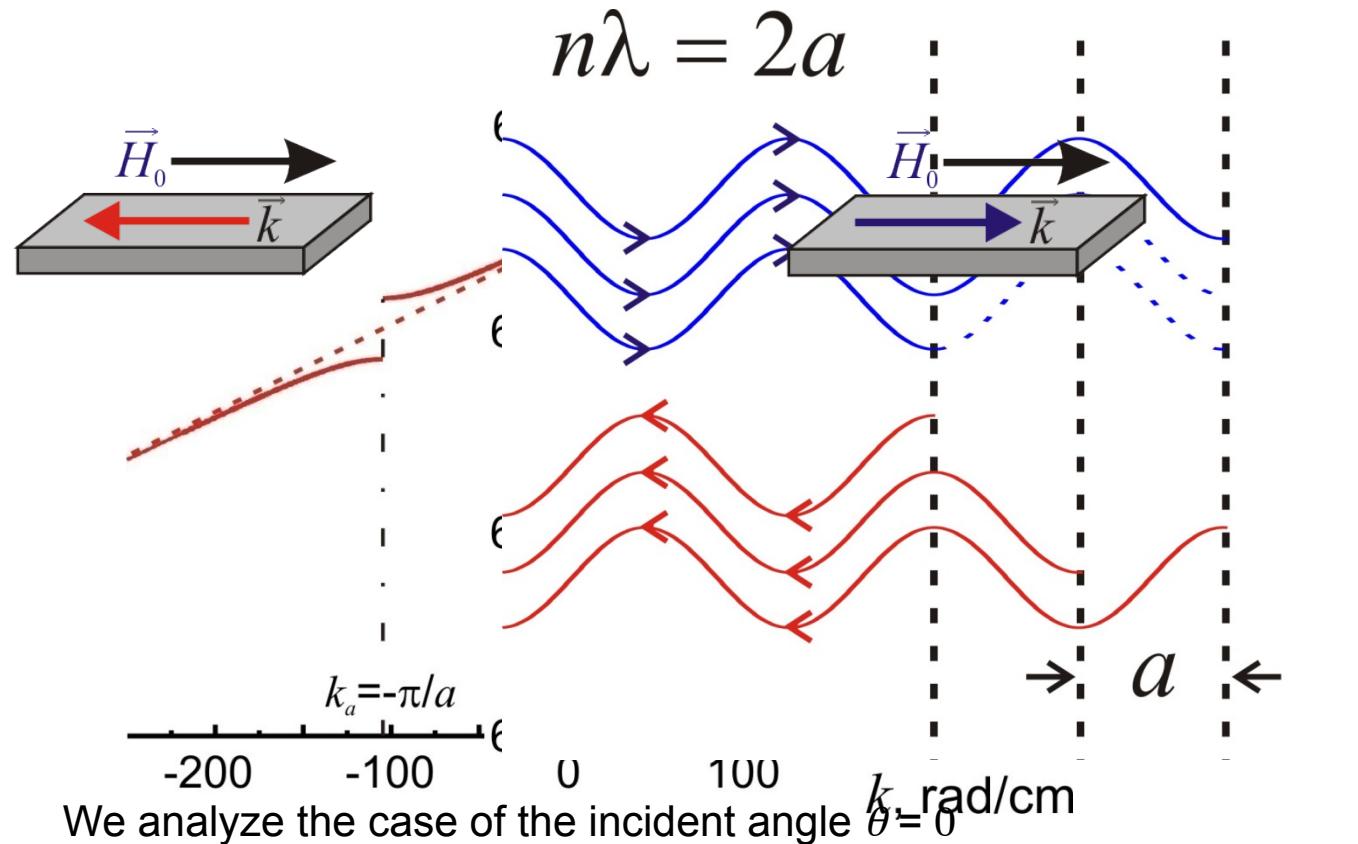


Magnonic-crystal are
engineered to have properties
that may not be found in nature

- ❖ analogous to **photonic and sonic** crystals but operates with spin waves in the GHz frequency range

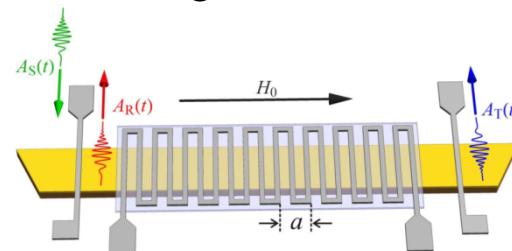
Band gap

Band gaps – regions of the spectrum over which waves are **not allowed** to propagate



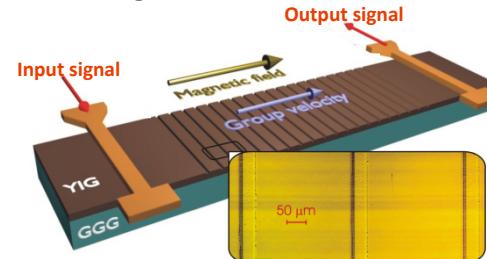
Which magnetic property do we modulate?

Bias magnetic field



Chumak et al., J. Phys. D **42**, 205005 (2009)

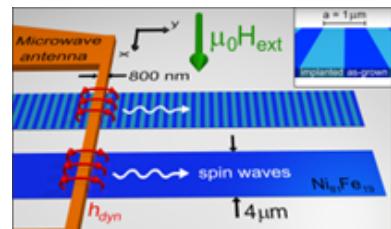
Waveguide thickness



Sykes et al., APL **29**, 388 (1976)
Chumak et al., APL **93**, (2008)

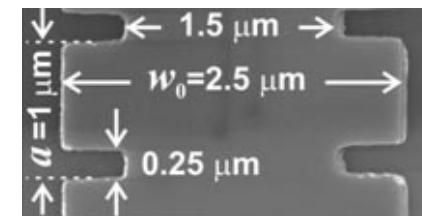
$$f(k) = \gamma \left(H_0 + 4\pi M_0 \frac{1 - \exp\{-\sqrt{(\pi/w)^2 + k^2}d\}}{\sqrt{(\pi/w)^2 + k^2}d} \right)$$

Effective saturation magnetization



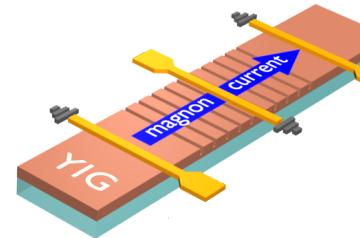
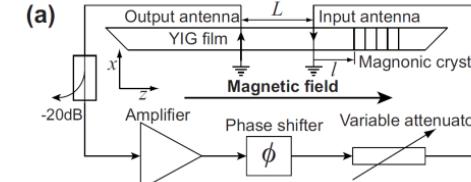
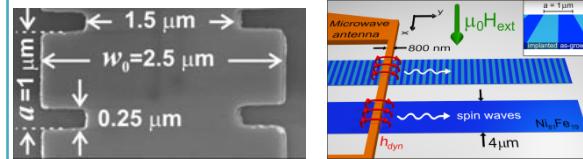
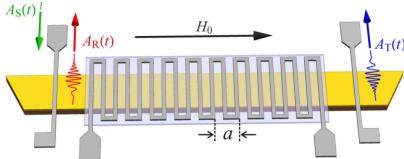
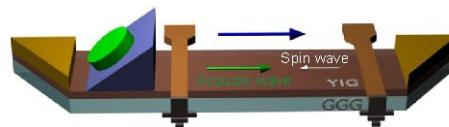
Chumak et al., PRB, **81**, 140404 (2010)
Obry et al., APL, **102**, 202403 (2013)

Waveguide width



Chumak et al., APL **95**, (2009)
Lee et al., PRL **102**, 127202 (2009)

Magnonic crystals - overview

microwave filter		Chumak et al., APL 93 , 022508 (2008) Chumak et al., APL 94 , 172511 (2009) Chumak et al., JAP 105 , 083906 (2009)
data storage element		Chumak et al., PRL, 108 , 257207 (2012)
stable microwave generator		Karenowska et al., APL 96 , 082505 (2010)
micro-sized crystal		Chumak et al., APL 95 , 262508 (2009) Ciubotaru et al., J.Phys.D 45 , 255002 (2012) Obry et al., APL, 102 , 202403 (2013) Ciubotaru, et al. PRB 88 , 134406 (2013)
switchable device		Chumak, et al., J.Phys.D 42 , 205005 (2009) Chumak, et al., Nat. Commun. 1 :141 (2010) Karenowska, et al., PRL 108 , 015505 (2012)
travelling crystal		Chumak et al., PRB, 81 , 140404 (2010)

Localized ion implantation

Localized ion implantation:

- Purely magnetic patterning
- No change in sample topography

State of the art studies (YIG):

G. Volluet, P. Hartemann,

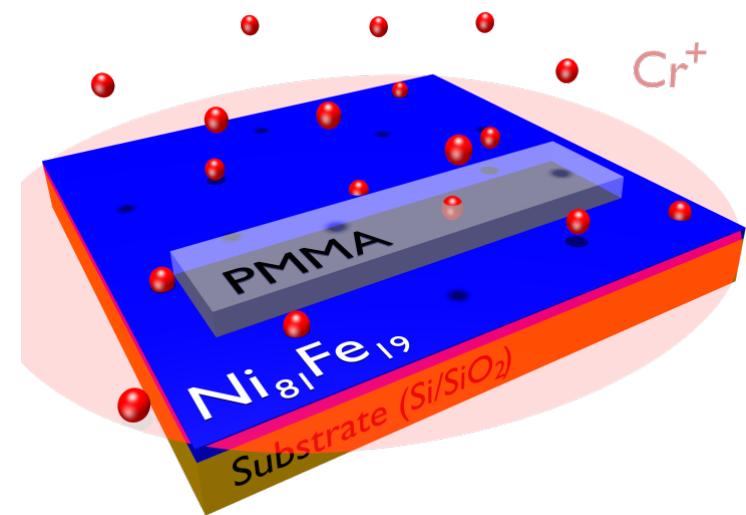
Reflection of magnetostatic forward volume waves by ion implanted gratings,
Proc. IEEE Ultrasonics Symp., 394 (1981).

R. L. Carter, J. M. Owens, C. V. Smith, Jr., K. W. Reed,
Ion-implanted magnetostatic wave reflective array filters,
J. Appl. Phys. 53 (1982), 2655.

Irradiation of $\text{Ni}_{81}\text{Fe}_{19}$ films with 30 keV
 Cr^+ ions:

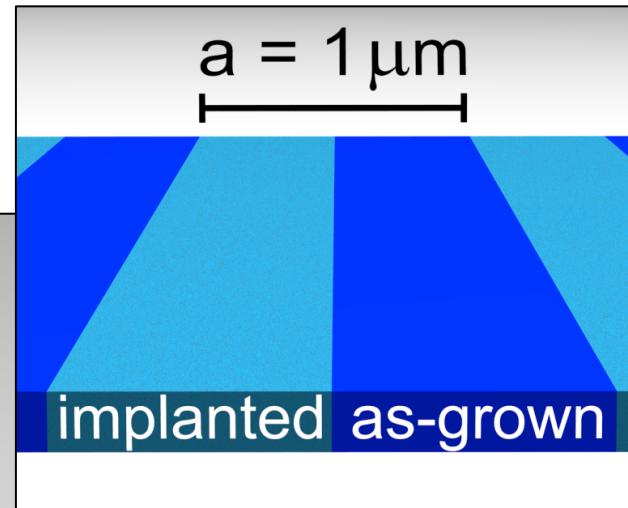
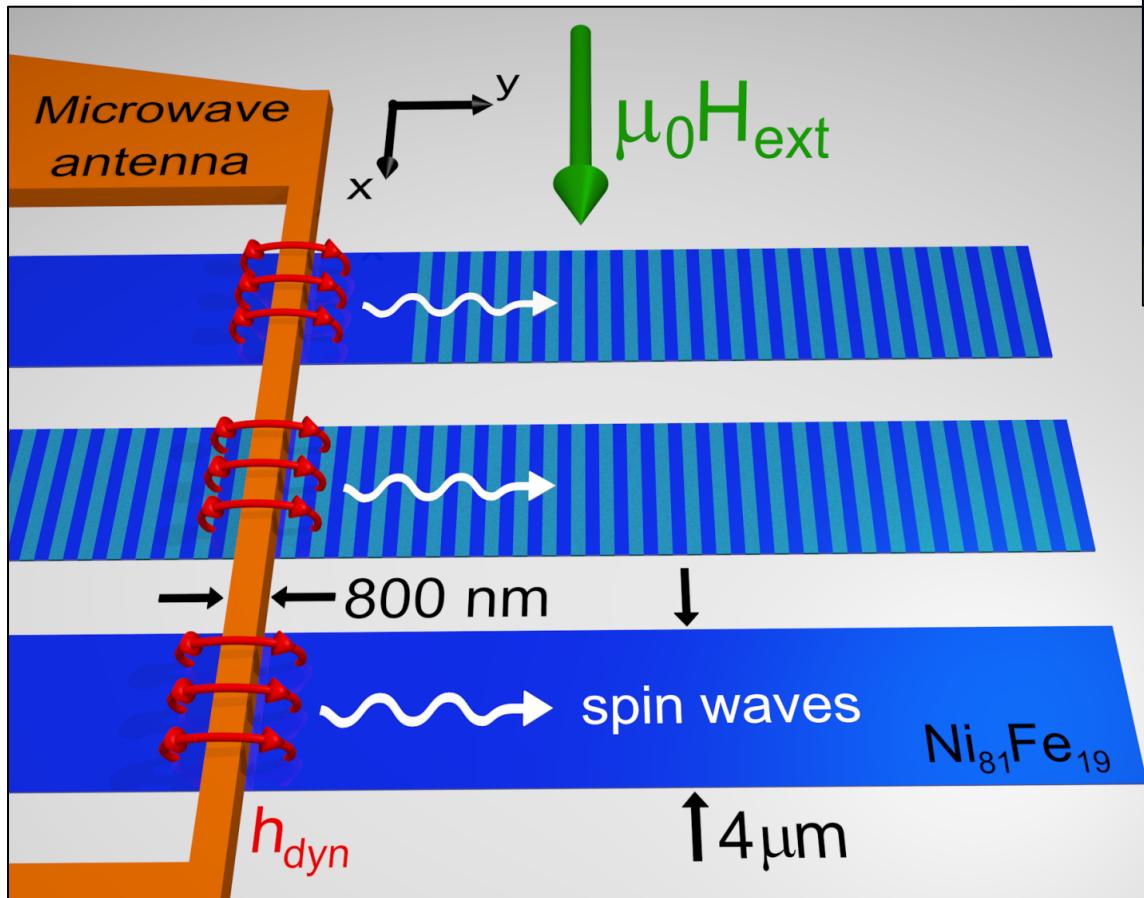
Control saturation magnetization M_s
and Gilbert damping α

Fassbender *et al.*, PRB 73, 184410 (2006)



Sample setup

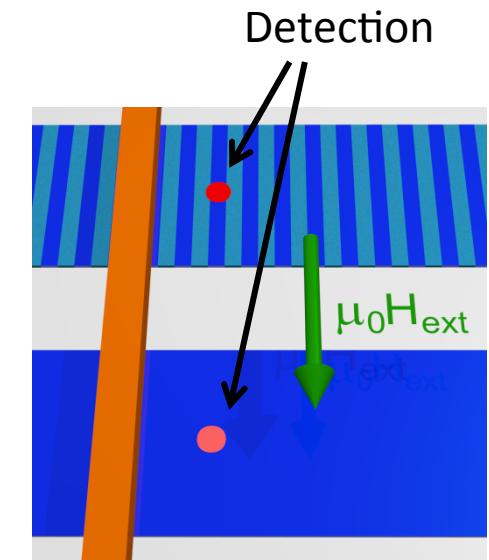
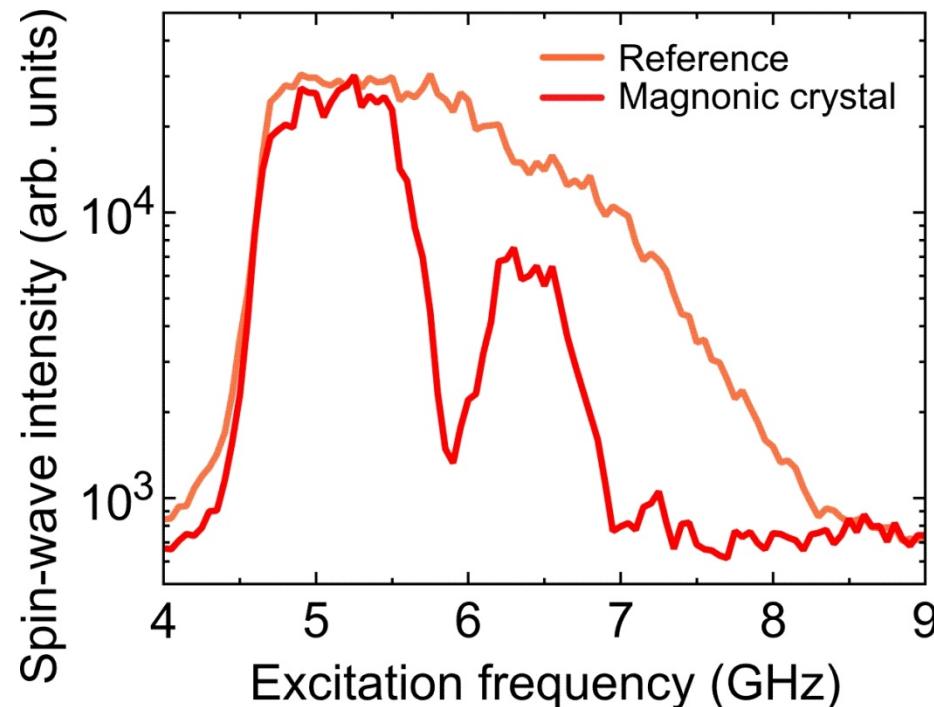
Fabrication of microscopic metallic magnonic crystal with periodic change in saturation magnetization M_s



- **Waveguides:**
MBE evaporation
Lift-off techniques

Transmission spectrum

Spin-wave excitation spectra:



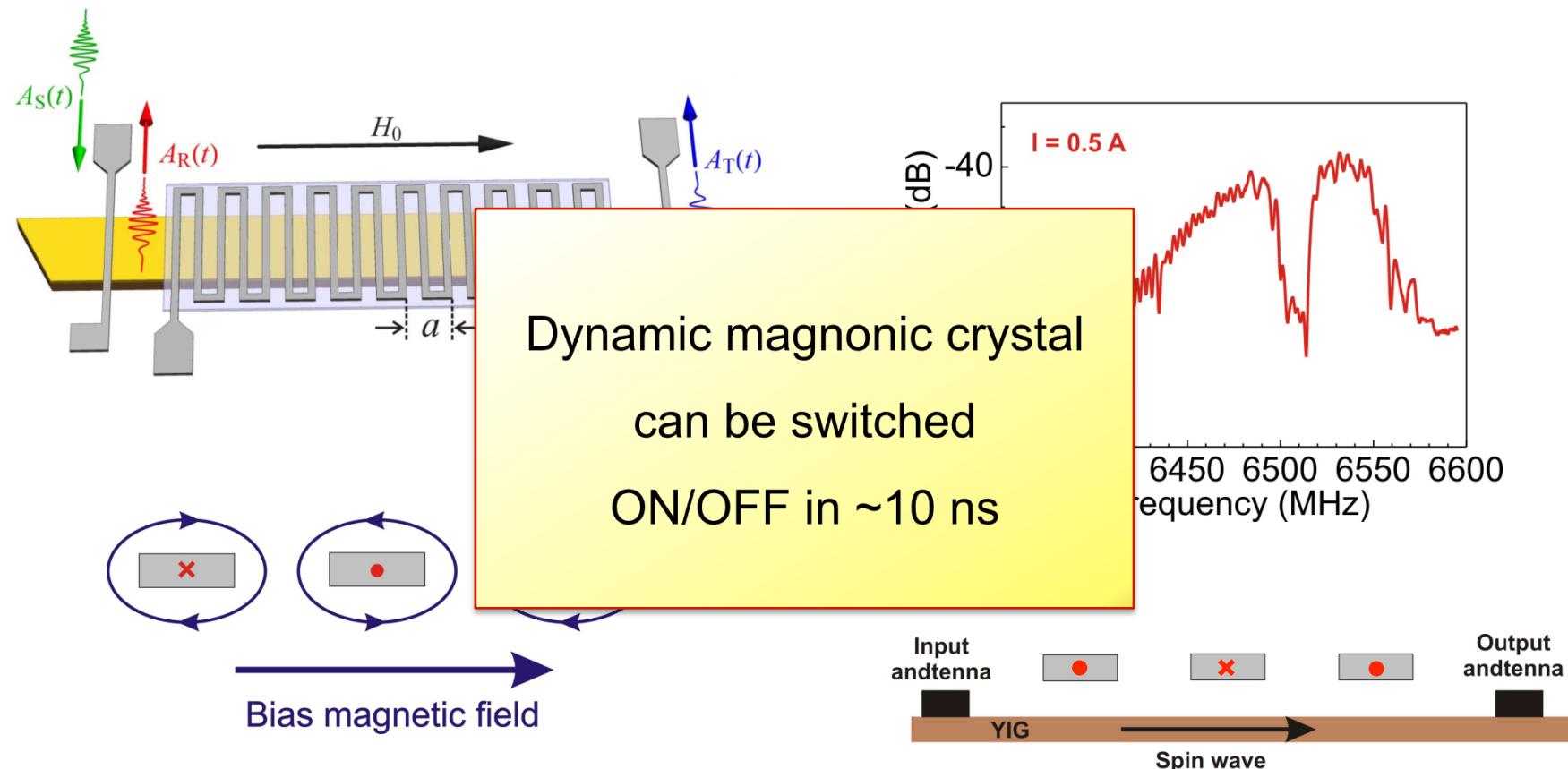
Two pronounced band gaps in transmission spectrum

Obry *et al.*, APL **102**, 202403 (2013)

Ciubotaru, et al. PRB **88**, 134406 (2013)

Dynamic magnonic crystal

Periodic modulation of the bias magnetic field by **current-carrying wires**

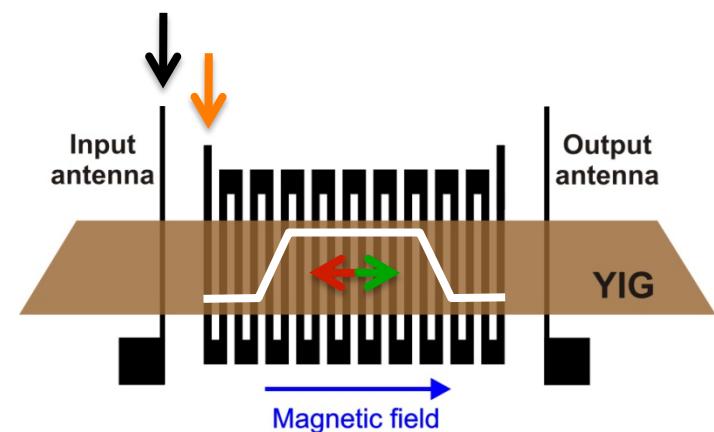
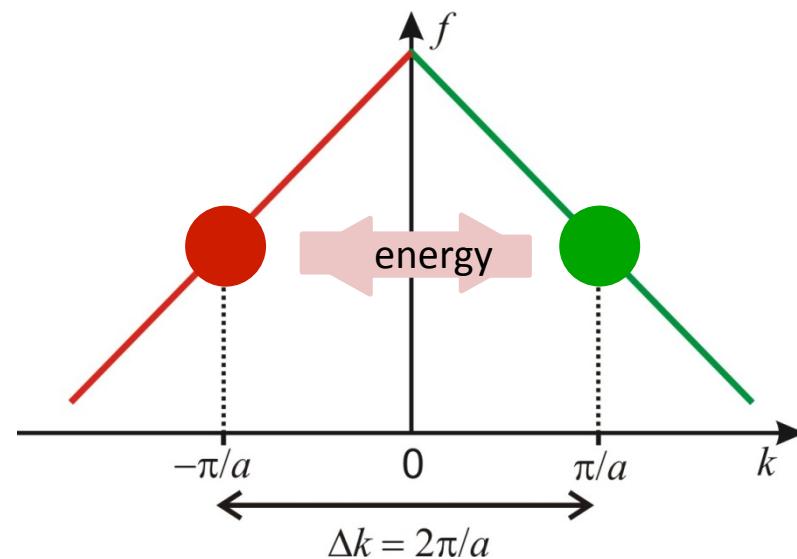


Lattice constant: $a = 300 \mu\text{m}$

Number of periods: $N_g = 20$

Chumak, et al., J. Phys. D **42**, 205005 (2009)

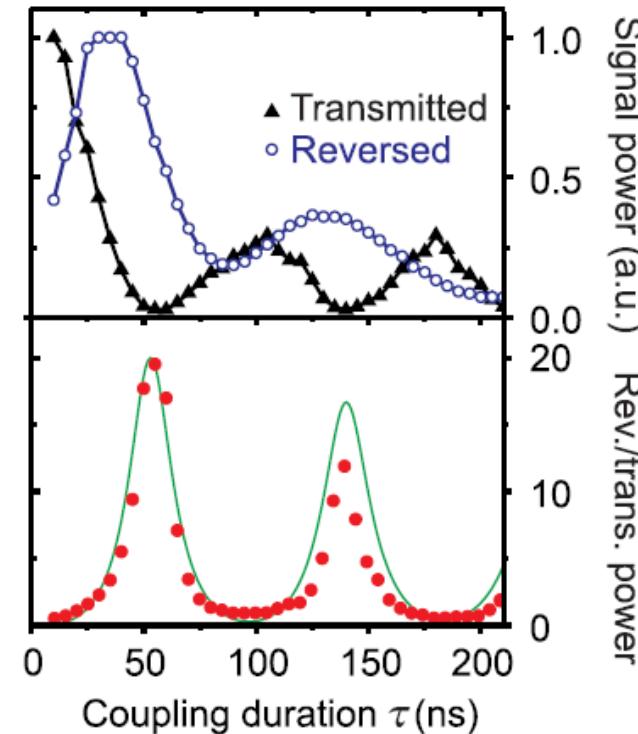
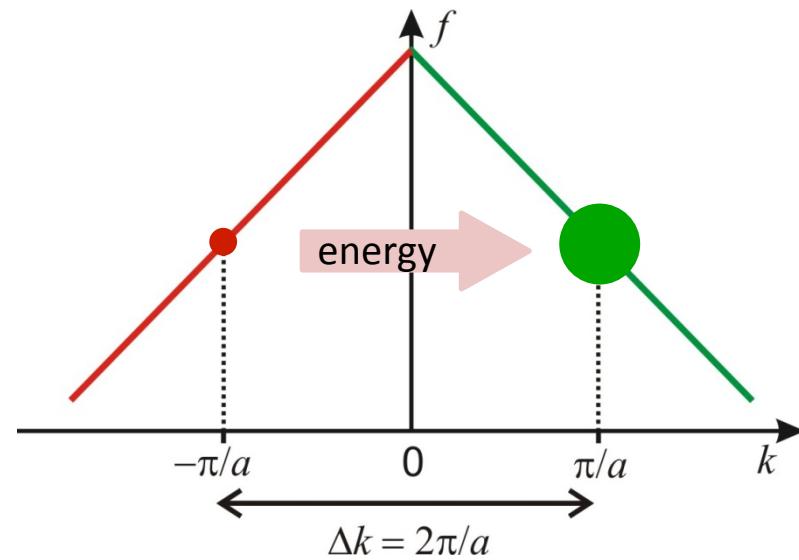
Spin-wave mode coupling by dynamic magnonic crystal



Two modes with $k = \pi/a$ and $k = -\pi/a$ are coupled by periodic variation of field

Coupling provides a mechanism for energy transfer

Spin-wave mode coupling by dynamic magnonic crystal



Chumak, et al., Nat. Commun. **1**:141 (2010)
 Karenowska, et al., PRL **108**, 015505 (2012)

Two modes with $k = \pi/a$ and $k = -\pi/a$ are coupled by periodic variation of field

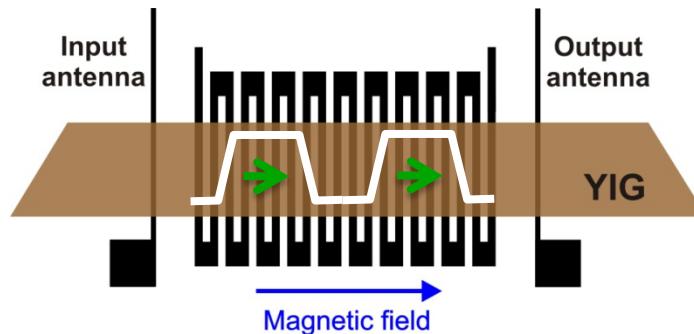
Coupling provides a mechanism for energy transfer

All-linear time reversal by DMC

Input signal:

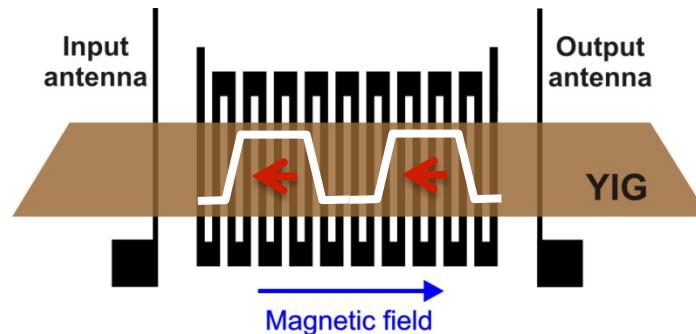
$$A_S(t) \sim \sum_{\Delta f} \exp(-i 2\pi \Delta f t)$$

Δf is a frequency shift from the Bragg frequency

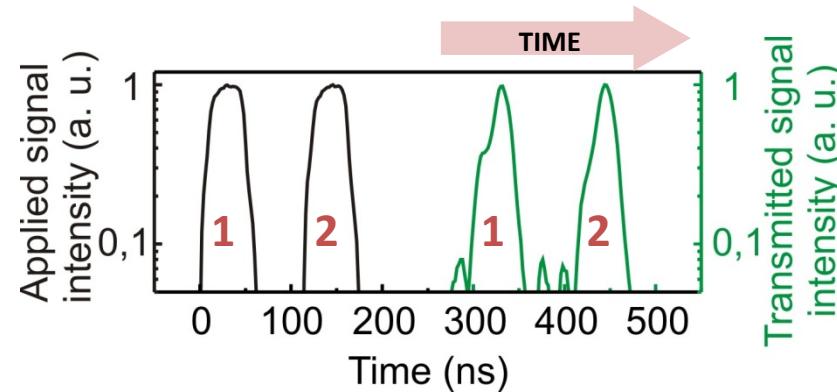


Reflected signal:

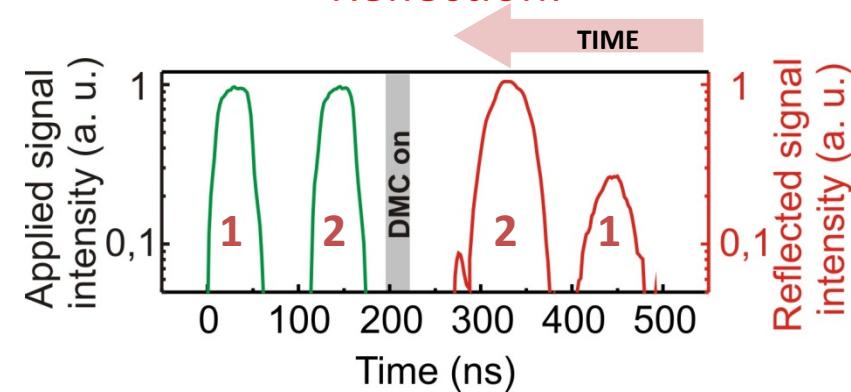
$$A_R(t) \sim \sum_{\Delta f} \exp(i 2\pi \Delta f t) \sim A_S(-t)$$



Transmission:



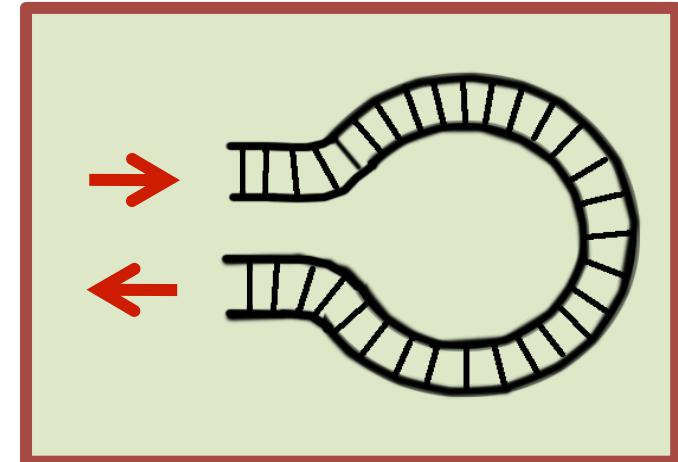
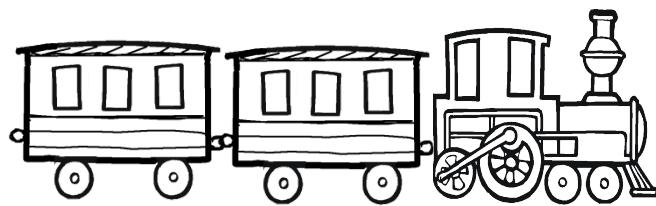
Reflection:



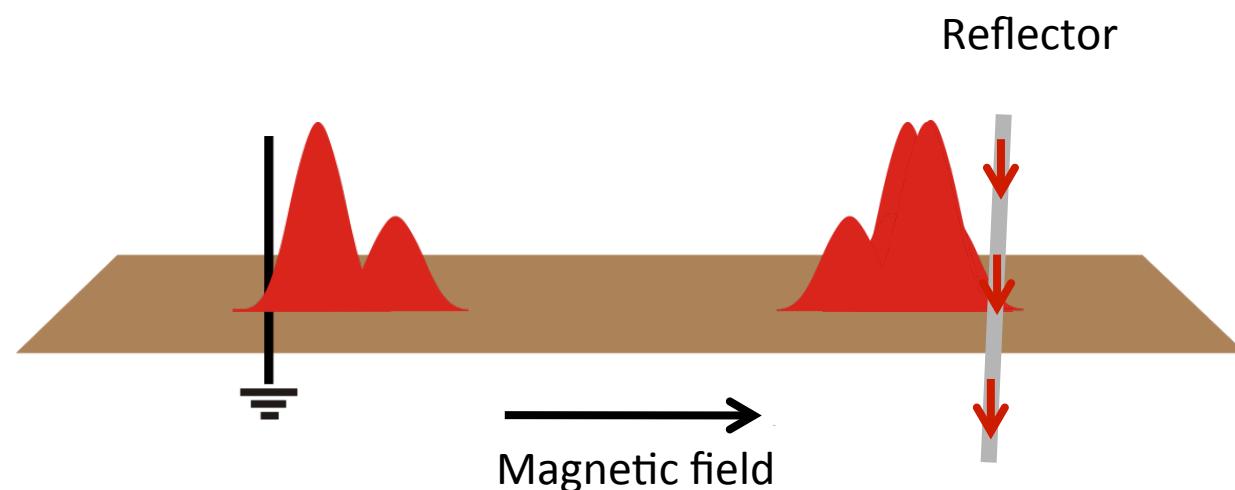
Chumak et al., *Nat. Commun.* 1:141 doi: 10.1038/ncomms1142 (2010)

Classical reflection from mirror

Railway analog:

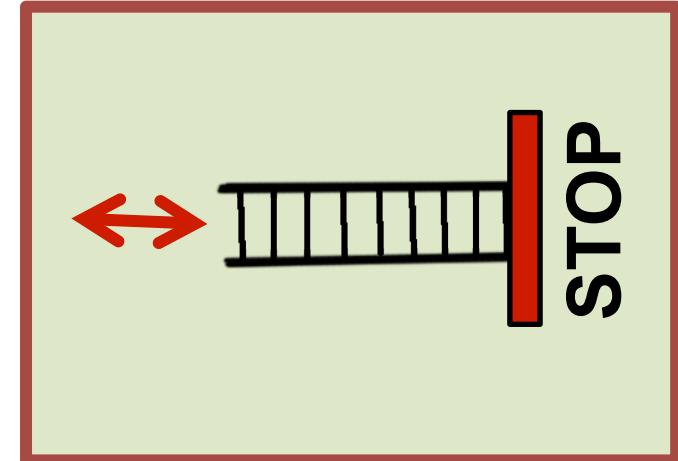
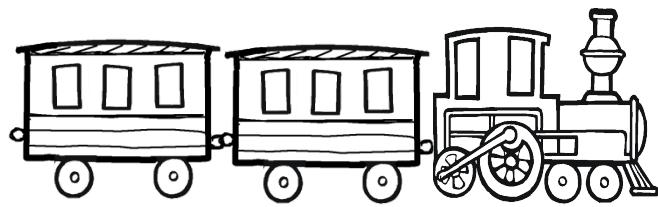


Spin-wave experiment:

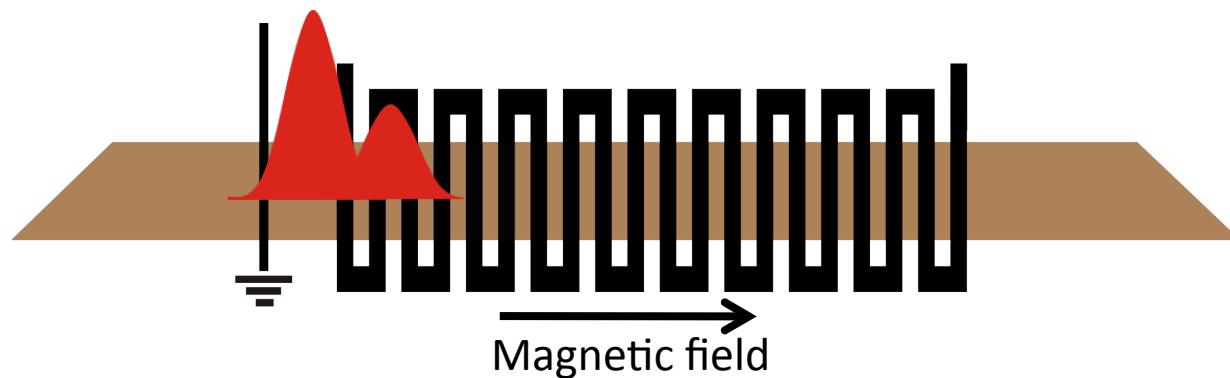


Reflection via time reversal

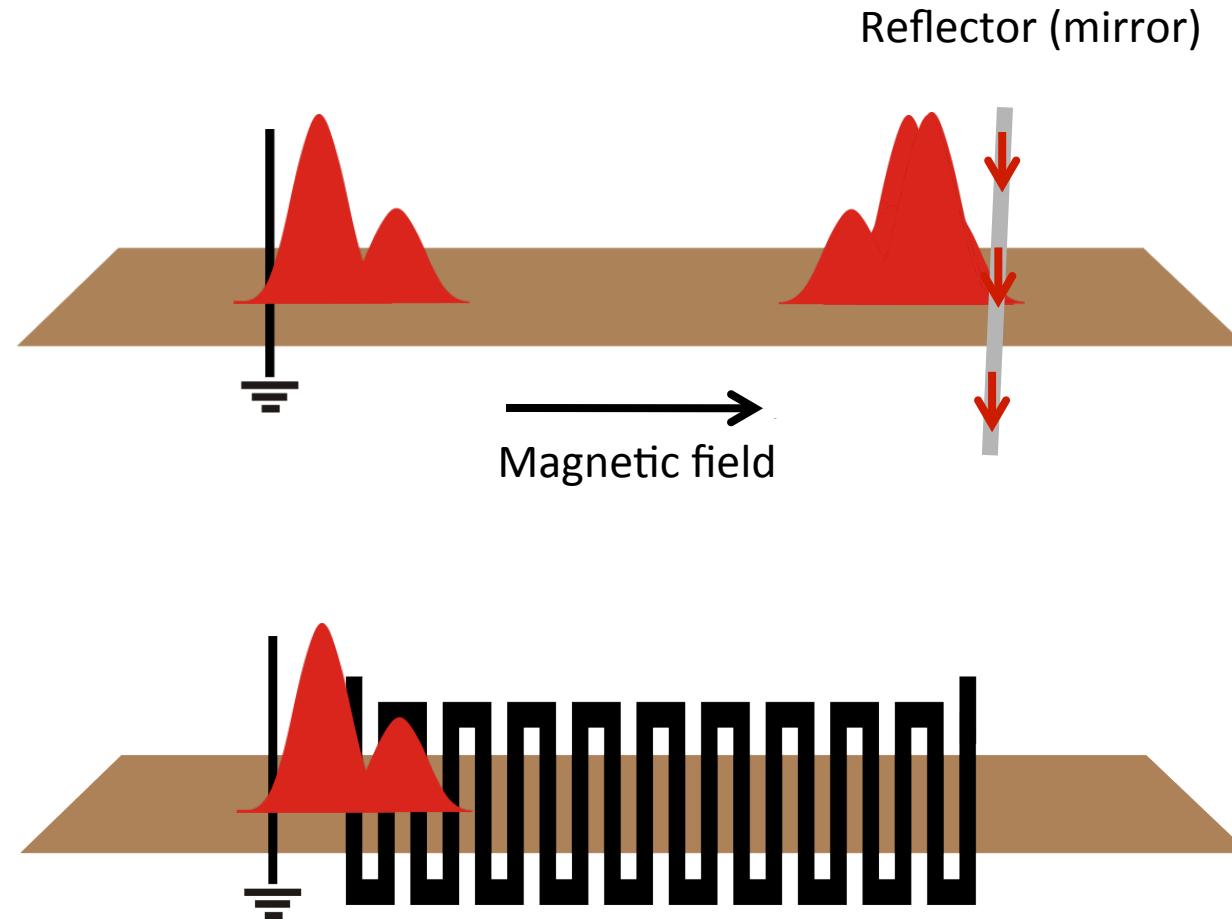
Railway analog:



Spin-wave experiment:



Reflection from a mirror



Chumak, et al., Nat. Commun. **1**:141 (2010)

Karenowska, et al., PRL **108**, 015505 (2012)

I. New materials for magnonics

- Commonly used material: Permalloy (Py, NiFe): $\alpha = 8 \times 10^{-3}$
 - CoFeB: low damping (both am. and crystalline phase): $\alpha = 4 \times 10^{-3}$
 - Novel Heusler compounds: $\alpha = 3 \times 10^{-3}$
 - Micro-structured Yttrium Iron Garnet $\text{Y}_3\text{Fe}_5\text{O}_{12}$ (YIG): $\alpha = 4 \times 10^{-5}$
 - Magnonic crystals: artificial magnetic materials (static and dynamic)
- Normally magnetized non-reciprocal materials

Simulation of a perpendicular magnetized Permalloy discs (OOMMF)

Parameters

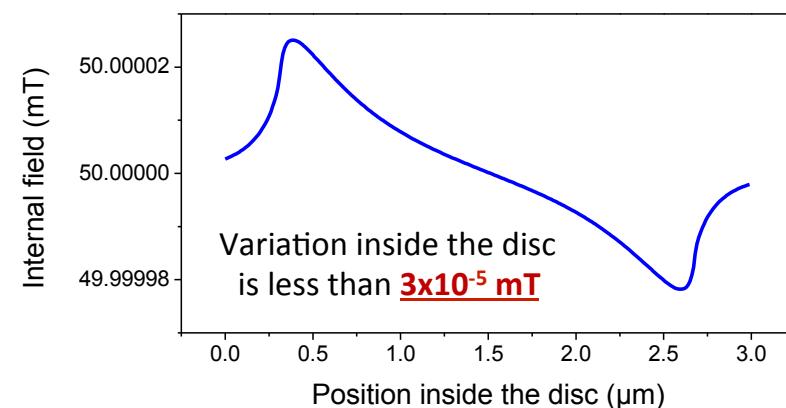
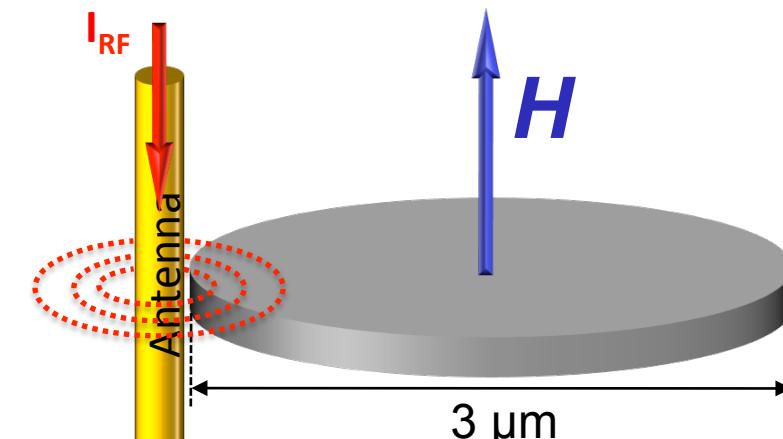
$$\begin{aligned}\mu_0 M_s &= 1 \text{ T} \\ H_k &= 0 \\ \alpha &= 0.007 \\ \mu_0 H &= 1.050 \text{ T}\end{aligned}$$

Spin waves are excited by the dynamic Oersted field created by injecting an RF current through antenna.

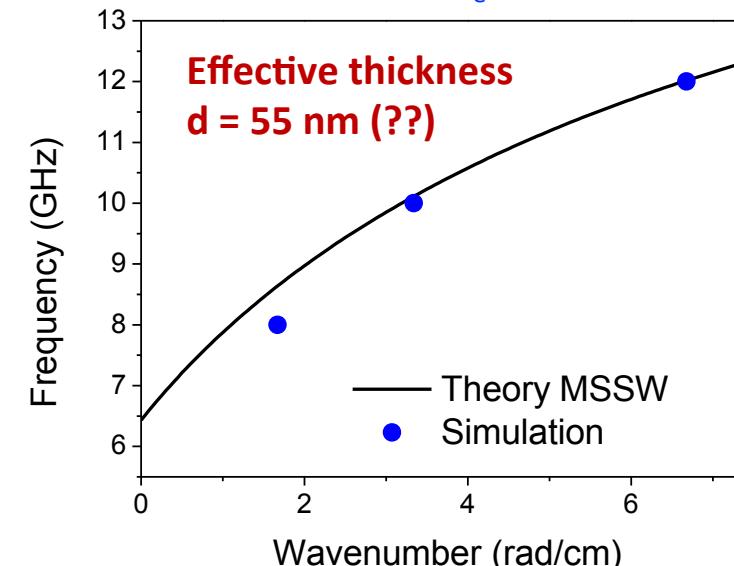
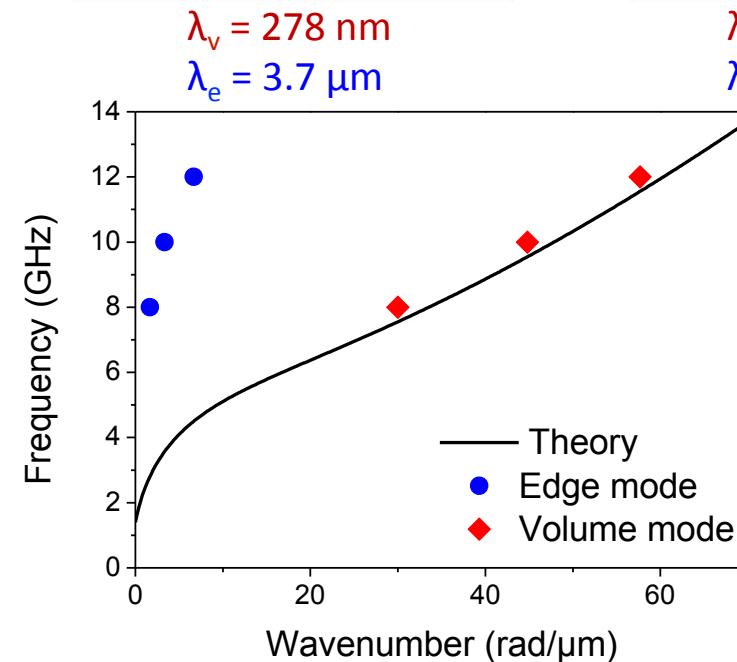
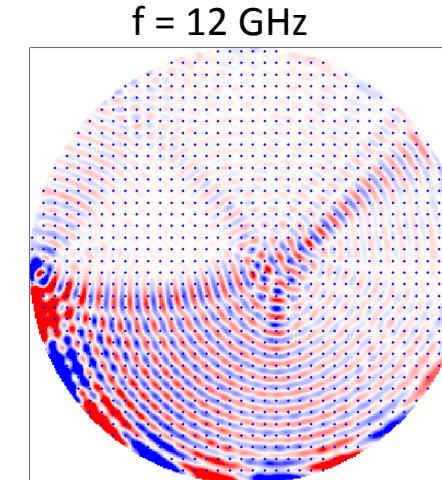
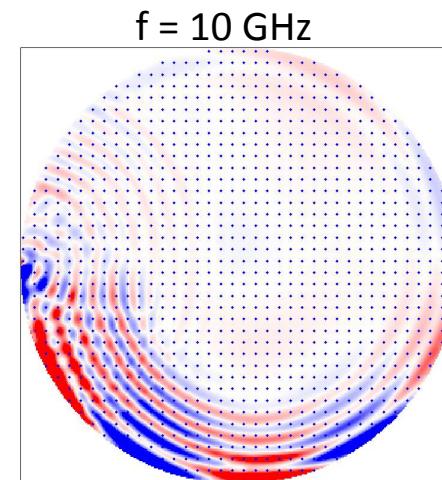
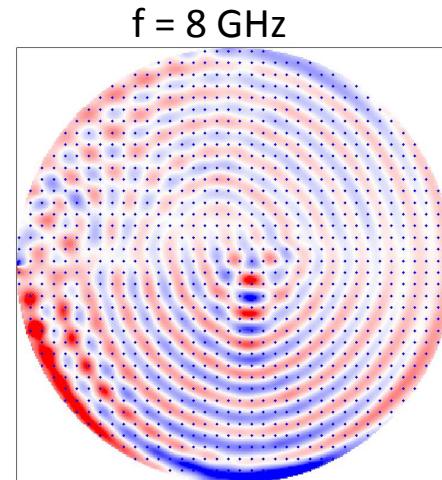
Diameter of antenna: 500 nm

Side charges were added in order to have an uniform distribution of the internal field.

Internal field = 50 mT



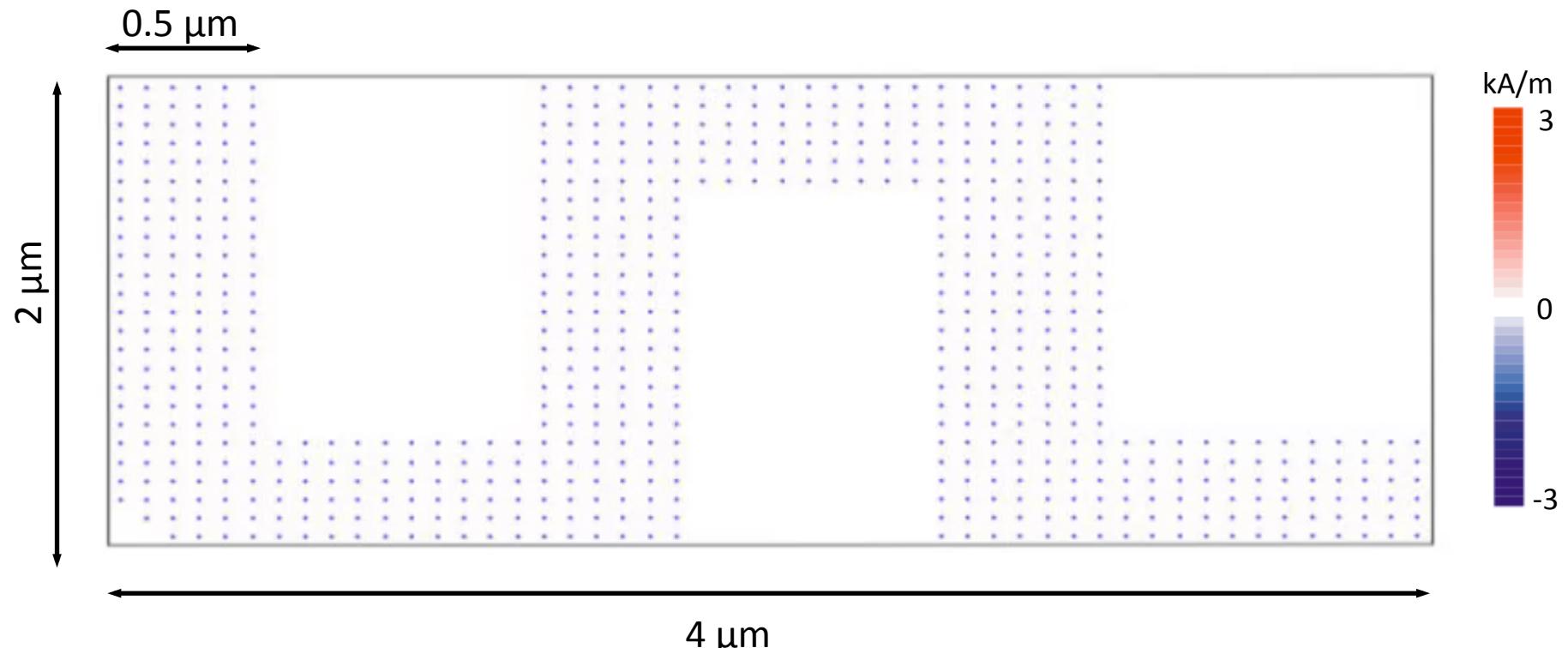
Spin-wave boundary mode



Spin-wave circuit

Thickness = 300 nm

Excitation frequency = 10 GHz



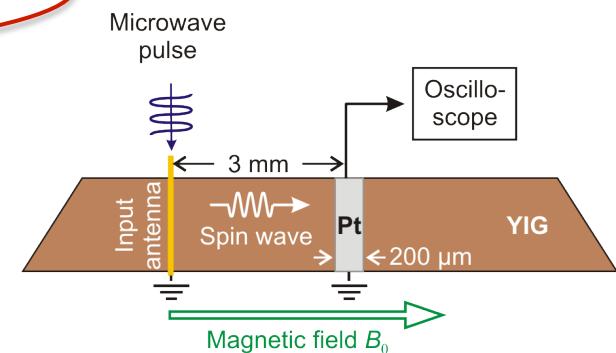
The boundary mode propagates around the corners
without changing its properties

I. New materials for magnonics

II. Novel means for magnon detection

III. Data processing using magnons

IV. Magnonic supercurrents



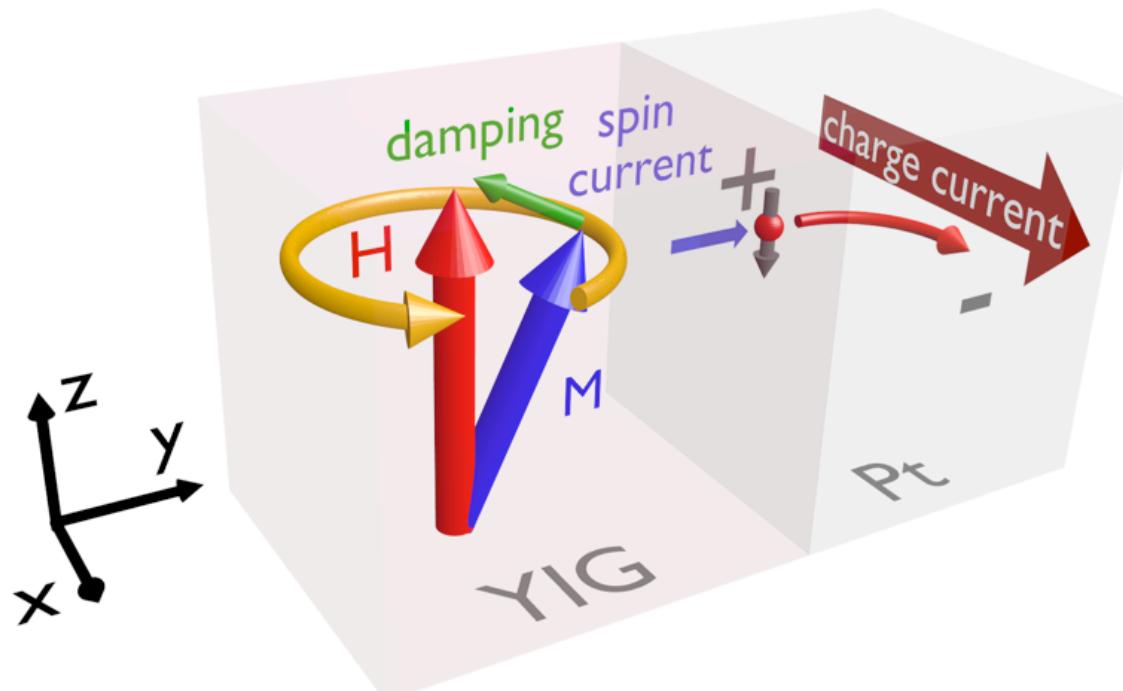
II. Novel means for magnon detection



Spin pumping + inverse spin Hall effect

Magnon to charge current conversion

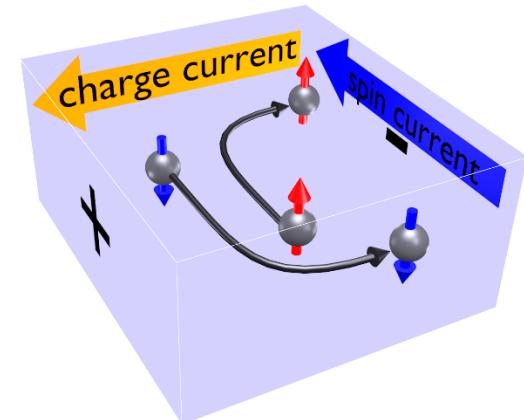
Detection of magnons by a combination of spin pumping and inverse spin Hall effect



Spin pumping

Tserkovnyak et al., PRL (2002)
Costache et al., PRL (2006)

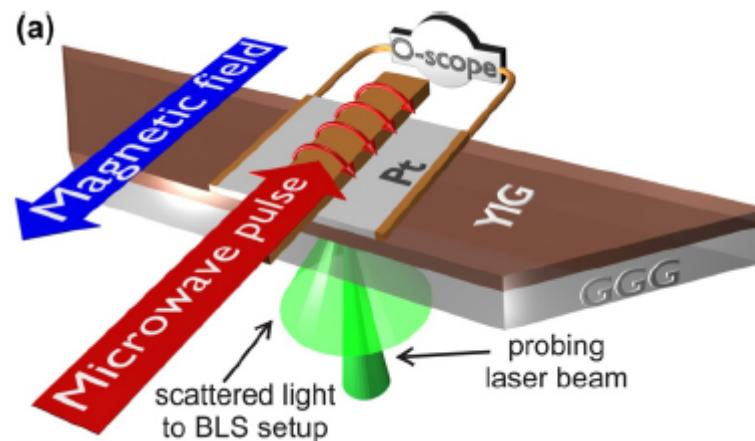
Inverse spin Hall effect (ISHE)



Hirsch, PRL (1999)
Saitoh et al., APL **88** 182509 (2006)

Time resolved ISHE voltage

Yttrium iron garnet (YIG) / platinum (Pt) bilayer was used



Parameters:

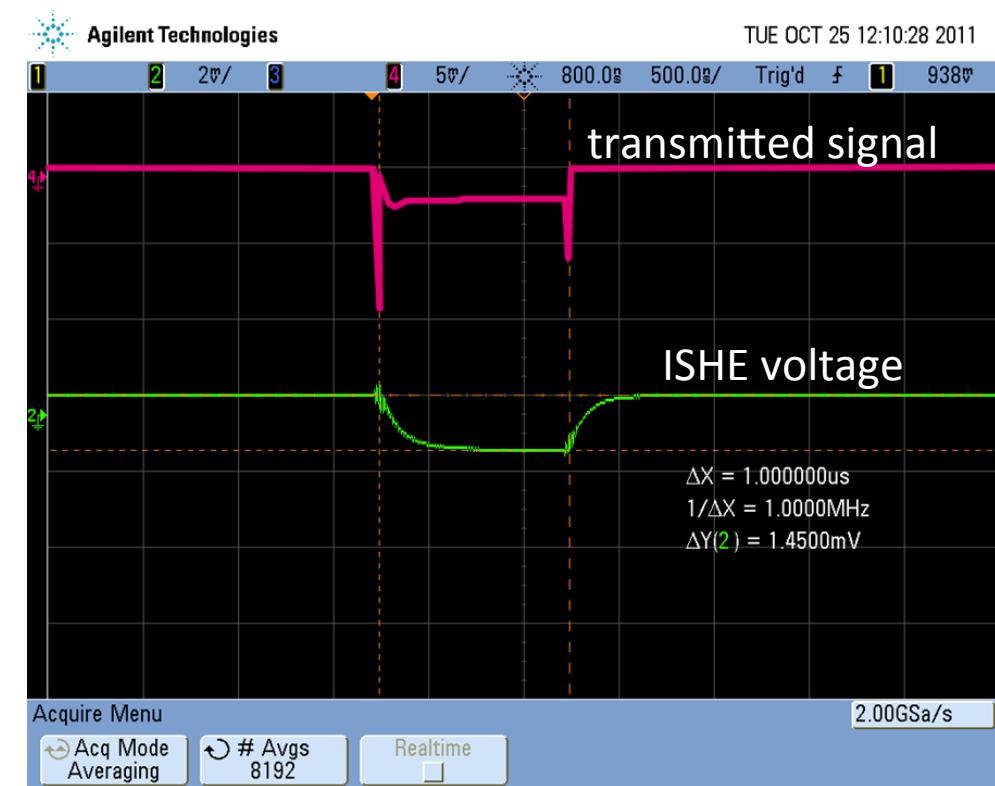
YIG thickness: $2.1\text{ }\mu\text{m}$

YIG/Pt width: 3 mm

Pt thickness: 10 nm

Magnetic field $B = -175.5\text{ mT}$

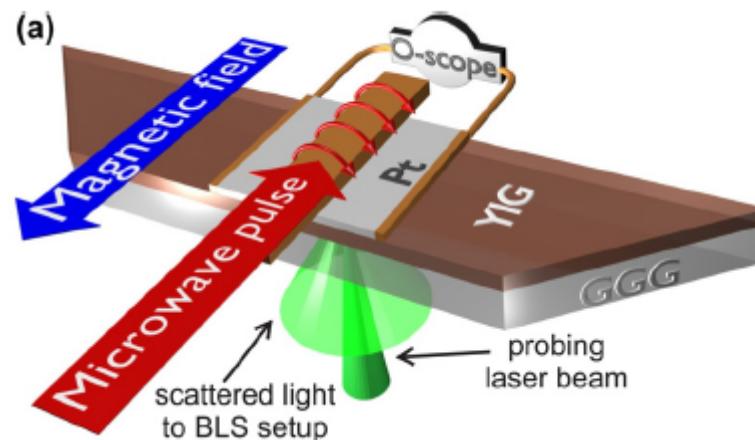
Microwave frequency 7 GHz



Jungfleisch et al., APL **99**, 182512 (2011)

Time resolved ISHE voltage

Yttrium iron garnet (YIG) / platinum (Pt) bilayer was used



Parameters:

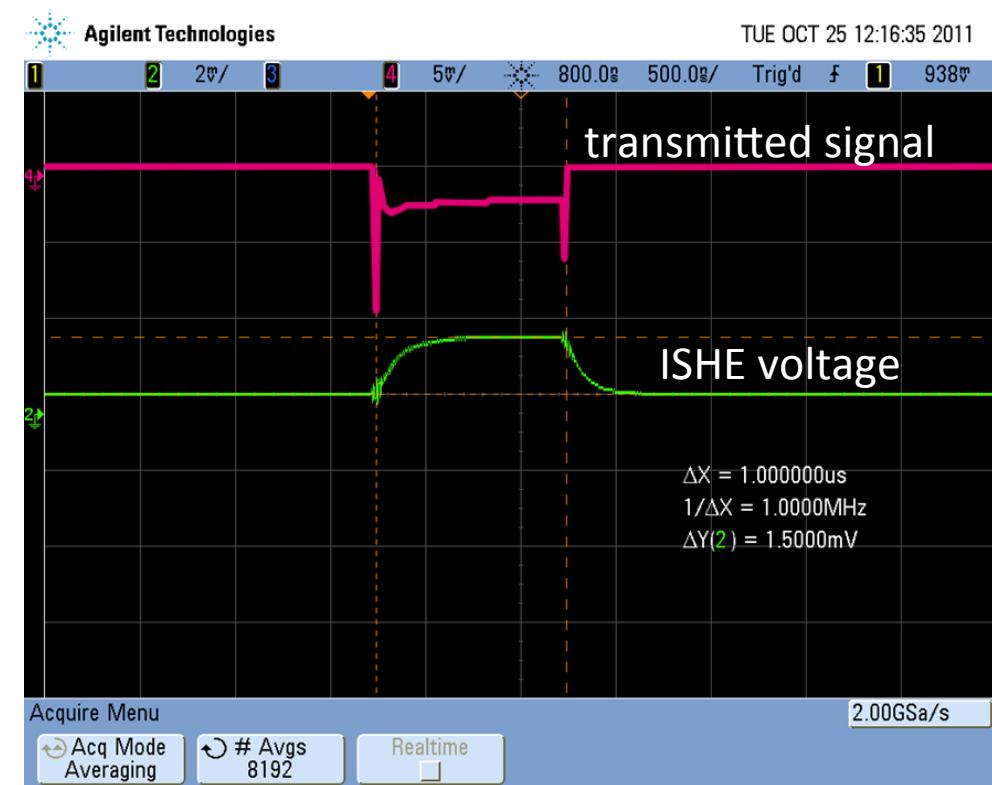
YIG thickness: $2.1\text{ }\mu\text{m}$

YIG/Pt width: 3 mm

Pt thickness: 10 nm

Magnetic field $B = +175.5\text{ mT}$

Microwave frequency 7 GHz

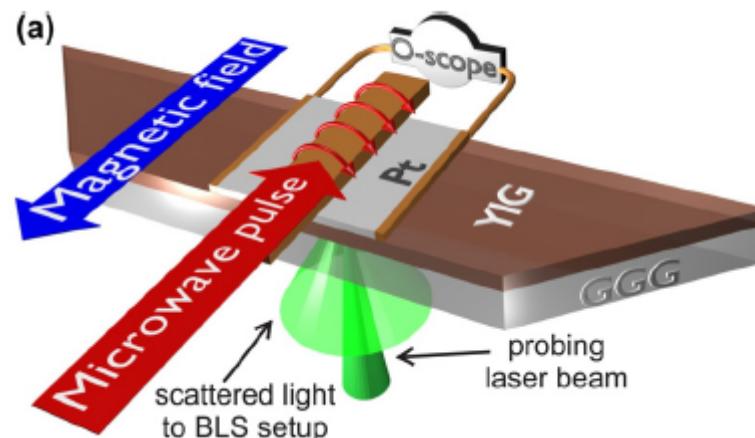


Jungfleisch et al., APL **99**, 182512 (2011)

Time resolved ISHE voltage

Yttrium iron garnet (YIG) / platinum (Pt) bilayer was used

Magnetic field $B = + 175.5$ mT
Microwave frequency 7 GHz

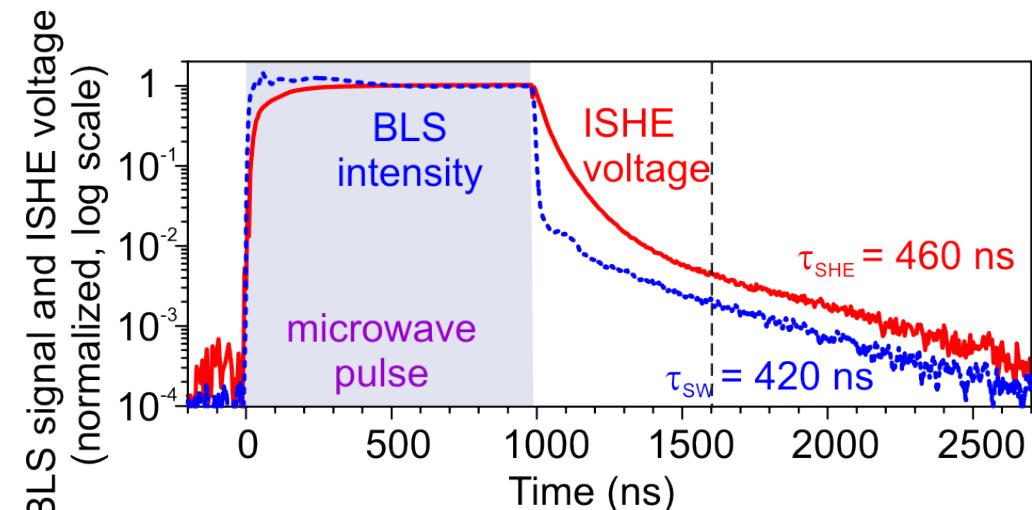


Parameters:

YIG thickness: 2.1 μm

YIG/Pt width: 3 mm

Pt thickness: 10 nm

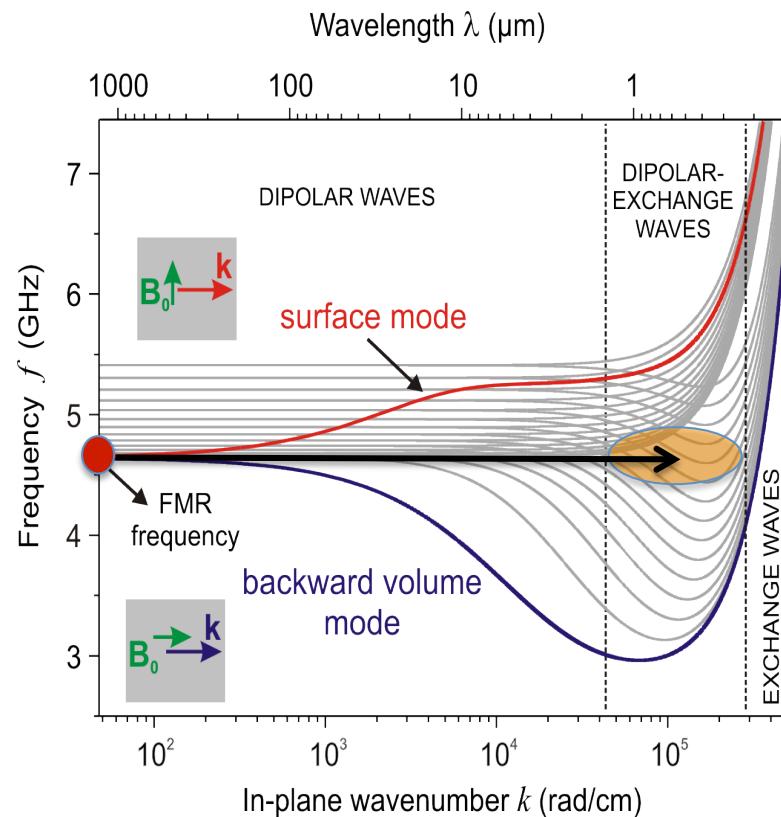


- ❖ ISHE pulse has a long rise and fall times
- ❖ Secondary magnons contribute to ISHE

Jungfleisch et al., APL 99, 182512 (2011)

Surface modes contribute to spin pumping

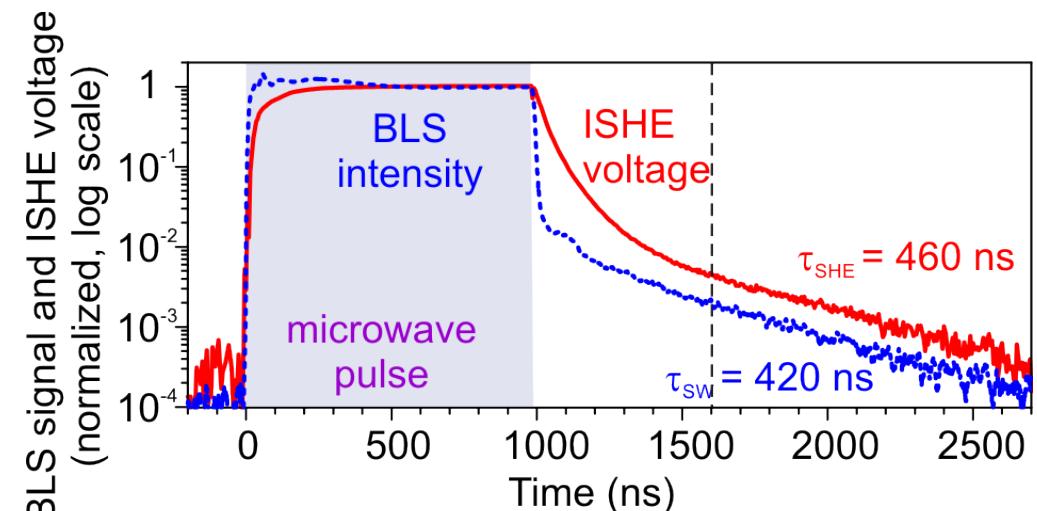
Yttrium iron garnet (YIG) / platinum (Pt) bilayer was used



Jungfleisch et al., APL 99, 182512 (2011)

Magnetic field $B = + 175.5$ mT

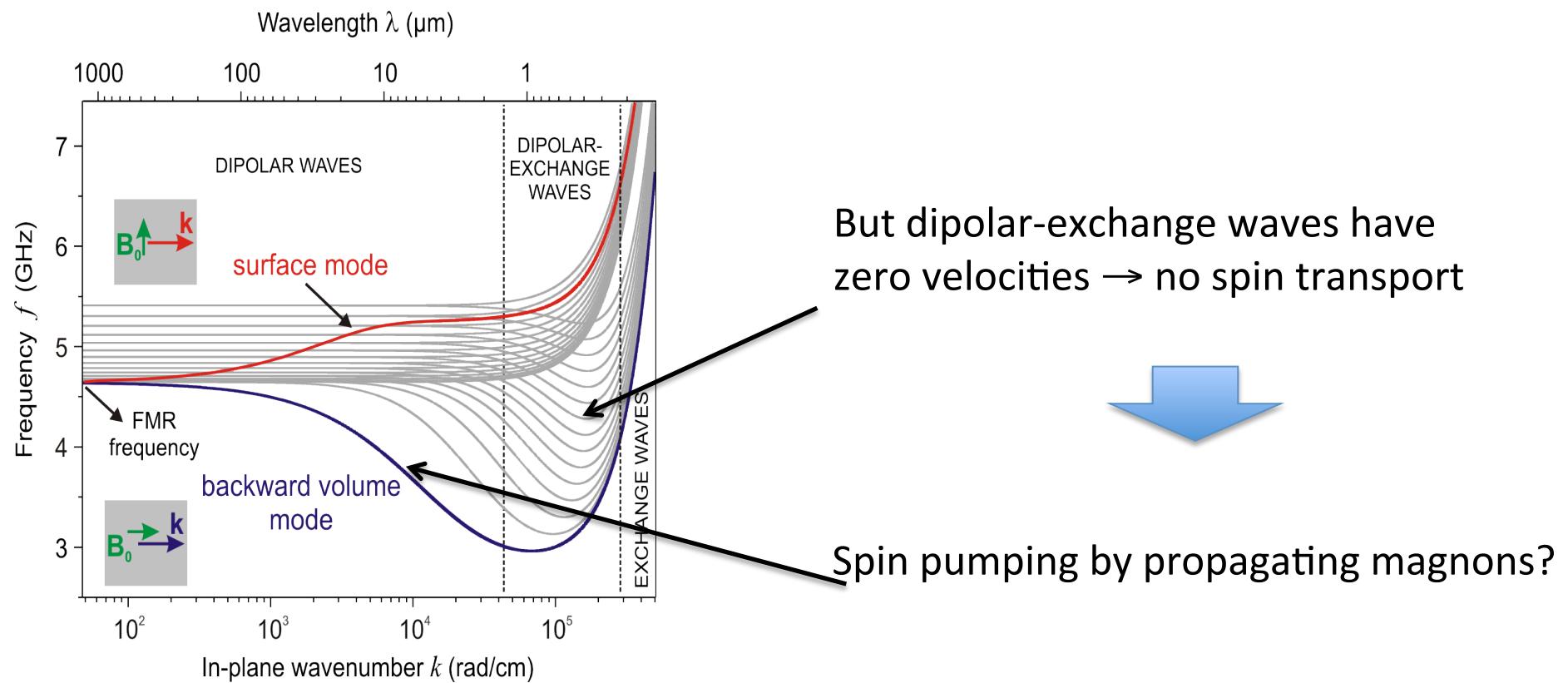
Dipolar-exchange waves
contribute to spin pumping!



- ❖ ISHE pulse has a long rise and fall times
- ❖ Secondary magnons contribute to ISHE

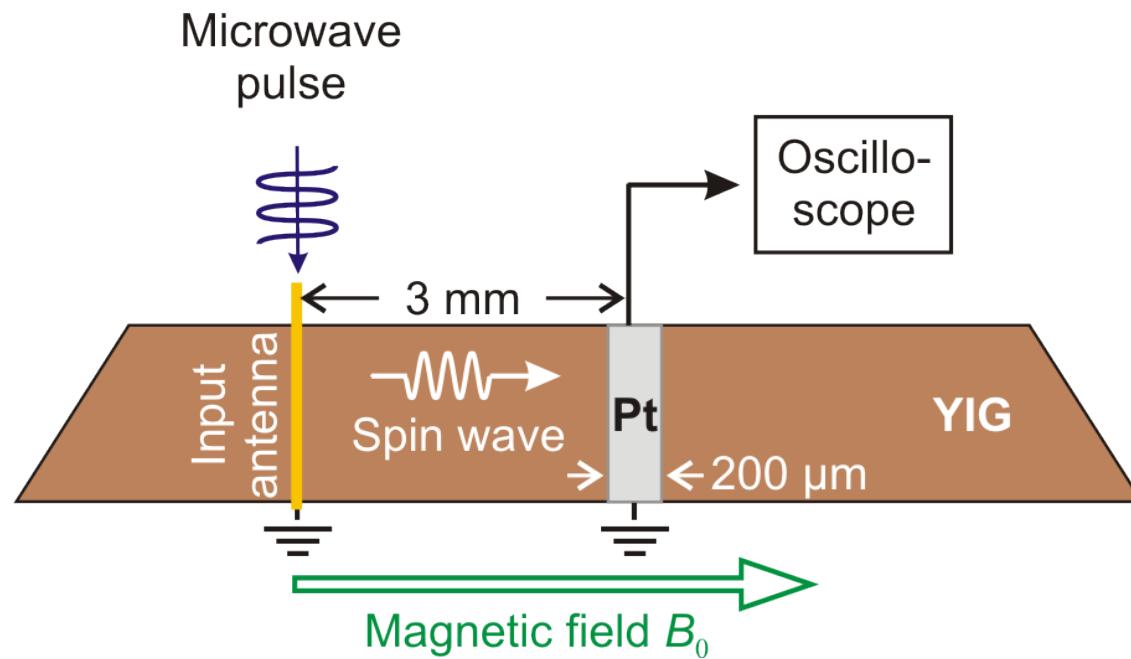
Surface modes contribute to spin pumping

Yttrium iron garnet (YIG) / platinum (Pt) bilayer was used



ISHE detection of propagating spin waves

Spin-wave source and Pt detector are separated in space



Parameters:

YIG thickness: 2.1 μm

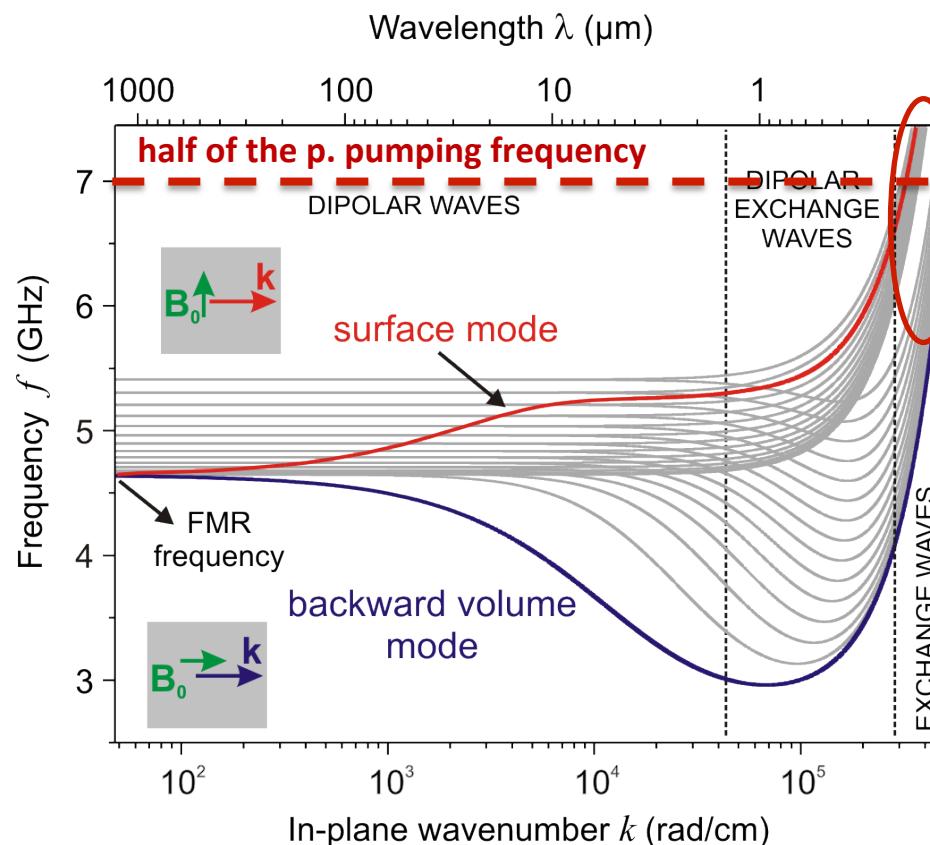
Pt size: 3 x 0.2 mm

Pt thickness: 10 nm

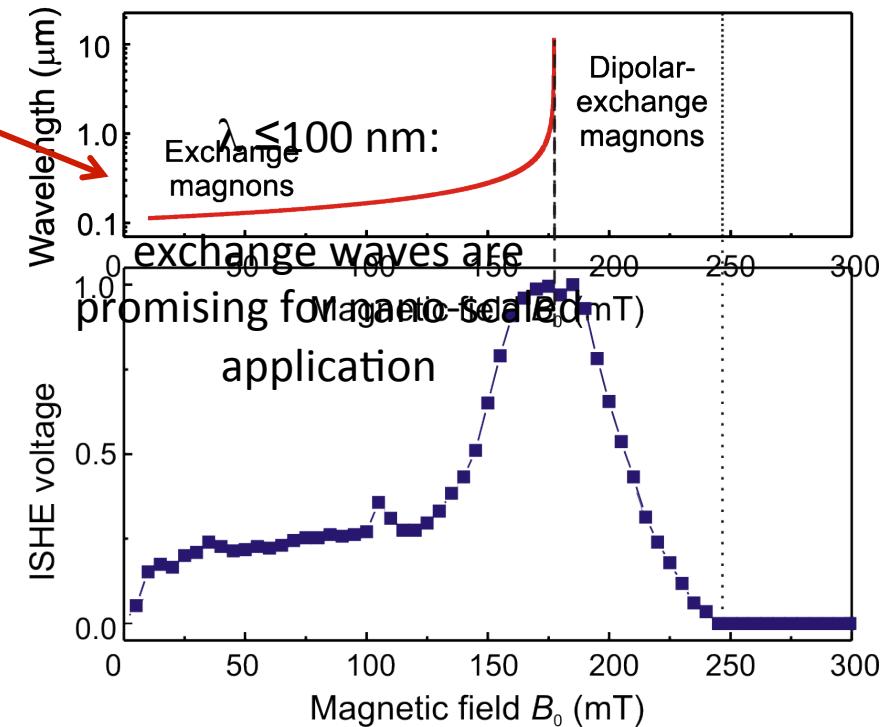
Signal frequency: 7 GHz

Spin pumping by exchange magnons

Parametric pumping at 14 GHz



Exchange magnons contribute to spin pumping!



Sandweg et al., PRL **106**, 216601 (2011)

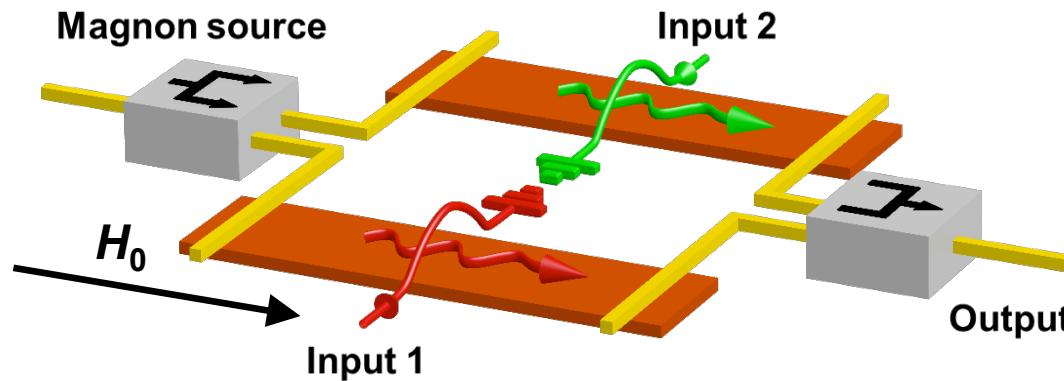
Kurebayashi, Dzyapko et al., APL **99**, 162502 (2011)
Ando et al., APL **99**, 092510 (2011)

III. Data processing using magnons

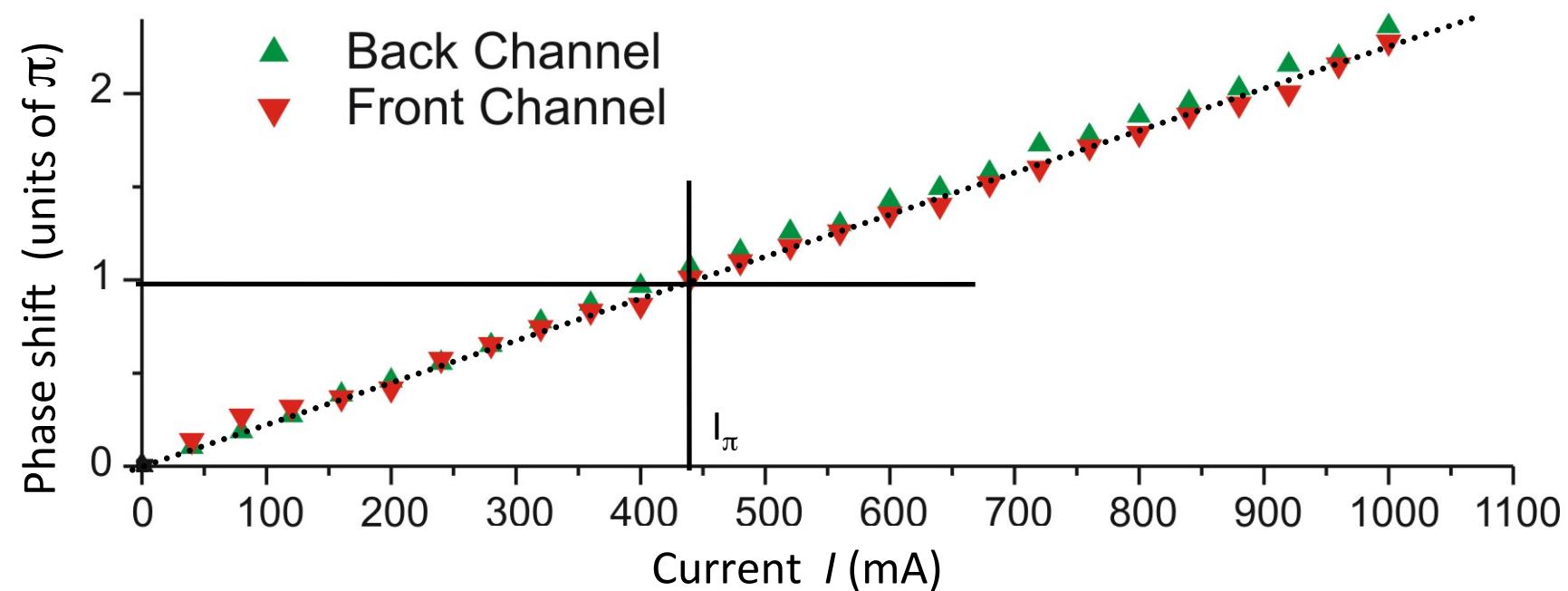


Spin-wave logic gates

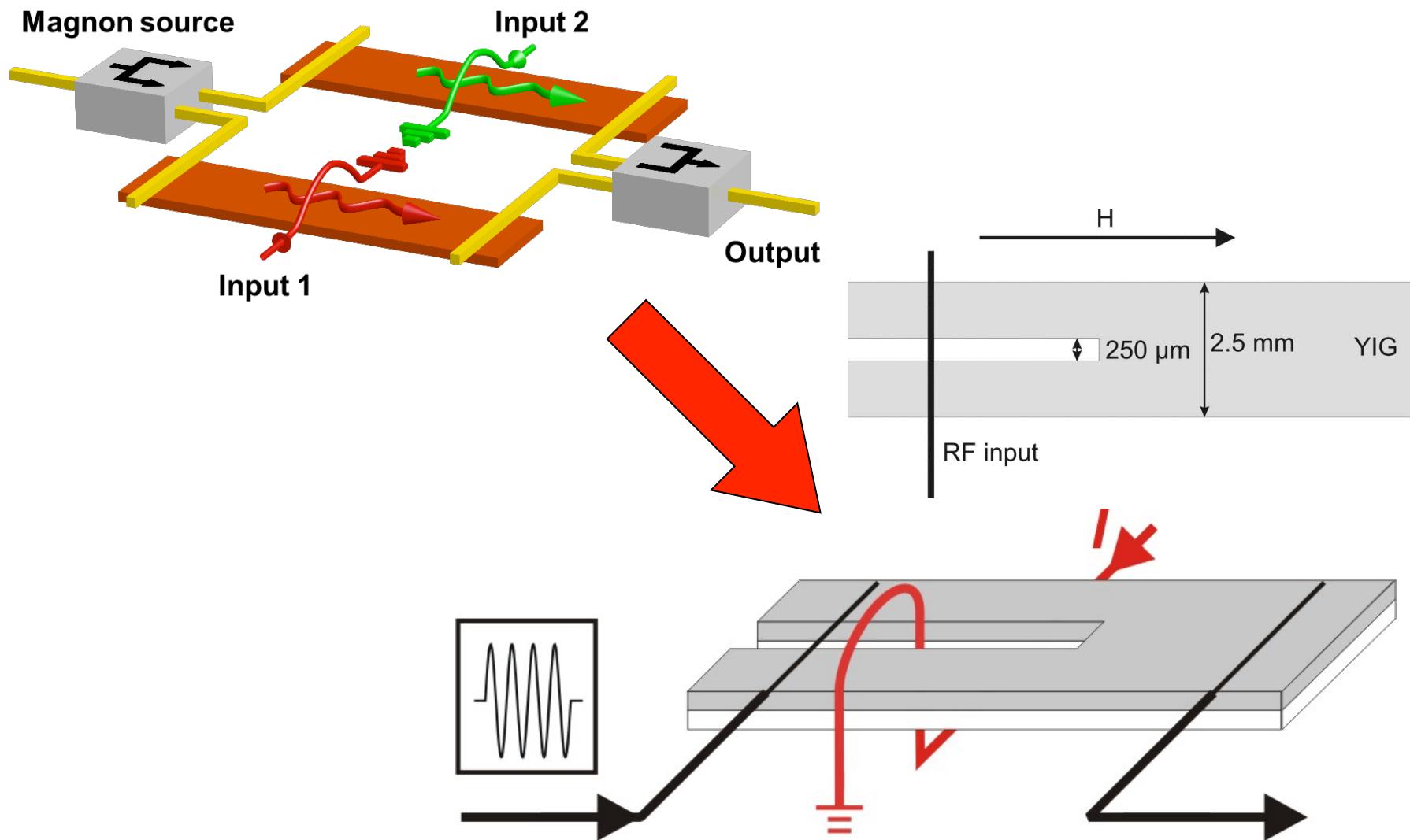
Mach-Zehnder interferometer based spin-wave logic gate



Kostylev et al., **APL** **87**, 153501 (2005)
 Schneider et al., **APL** **92**, 022505 (2008)

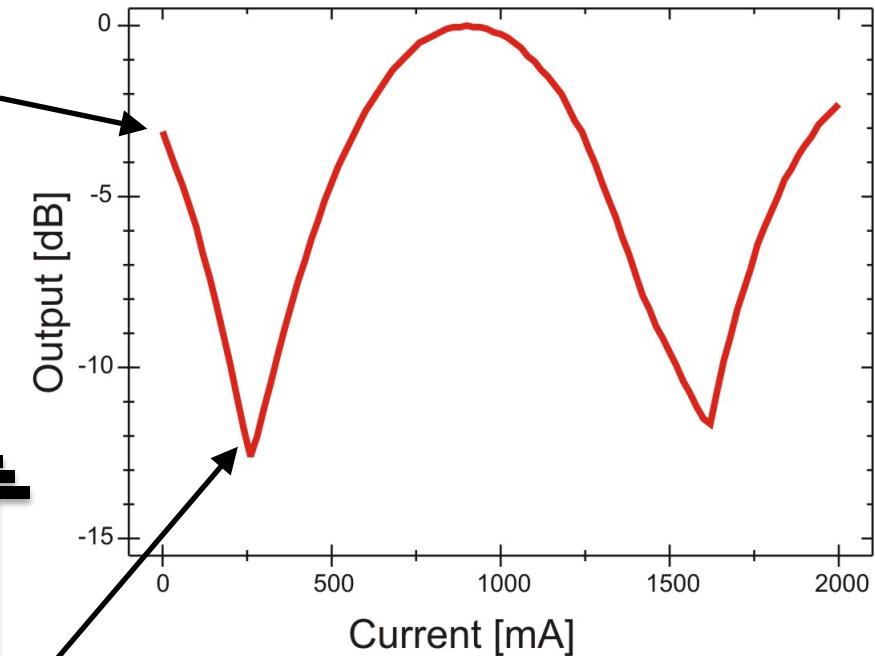
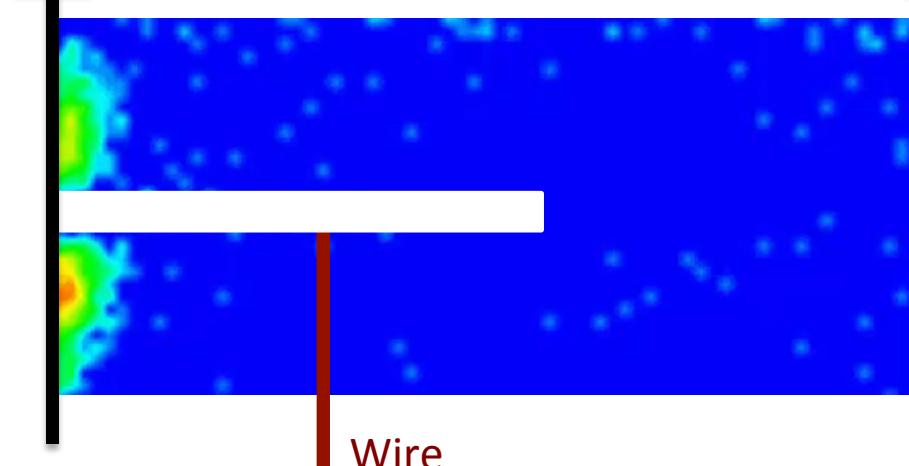


“Interferometer on a waveguide“

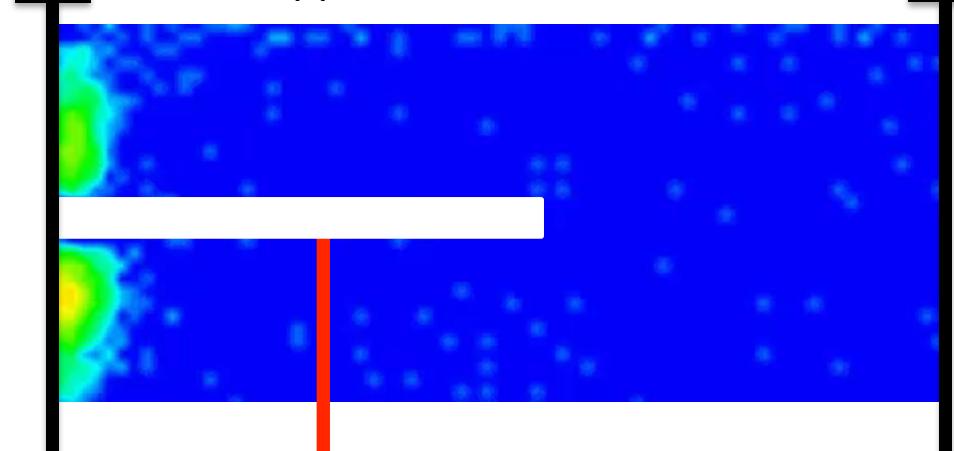


“Interferometer on a waveguide” – BLS measurement

No current applied to the wire

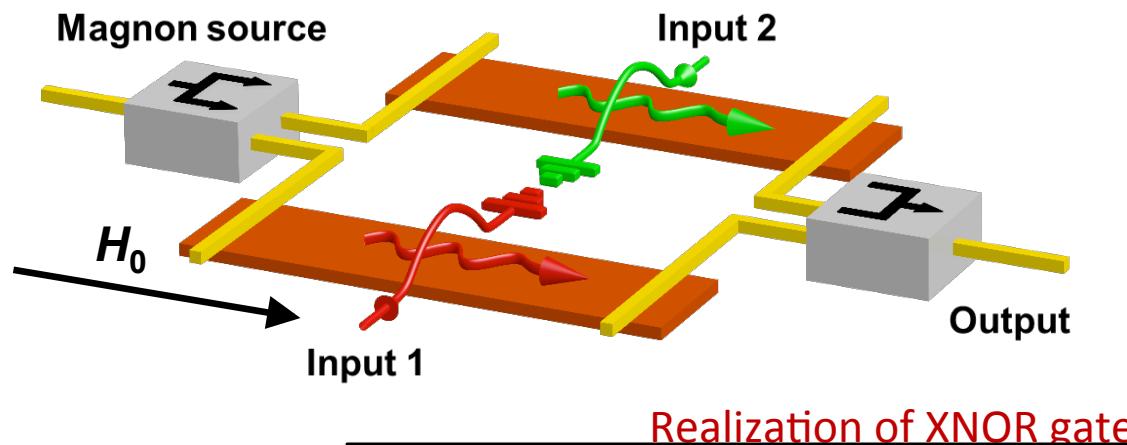


240 mA applied to the wire

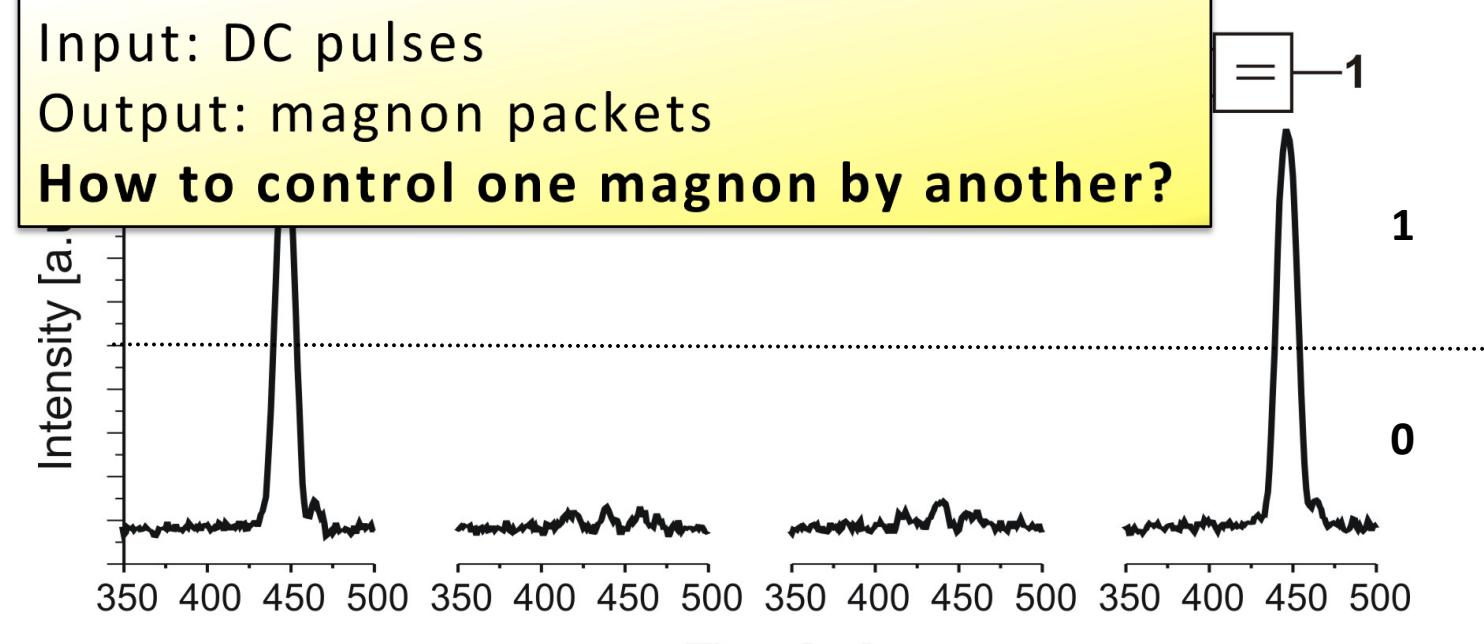


Schneider et al., J. of Nanoelectronics and
Optoelectronics 3, 69 (2008)

Mach-Zehnder interferometer based spin-wave logic gate



Inputs		Output
A (I_1)	B (I_2)	
0 (0)	0 (0)	1
0 (0)	1 (I_π)	0
1 (I_π)	0 (0)	0
1 (I_π)	1 (I_π)	1



Schneider et al., APL 92, 022505 (2008)

Burkard Hillebrands

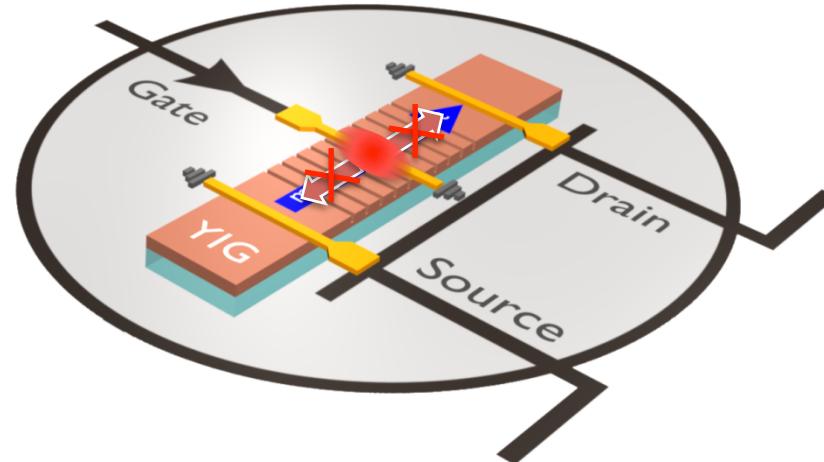
IEEE Magnetics Society Summer School - Minneapolis

June 14-19, 2015

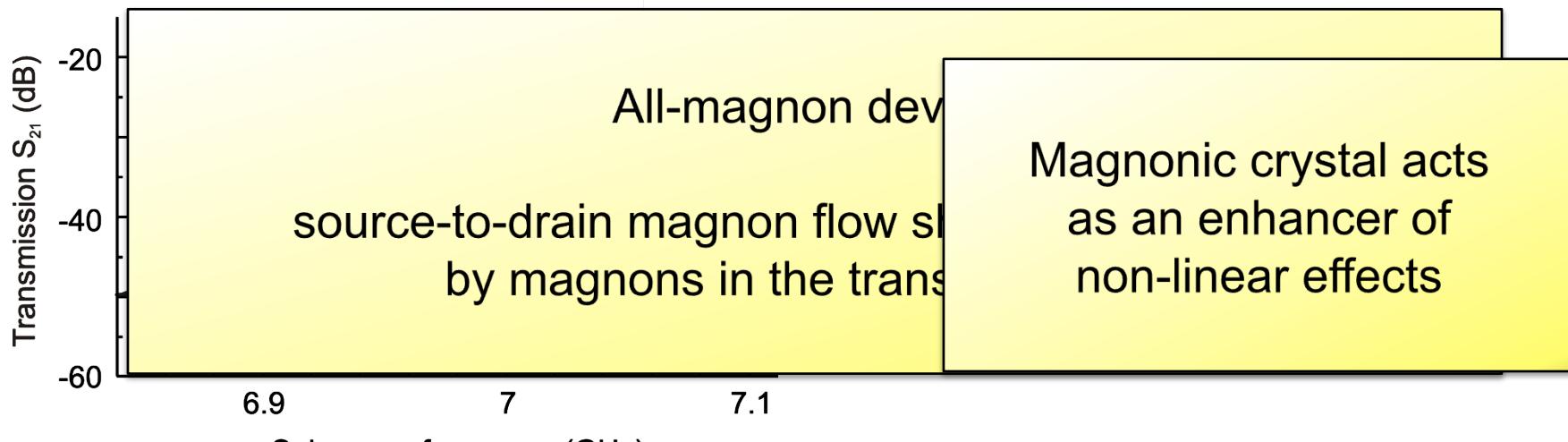
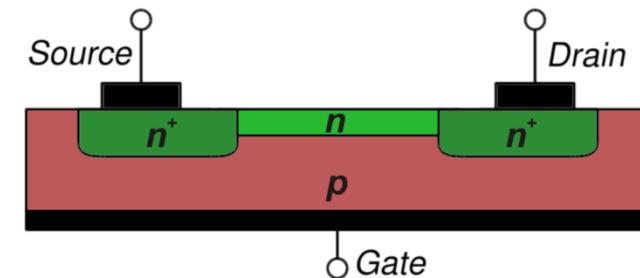
III. Data processing using magnons

- Spin-wave logic gates
- Magnon transistor

Magnon transistor



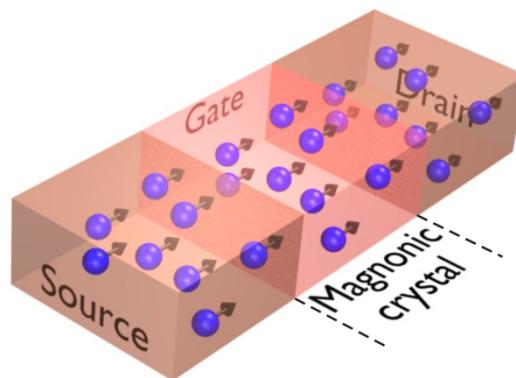
Semiconductor field-effect transistor:



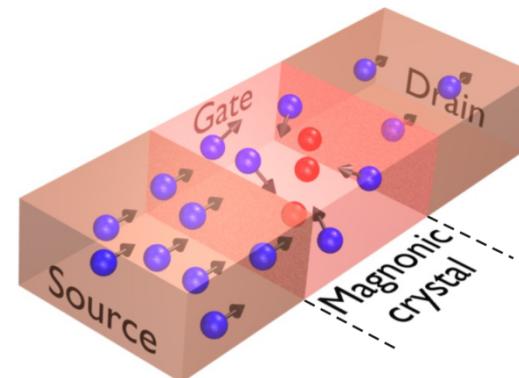
A.V. Chumak et al., Nat. Commun. 5:4700 (2014)

Magnon transistor

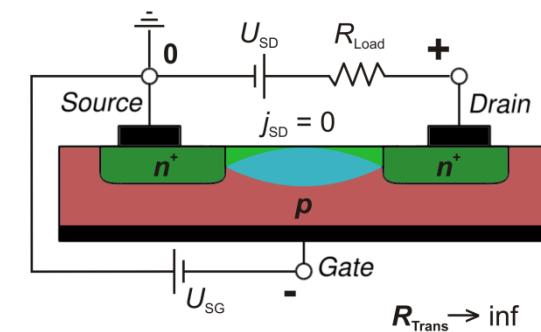
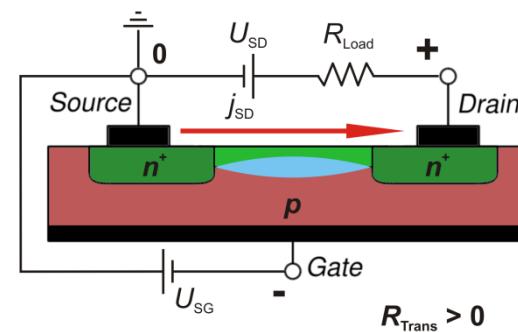
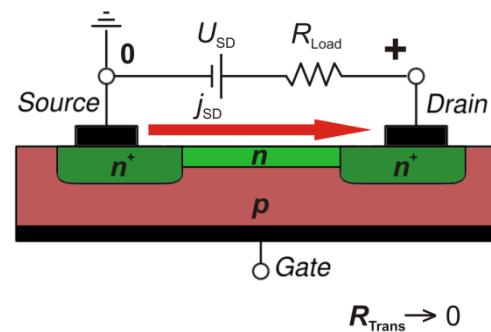
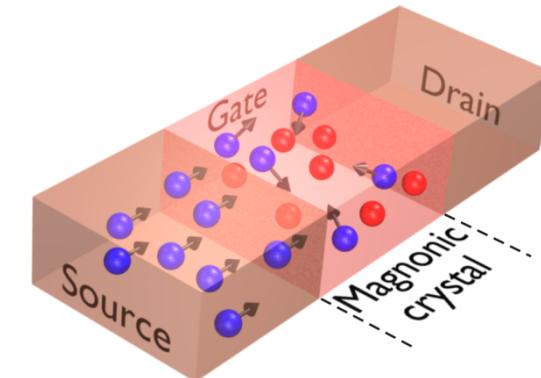
Opened: $R \rightarrow 0$
 Gate magnon density
 $n_G = 0$



Semi-closed: $R > 0$
 Gate magnon density
 $n_G > 0$



Closed: $R \rightarrow \infty$
 Gate magnon density
 $n_G \gg 0$

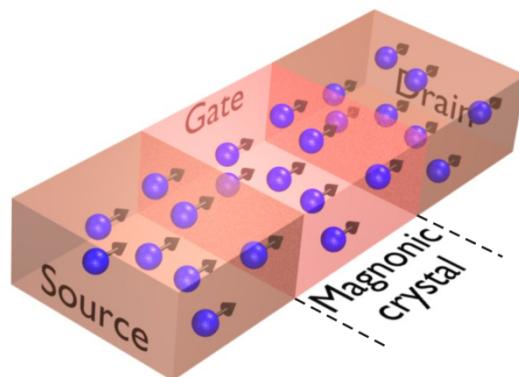


Magnon transistor

Opened: $R \rightarrow 0$

Gate magnon density

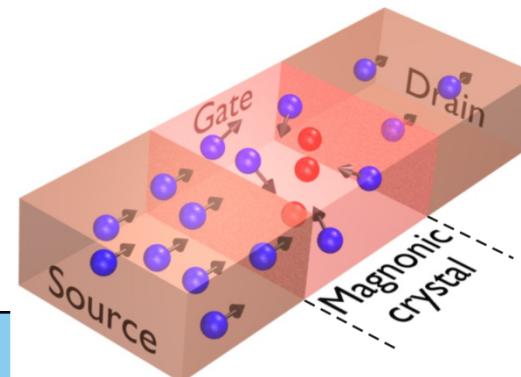
$$n_G = 0$$



Semi-closed: $R > 0$

Gate magnon density

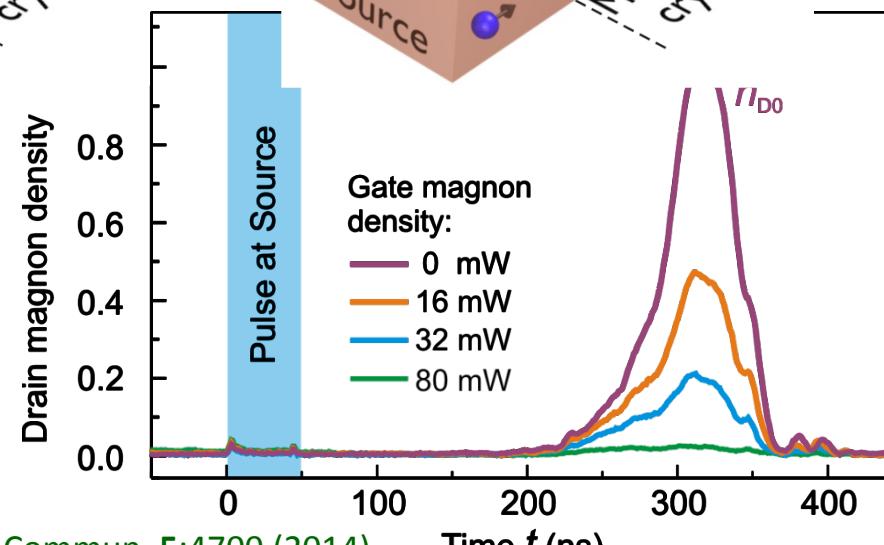
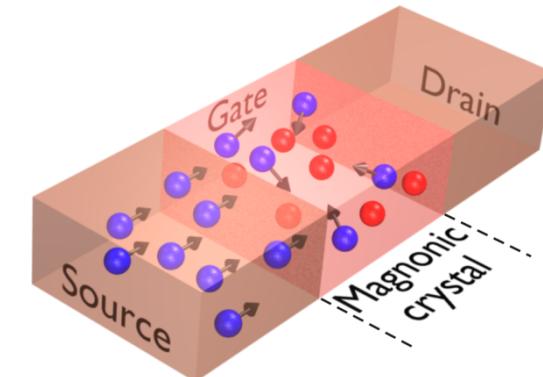
$$n_G > 0$$



Closed: $R \rightarrow \infty$

Gate magnon density

$$n_G \gg 0$$



A.V. Chumak et al., Nat. Commun. 5:4700 (2014)

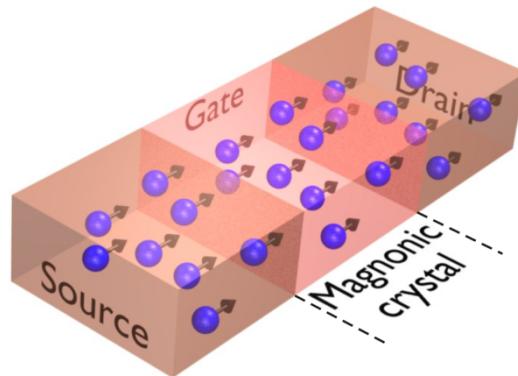
Burkard Hillebrands

IEEE Magnetics Society Summer School - Minneapolis

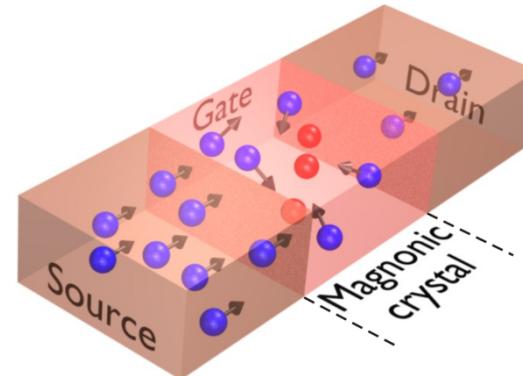
June 14-19, 2015

Magnon transistor

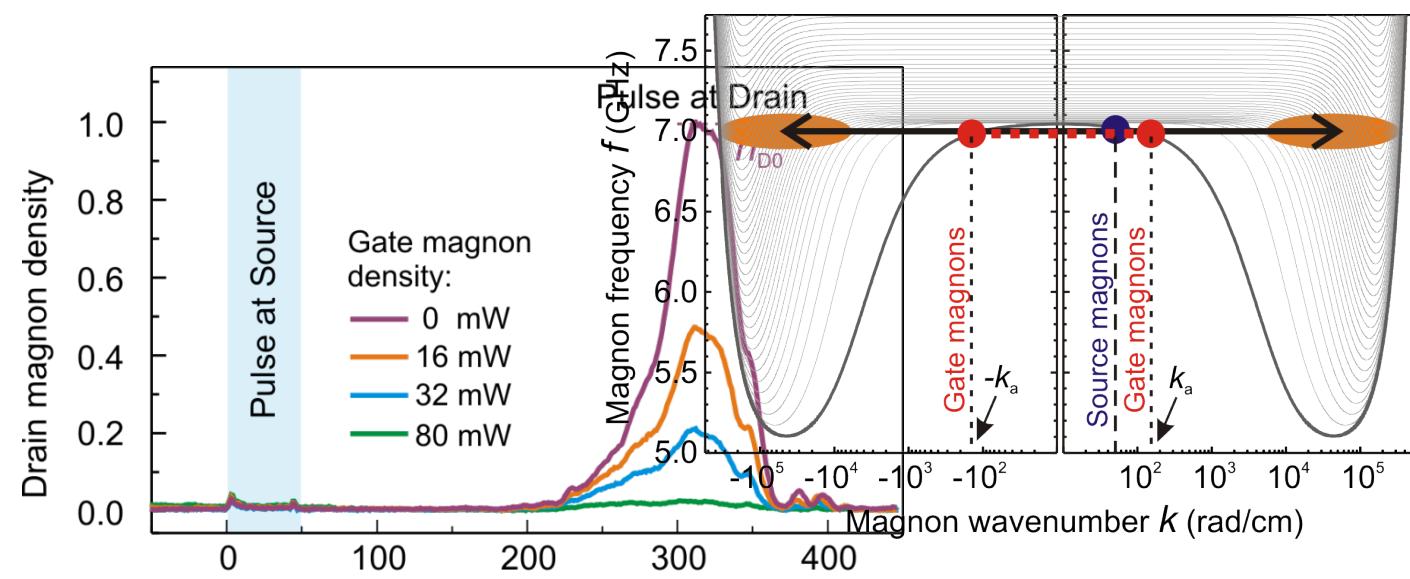
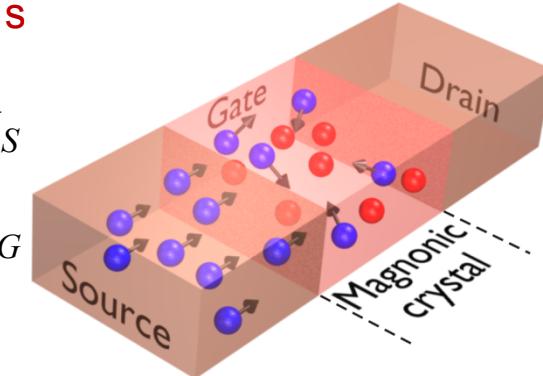
Opened: $R \rightarrow 0$



Semi-closed: $R > 0$



Closed: $R \rightarrow \infty$



A.V. Chumak et al., Nat. Commun. 5:4700 (2014)

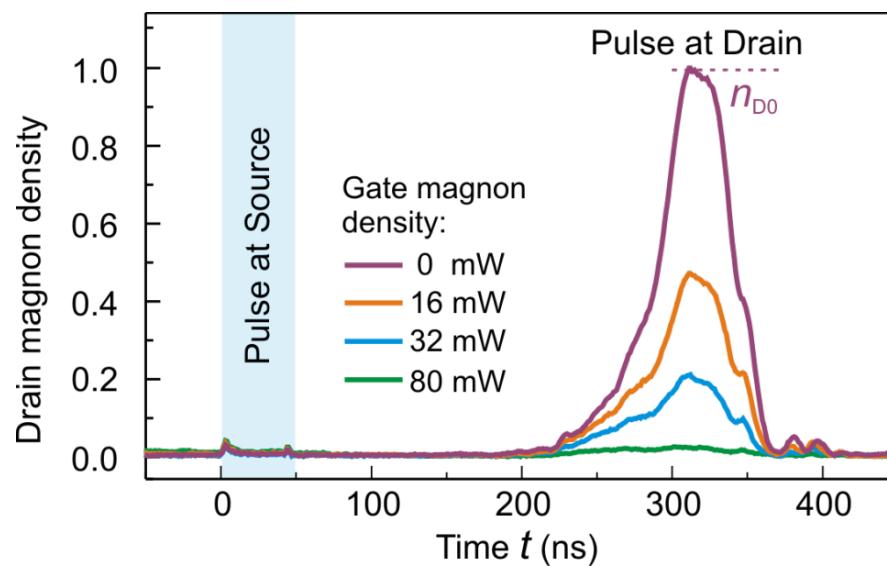
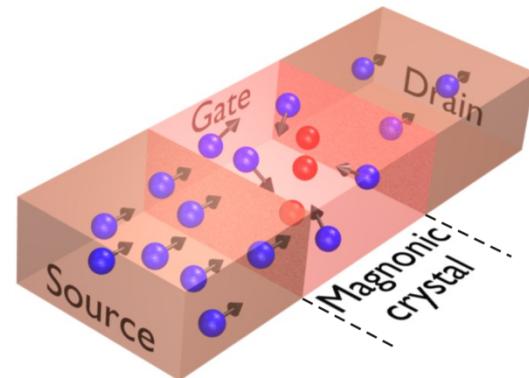
Time t (ns)

Burkard Hillebrands

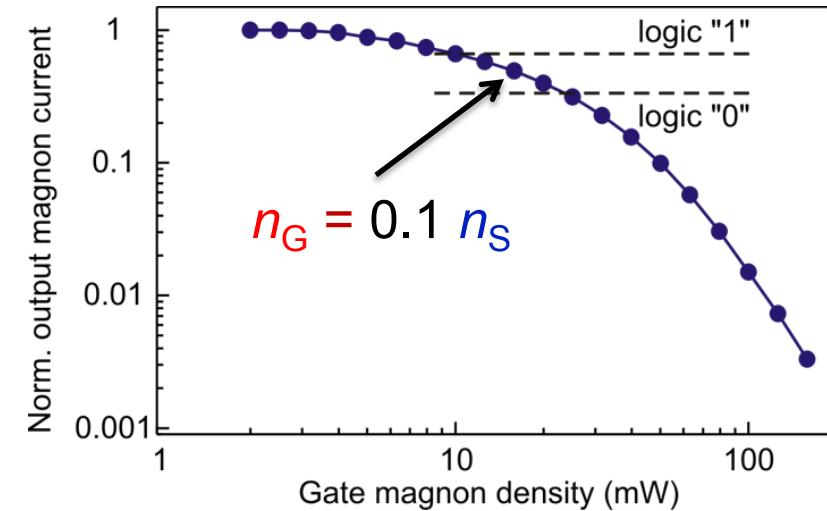
IEEE Magnetics Society Summer School - Minneapolis

June 14-19, 2015

Magnon transistor



A.V. Chumak et al., Nat. Commun. 5:4700 (2014)



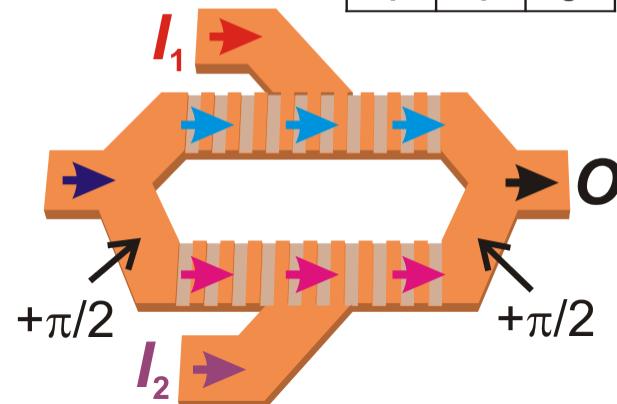
“magnon control by magnon” principle was realized:
data can be processed on the same magnetic chip

XOR logic gate



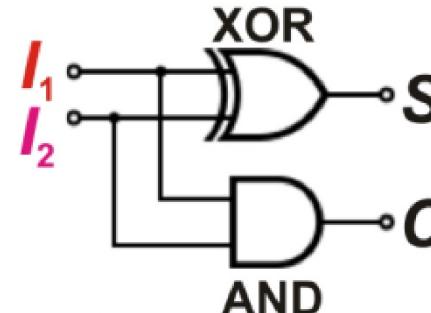
XOR

I ₁	I ₂	O
0	0	0
0	1	1
1	0	1
1	1	0

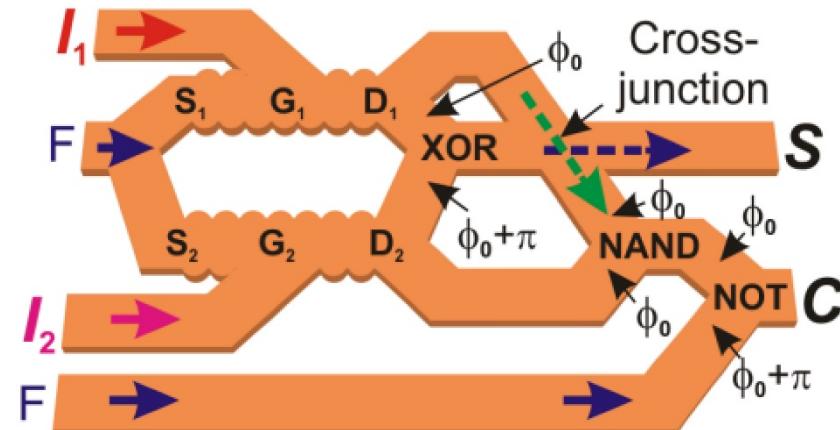


XOR gate requires 2 transistors
instead of 8 FET in CMOS

Half adder



I ₁	I ₂	C	S
0	0	0	0
1	0	0	1
0	1	0	1
1	1	1	0



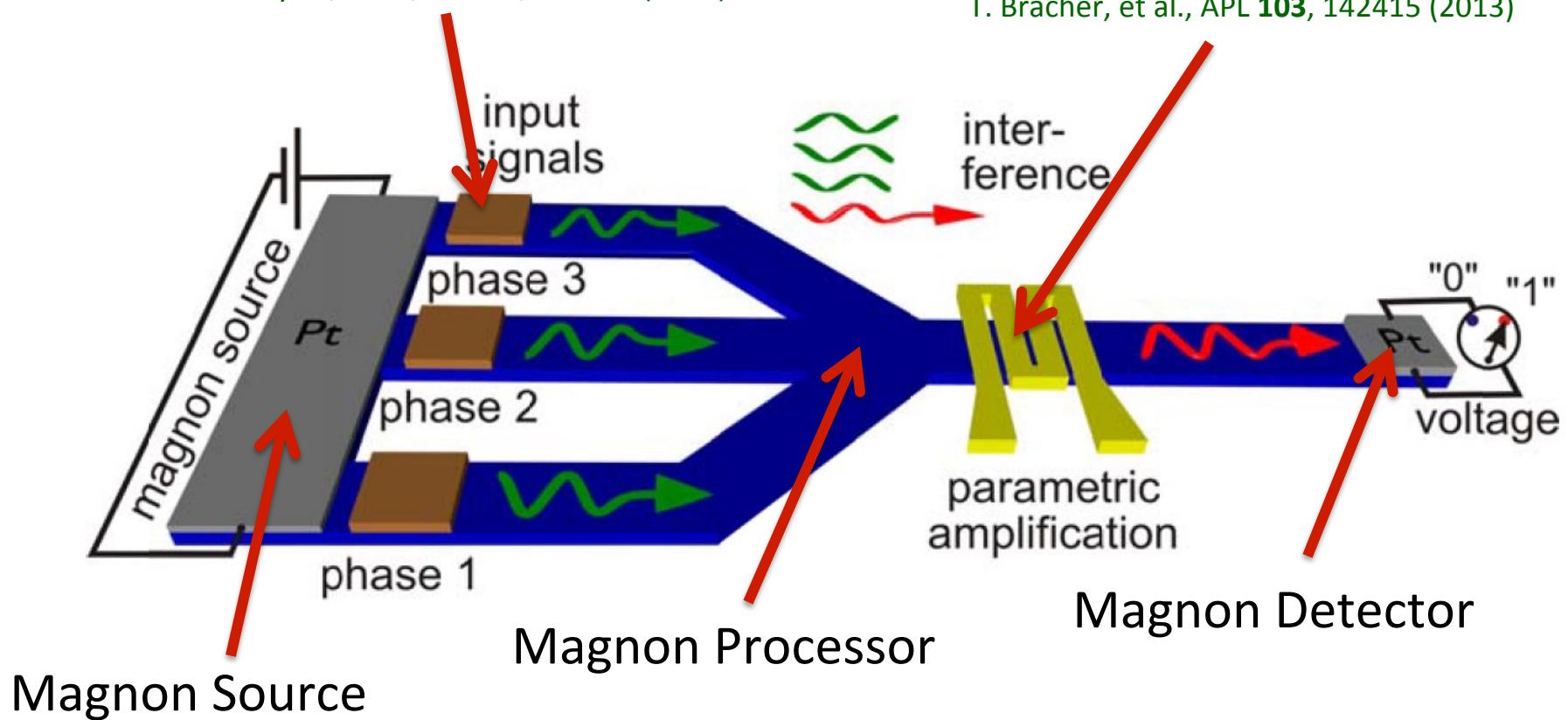
III. Data processing using magnons

- Spin-wave logic gates
- Magnon transistor
- Magnon majority gate

Majority gate

Phase Shifter

T. Neumann, et al. APL **94**, 042503 (2009)
M.P. Kostylev, et al., PRB **76**, 184419 (2007)



Parametric Amplifier

T. Brächer, et al. APL **104**, 092418 (2014)
T. Brächer, et al., APL **103**, 142415 (2013)

Design of a spin-wave majority gate employing mode selection

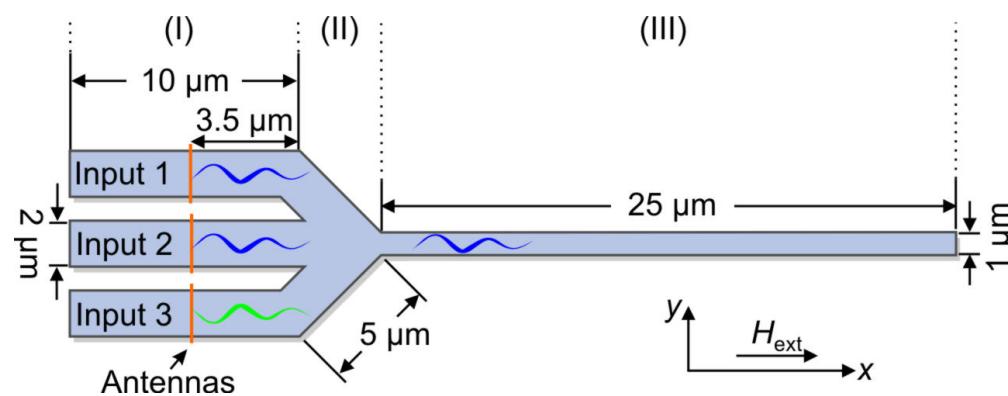
S. Klingler,^{1,*} P. Pirro,¹ T. Brächer,¹ B. Leven,¹ B. Hillebrands,¹ and A. V. Chumak¹

¹*Fachbereich Physik and Landesforschungszentrum OPTIMAS,
Technische Universität Kaiserslautern, 67663 Kaiserslautern, Germany*

(Dated: August 15, 2014)

Majority gate design:

100 nm YIG, [Pirro et al., APL 104, 012402 \(2014\)](#)



Truth-table:

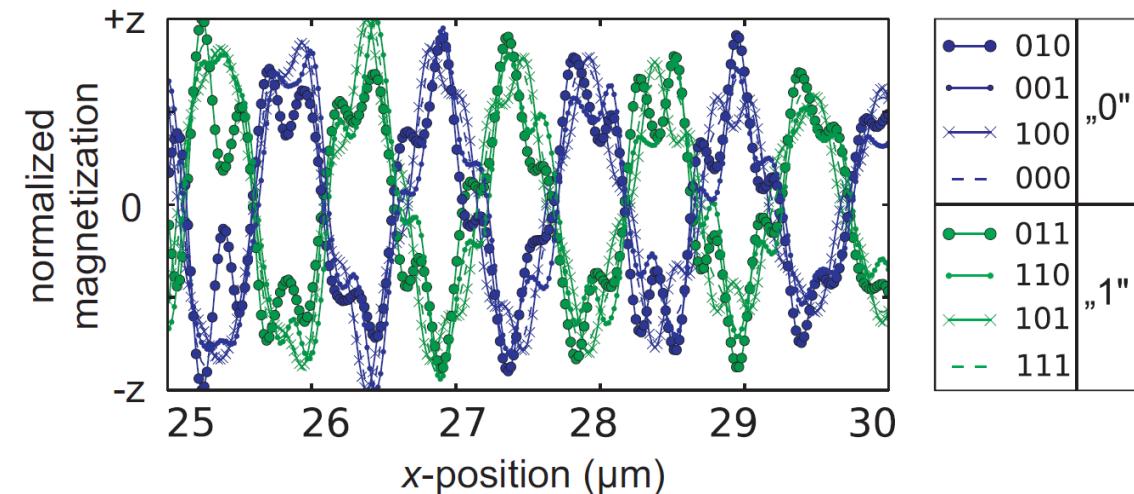
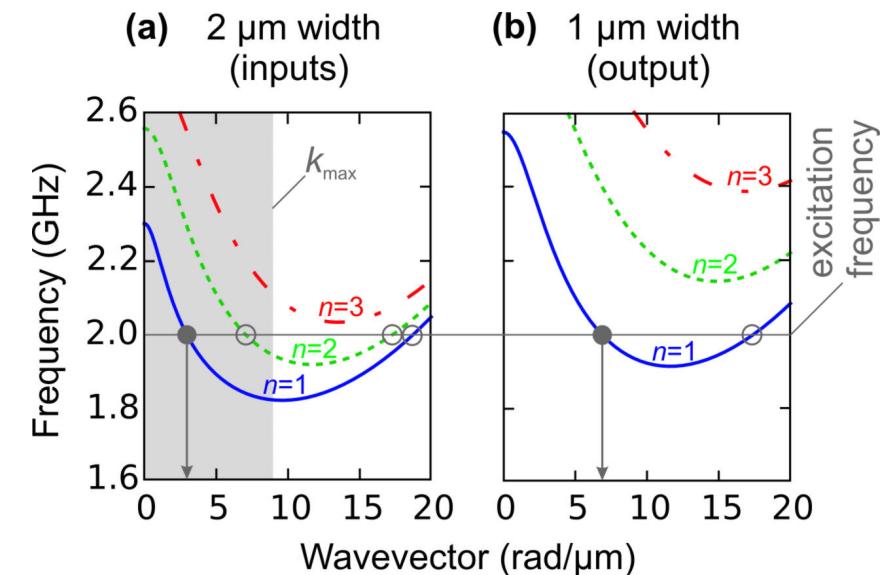
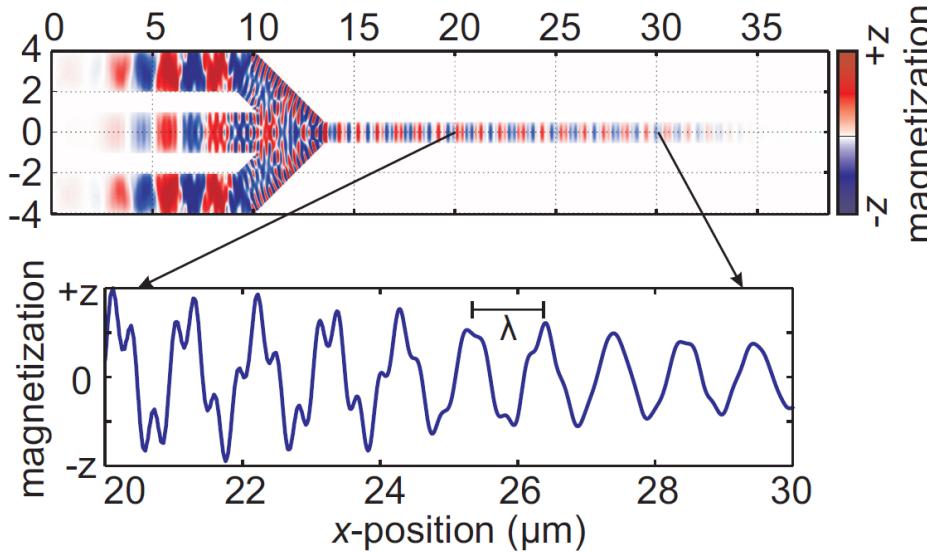
Input 1 Signal	Input 2 Signal	Input 3 Control	Output
0	0	0	0
1	0	0	0
0	1	0	0
1	1	0	1
AND			}
0	0	1	0
1	0	1	1
0	1	1	1
1	1	1	1
OR			

In CMOS **56 transistors** (3 NOT, 4 AND, 3 OR) are needed for majority gate

Data is coded into spin-wave phase

S. Klingler et al., arXiv: 1408.3235

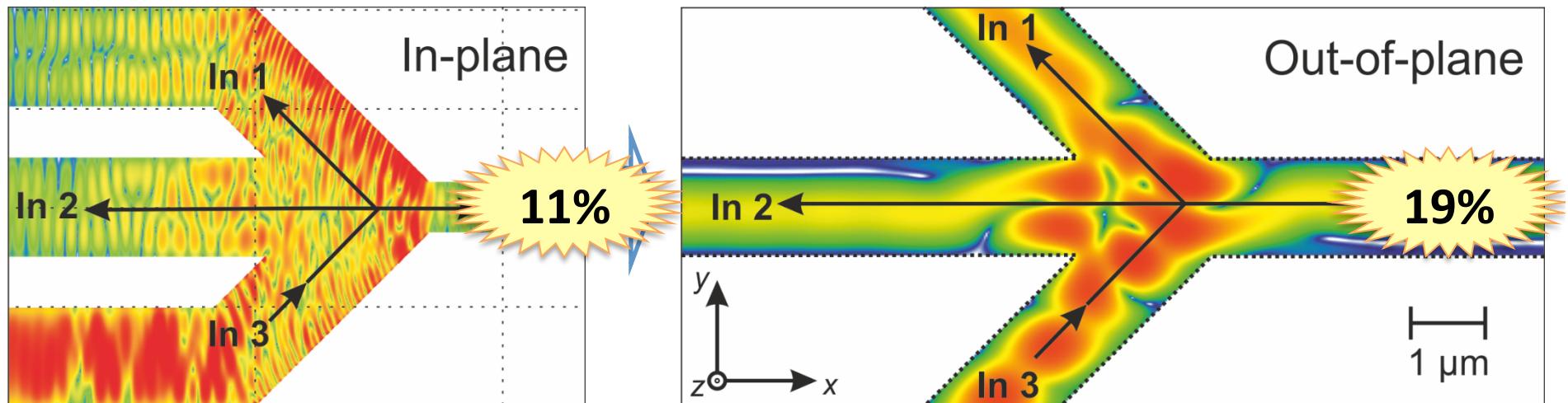
Majority gates: Simulations



Majority gates: Out-of-plane magnetization

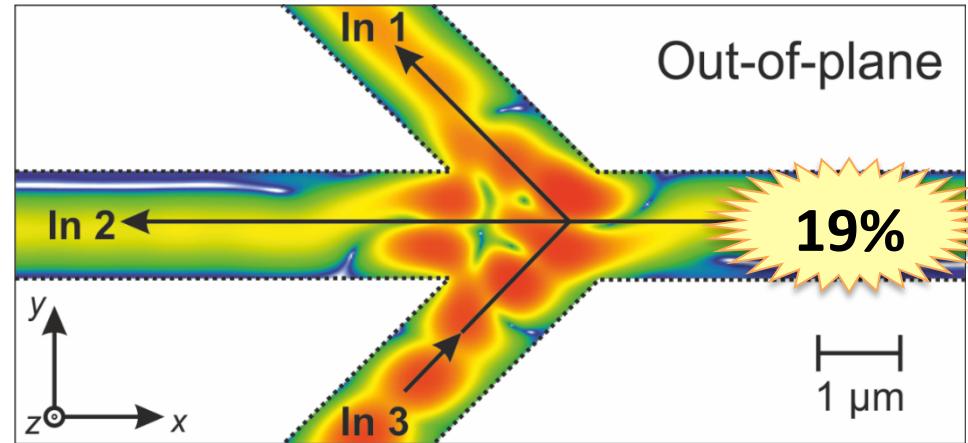
Switch from the in-plane magnetization to the out-of-plane
(see e.g. T. Schwarze, et al., Phys. Rev. B 85, 134448 (2012))

Forward volume spin waves are isotropic (always transverse to field)

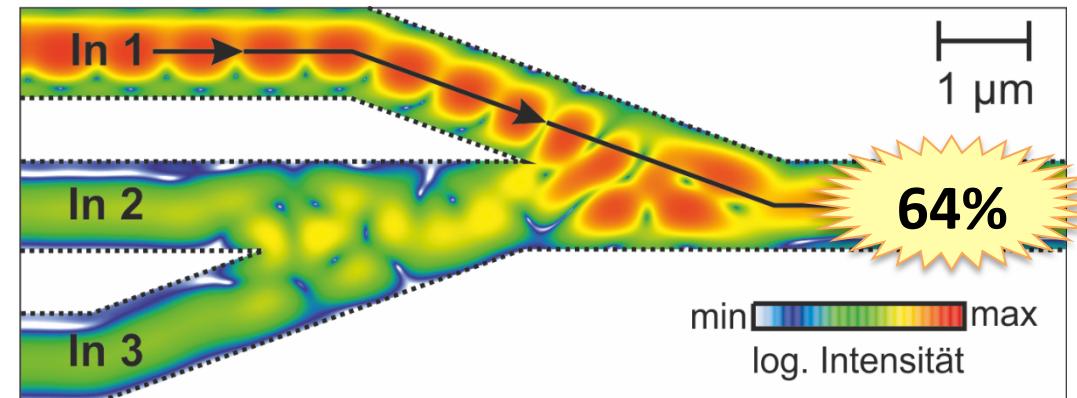


but no parasitic generation of high-wavenumber magnons

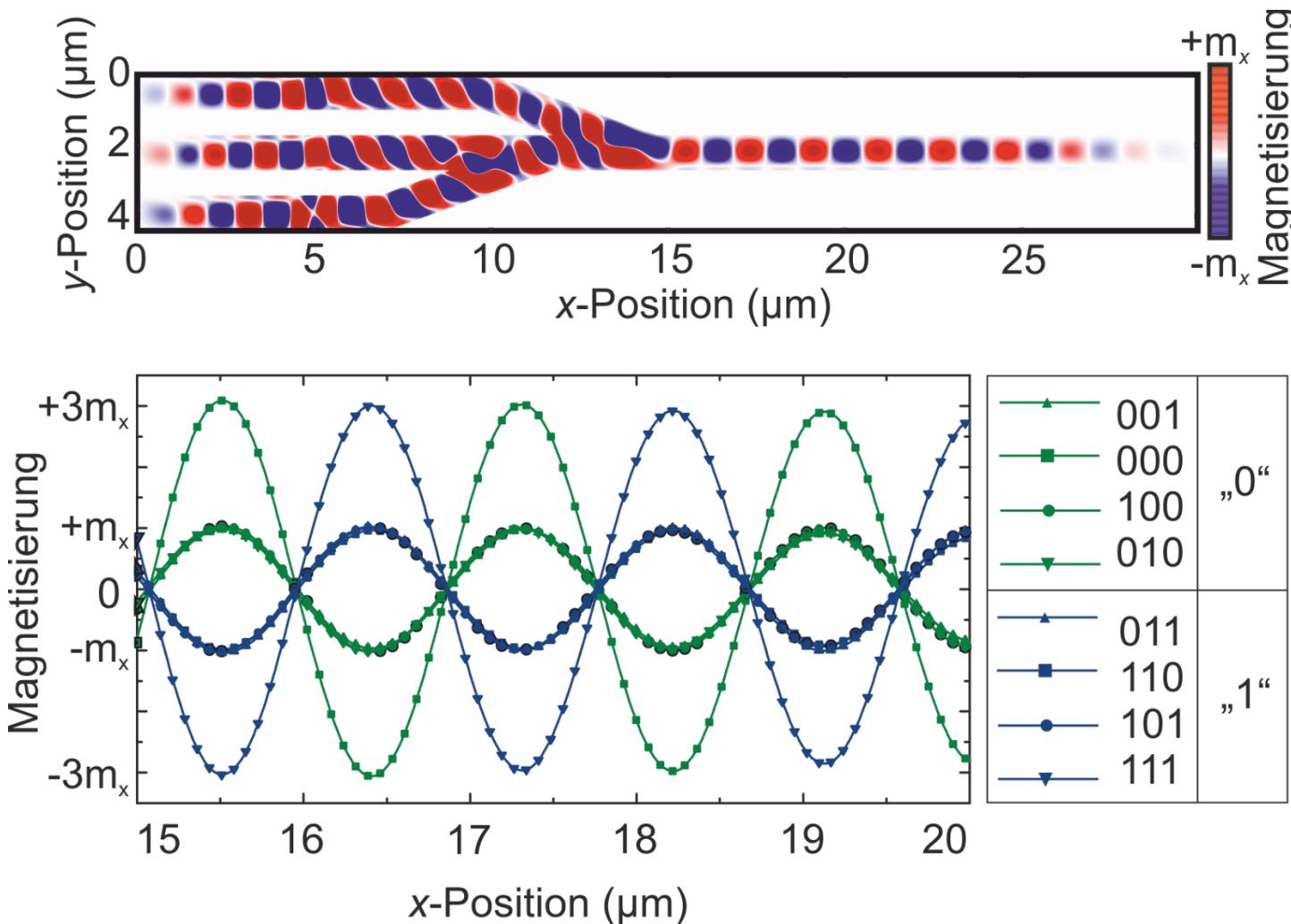
Majority gates: Out-of-plane magnetization



How to increase
this value?



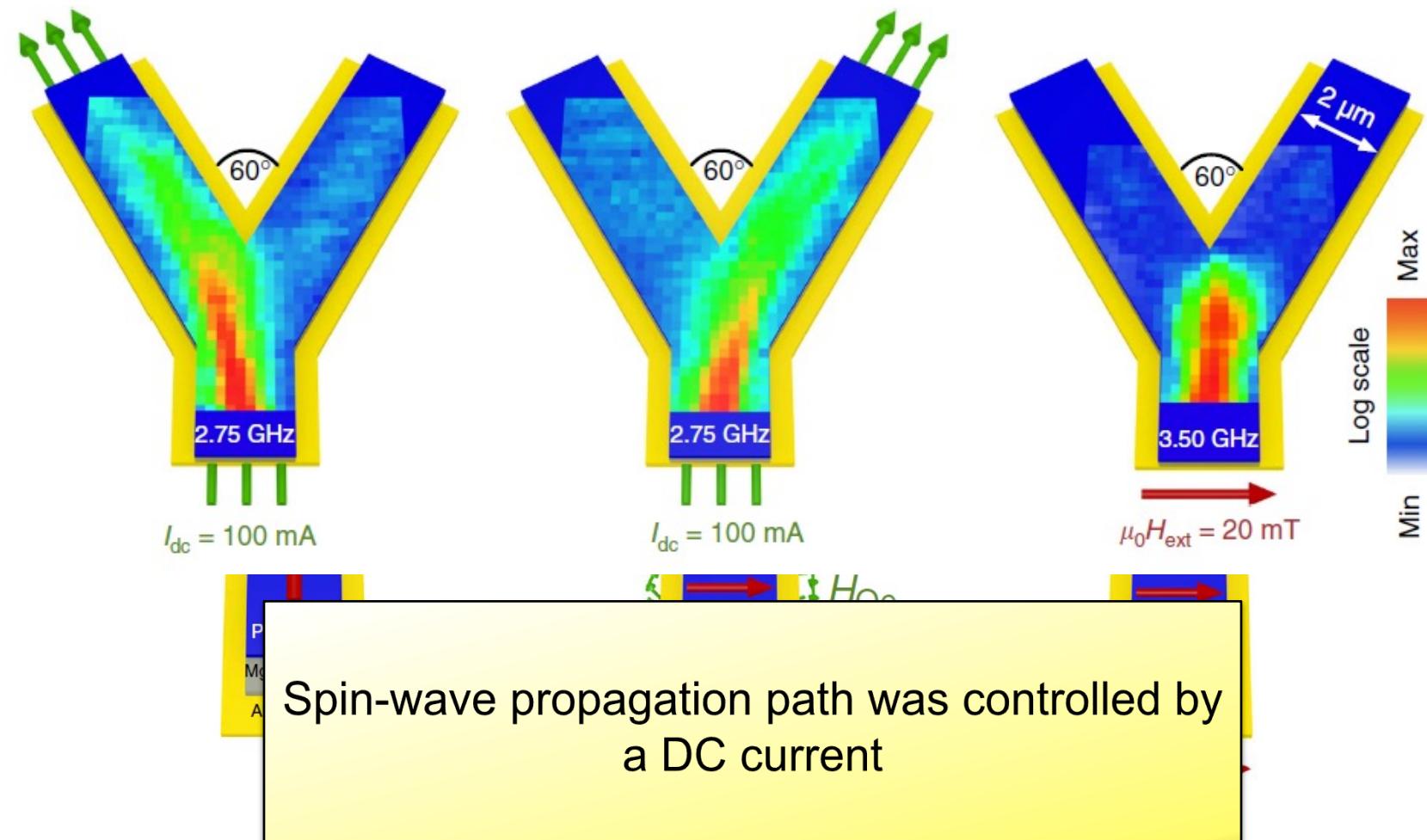
Results of numerical simulation (MuMax 2)



III. Data processing using magnons

- Spin-wave logic gates
- Magnon transistor
- Magnon majority gate
- ➡ Magnon multiplexer

Magnon multiplexer



Vogt, et al., Nat. Commun. 5:3727 (2014)

Burkard Hillebrands

IEEE Magnetics Society Summer School - Minneapolis

June 14-19, 2015

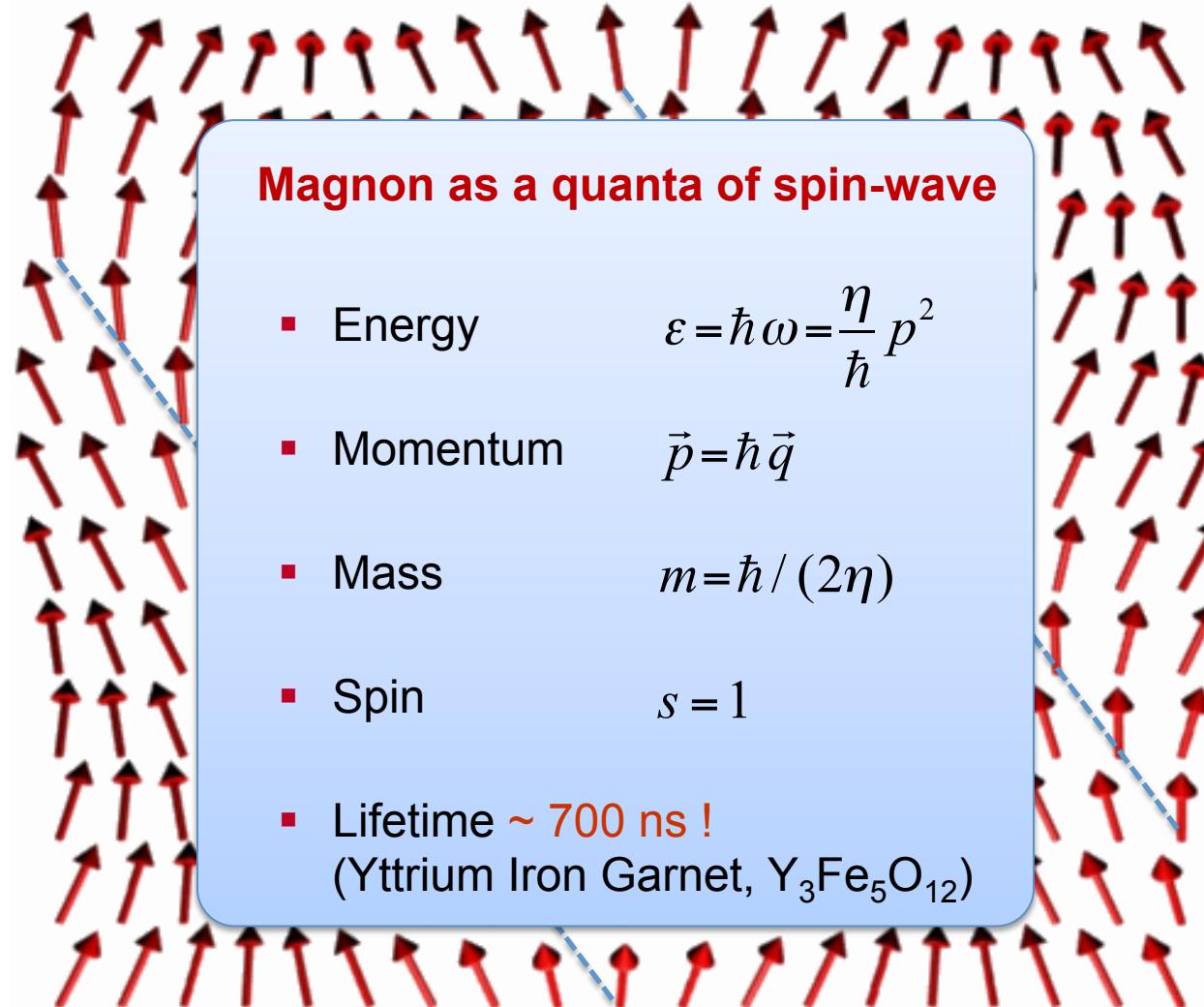
- Classical Computing
 - Scalar variable
 - Boolean logic
- Wave Computing
 - Vector variable
 - Special task data processing
- Quantum Computing
 - Vector state variable
 - Entanglement



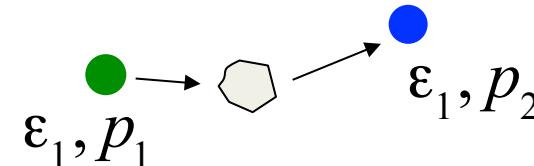
IV. Magnonic supercurrents

Main idea: find macroscopic magnonic quantum states for information transfer and processing

- analogous to superconductivity (Josephson currents), and to superfluidity in ^3He and ^4He
 - free of dissipation (apart from magnon-phonon and magnon-electron coupling)
-
- Bose-Einstein Condensation (BEC) of magnons
 - Supercurrents in magnon condensates

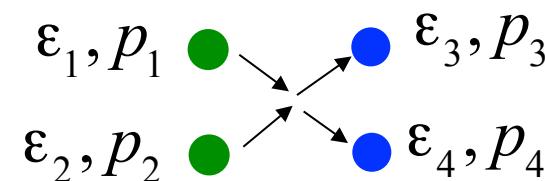


TWO-MAGNON SCATTERING



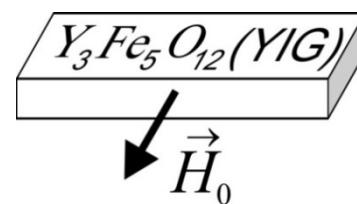
Magnon gas of interacting quasiparticles
Number of particles is conserved

FOUR-MAGNON SCATTERING



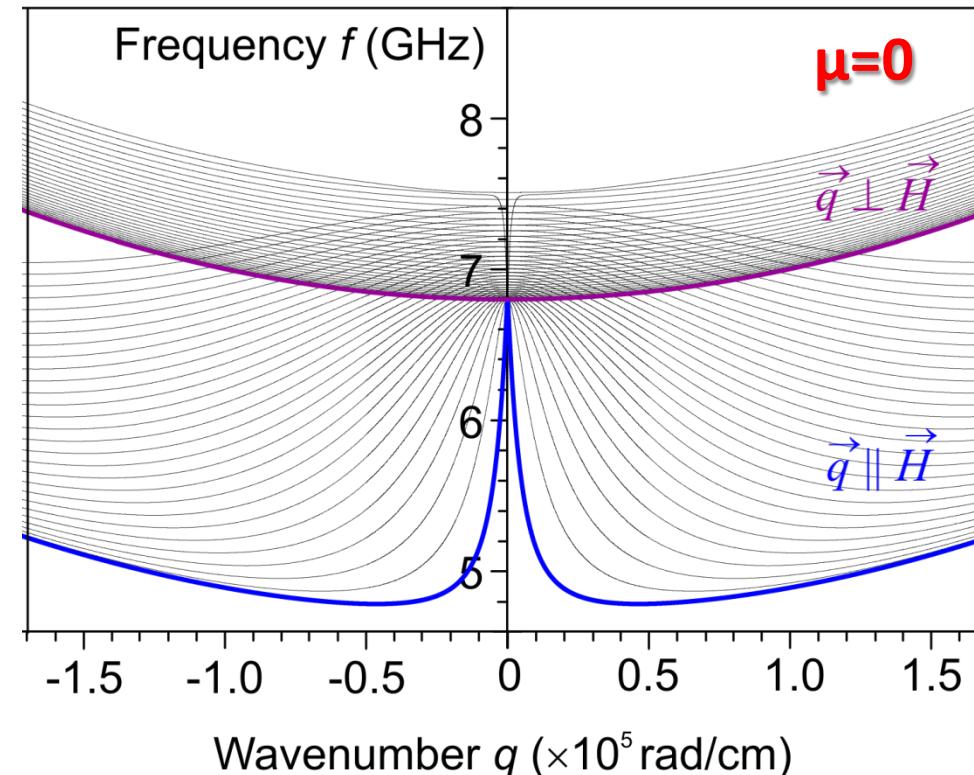
Magnon distribution

Magnons are **bosons** ($s=1$)
 and thus as any quasi-particles
 are described by Bose-Einstein distribution
 with **zero chemical potential**



**Bose-Einstein
distribution**

$$\rho(f) = \frac{D(f)}{\exp\left(\frac{hf - \mu}{k_B T}\right) - 1}$$

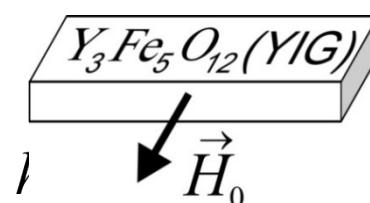


Control of magnon gas density by parametric pumping

Energy and momentum conservation laws

$$\begin{cases} \vec{q}_{\text{sw}} + \vec{q}'_{\text{sw}} = \vec{q}_p \simeq 0 \\ f_{\text{sw}} + f'_{\text{sw}} = f_p \end{cases}$$

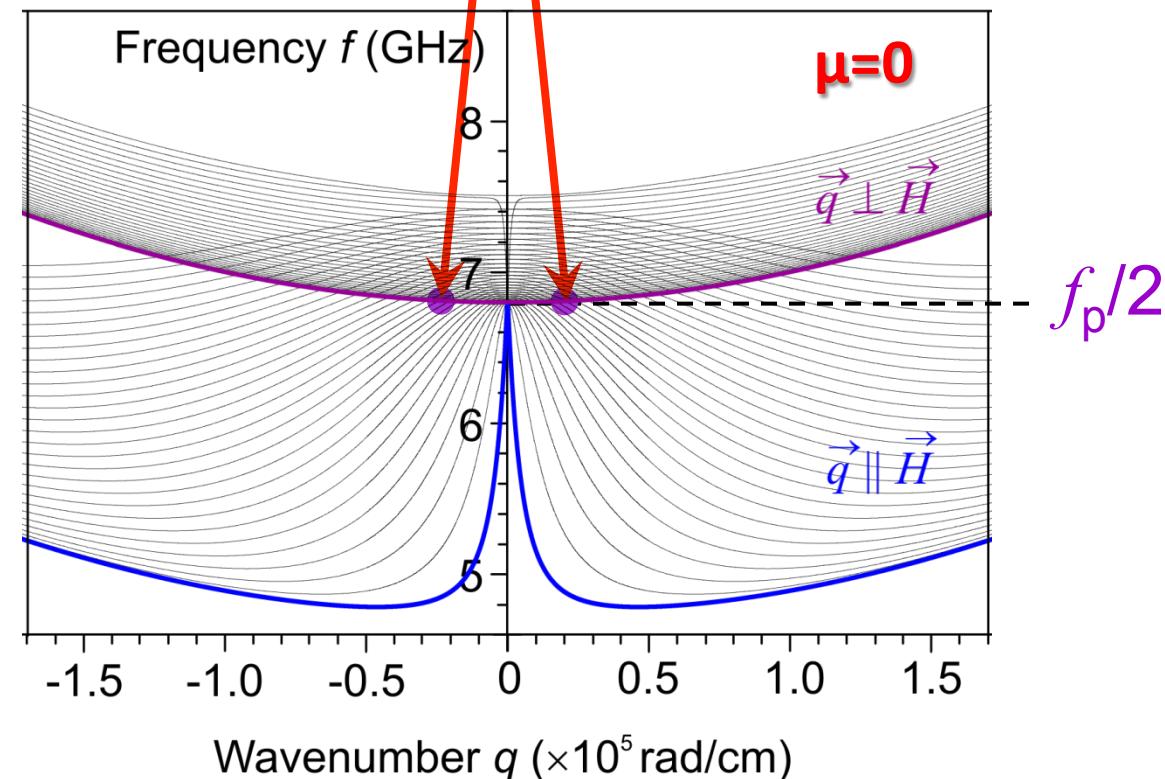
f_p Parametric pumping by electromagnetic wave at microwave frequency



Bose-Einstein distribution

$$\rho(f) = \frac{D(f)}{\exp\left(\frac{hf - \mu}{k_B T}\right) - 1}$$

μ : chemical potential

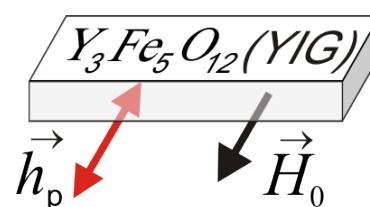


Control of magnon gas density by parametric pumping

Energy and momentum conservation laws

$$\begin{cases} \vec{q}_{\text{sw}} + \vec{q}'_{\text{sw}} = \vec{q}_p \approx 0 \\ f_{\text{sw}} + f'_{\text{sw}} = f_p \end{cases}$$

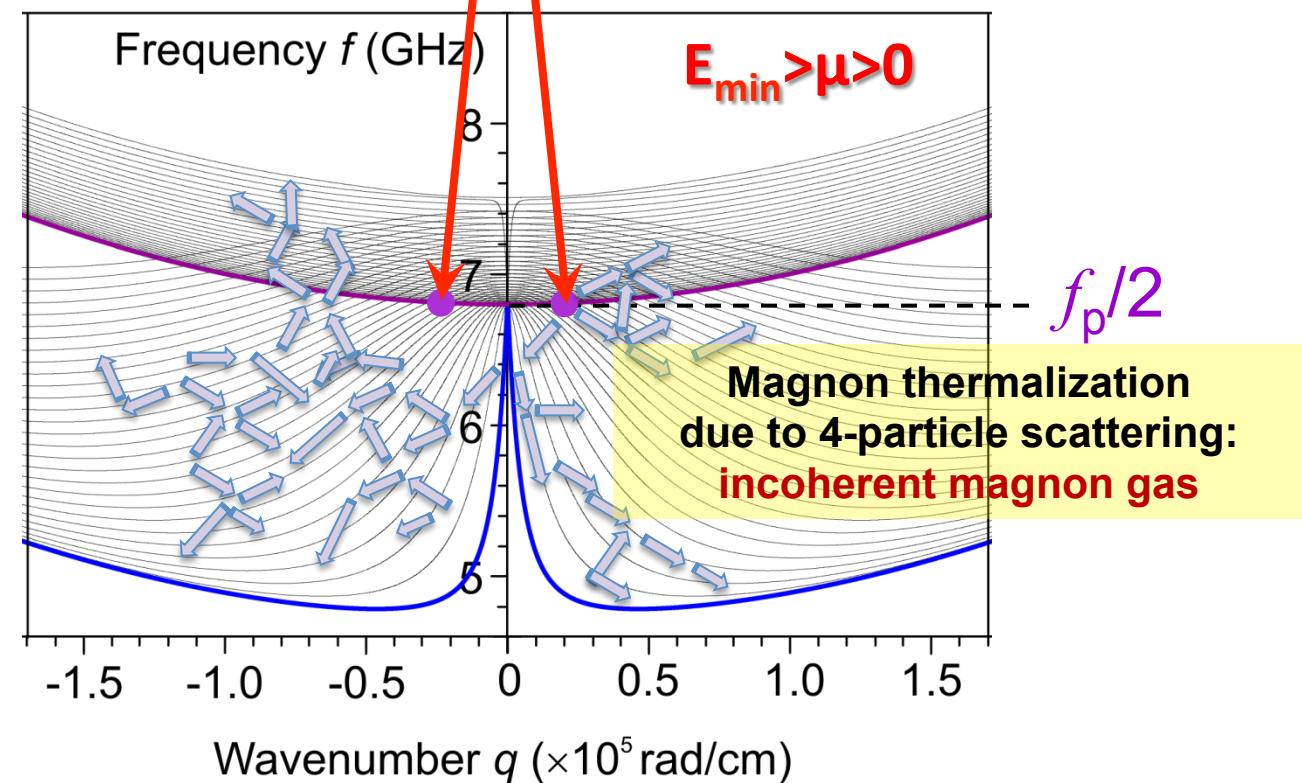
f_p Parametric pumping by electromagnetic wave at microwave frequency



Bose-Einstein distribution

$$\rho(f) = \frac{D(f)}{\exp\left(\frac{hf - \mu}{k_B T}\right) - 1}$$

μ : chemical potential

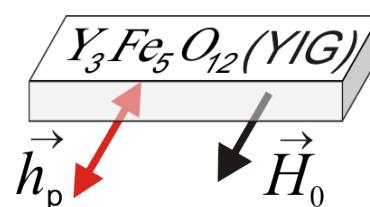


Control of magnon gas density by parametric pumping

Energy and momentum conservation laws

$$\begin{cases} \vec{q}_{\text{sw}} + \vec{q}'_{\text{sw}} = \vec{q}_p \approx 0 \\ f_{\text{sw}} + f'_{\text{sw}} = f_p \end{cases}$$

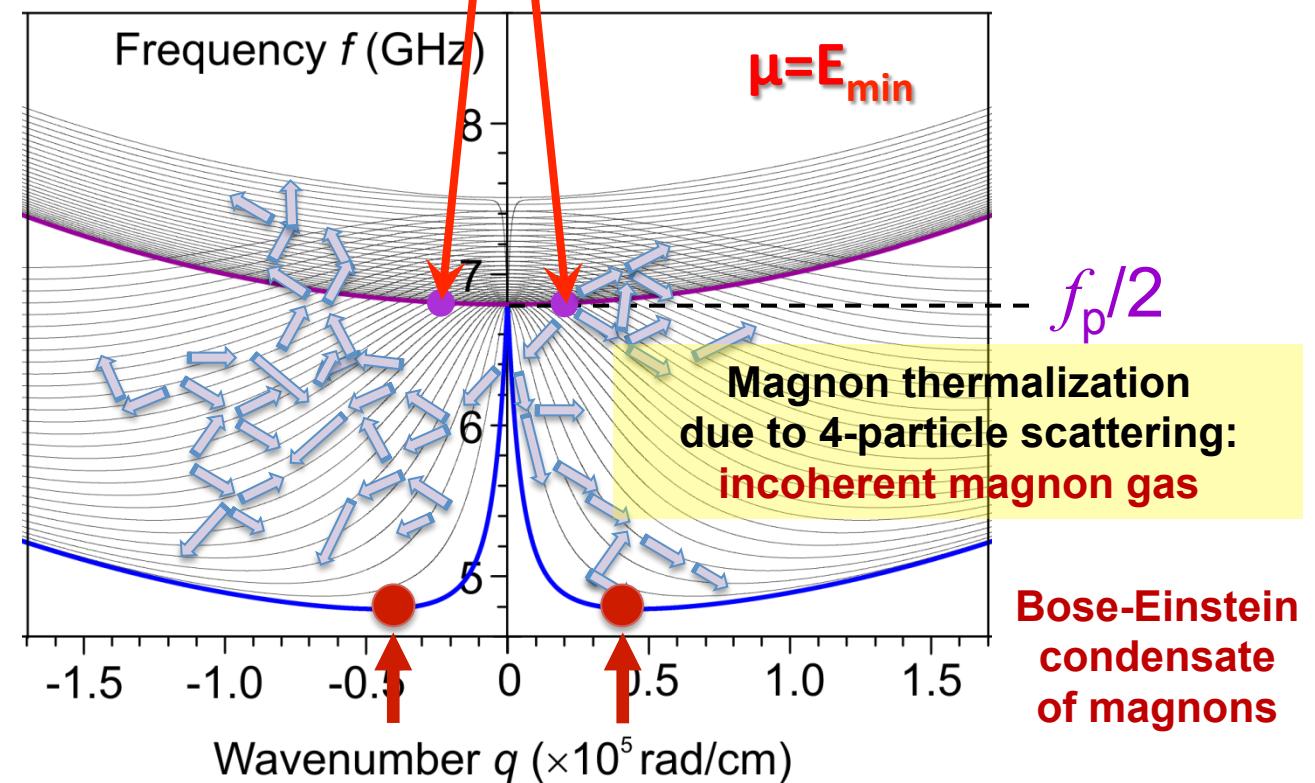
Parametric pumping by electromagnetic wave at microwave frequency



Bose-Einstein distribution

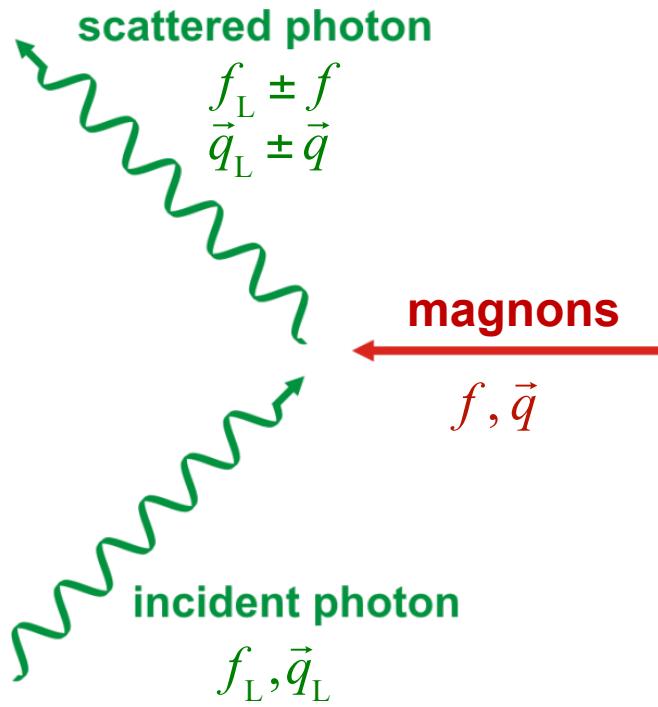
$$\rho(f) = \frac{D(f)}{\exp\left(\frac{hf - \mu}{k_B T}\right) - 1}$$

μ : chemical potential



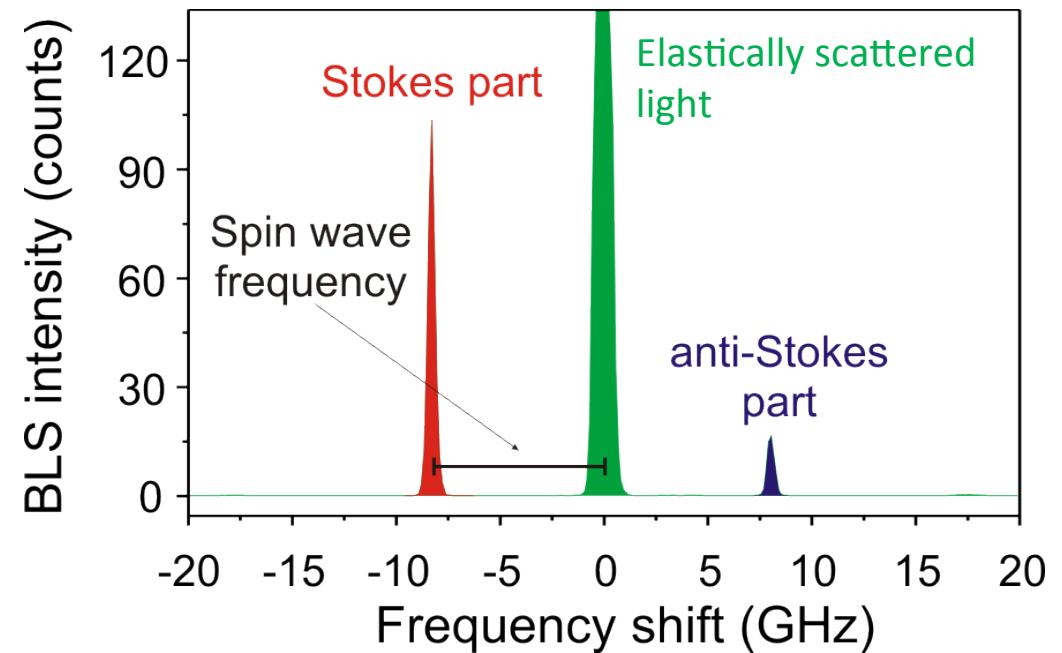
Brillouin light scattering process

= inelastic scattering of photons from spin waves

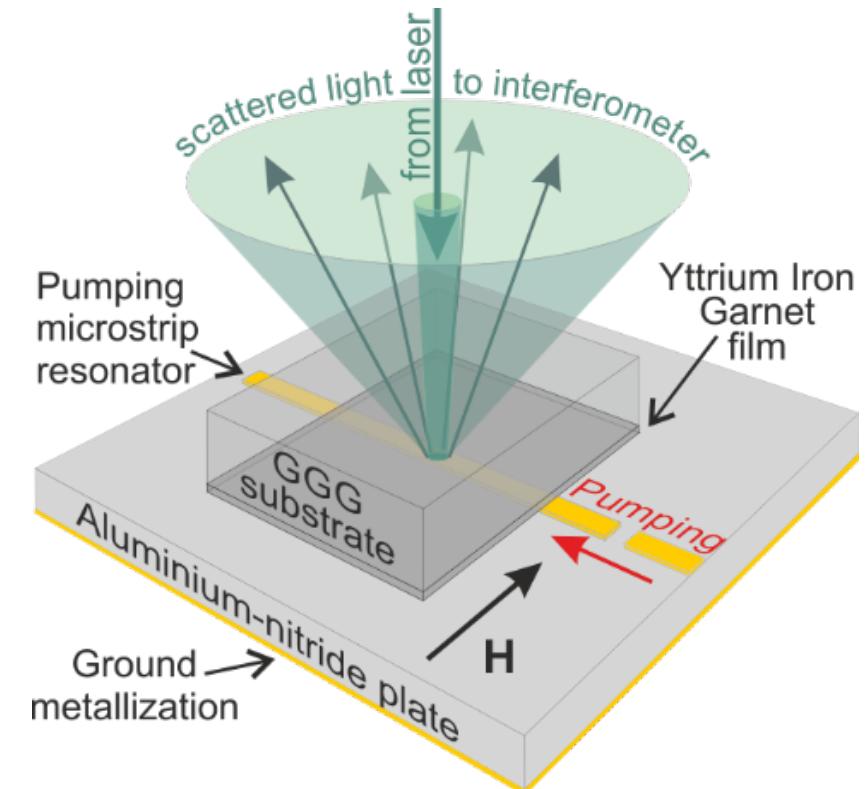
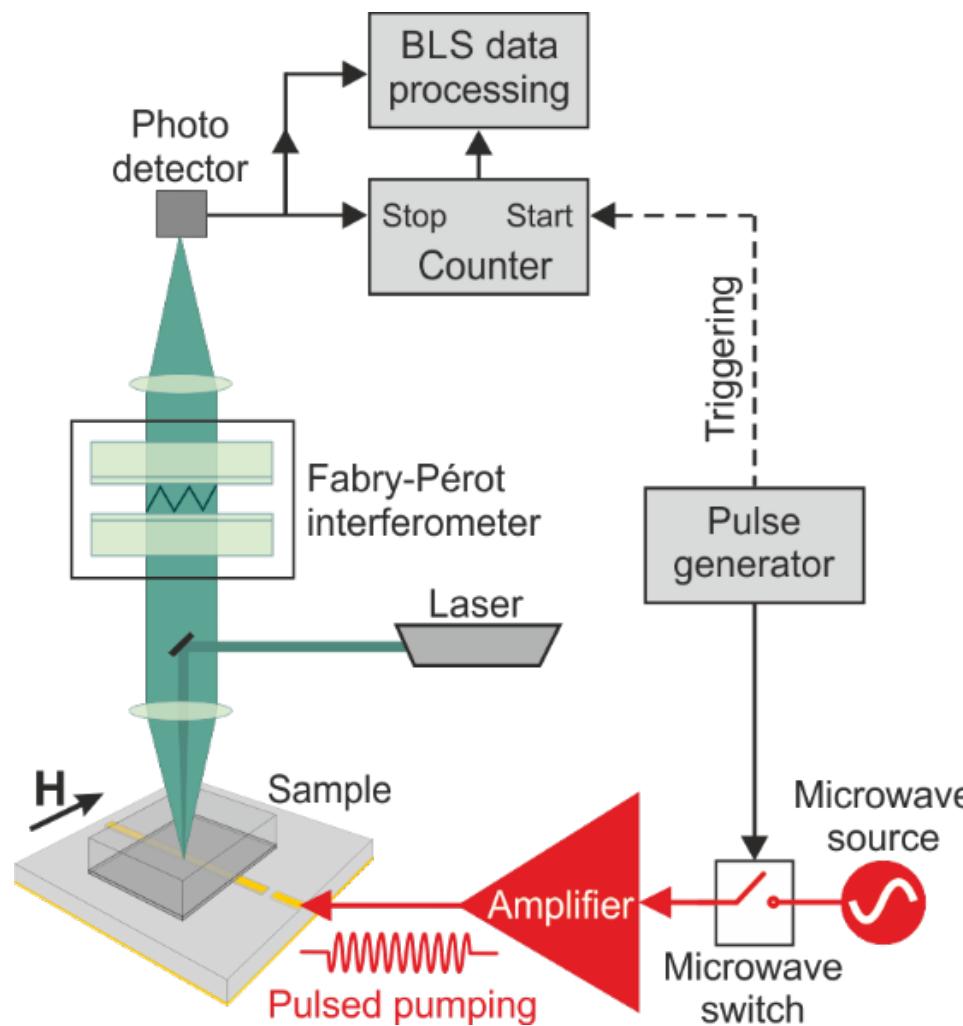


$$f_{\text{scattered L}} = f_L \pm f$$

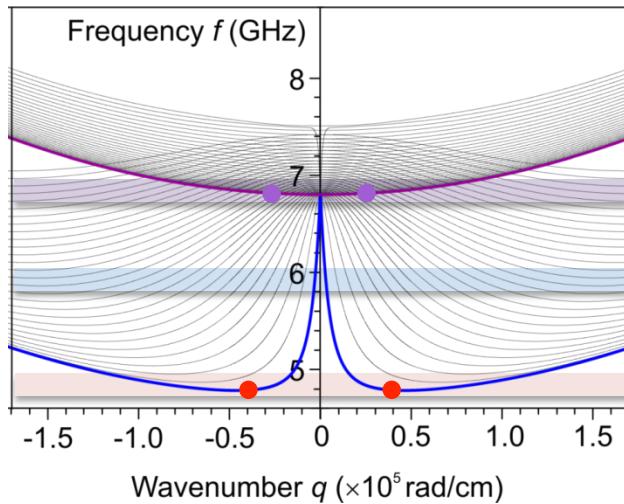
$$\vec{q}_{\text{scattered L}} = \vec{q}_L \pm \vec{q}$$



Time-resolved Brillouin light scattering spectroscopy



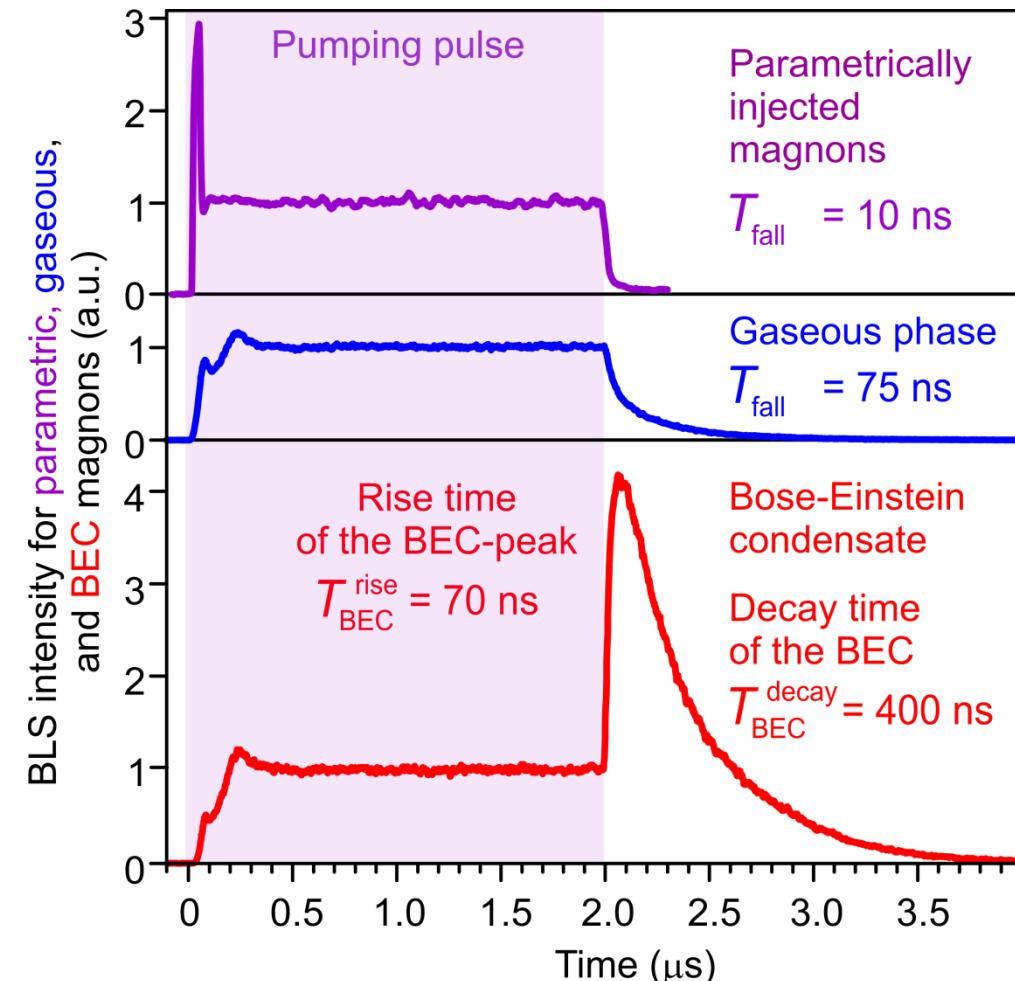
Parametric magnons, gaseous phase, and magnon BEC



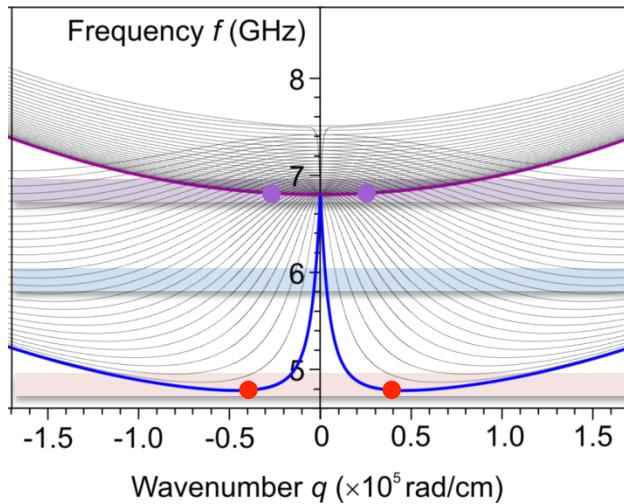
Decrease of density of parametric magnons and gaseous magnon phase



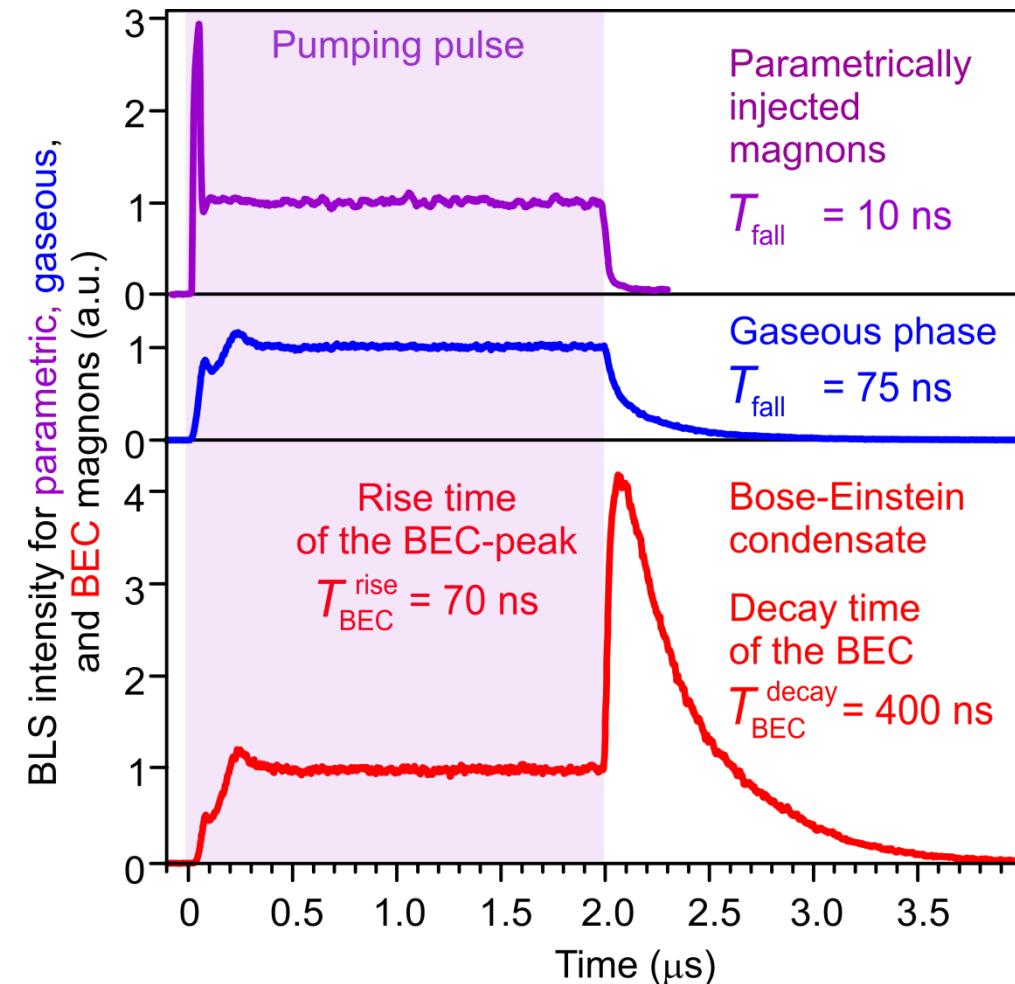
Sharp increase of intensity of pumping free BEC of magnons



Parametric magnons, gaseous phase, and magnon BEC



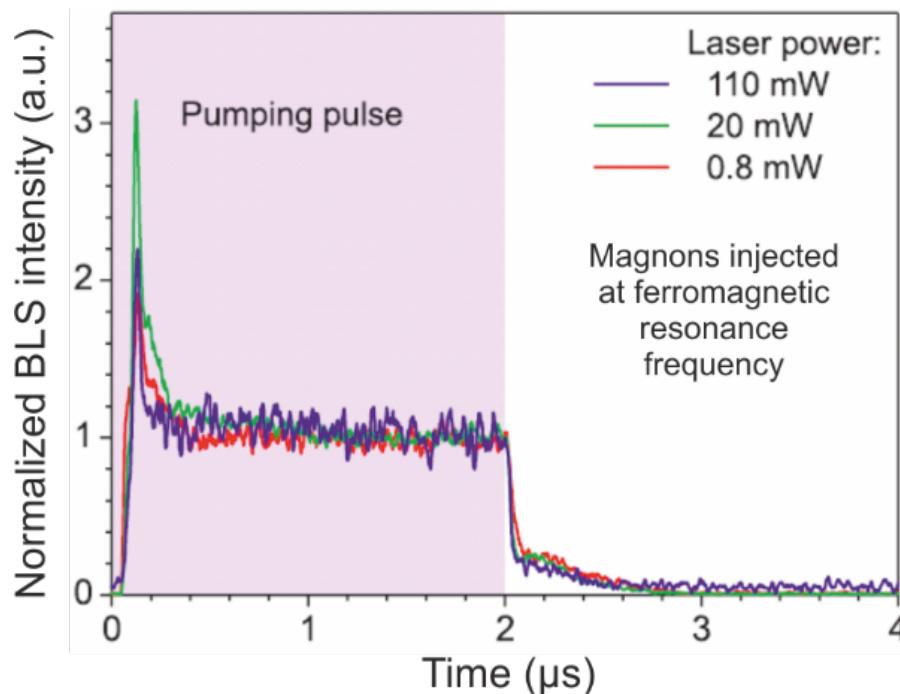
Evaporative supercooling of strongly overheated low energy area of the magnon gas



Serga et al., Nature Communications **5**, 4452 (2014)

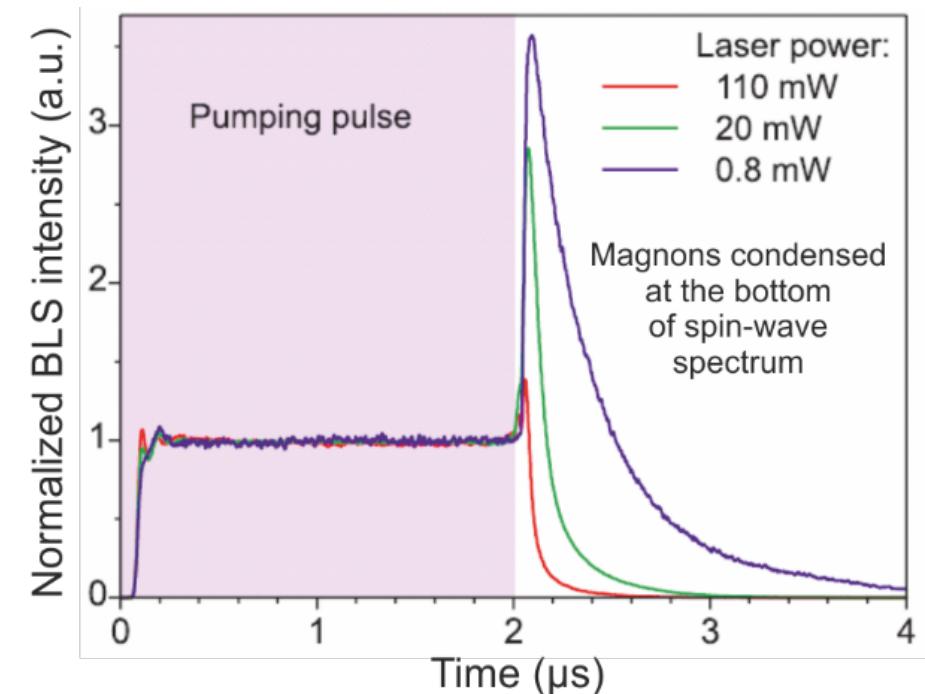
Parametrically injected magnons

- No influence from laser power on temporal dynamics !
- Laser heating decreases magnetization and thus strongly shifts down the magnon dispersion branch

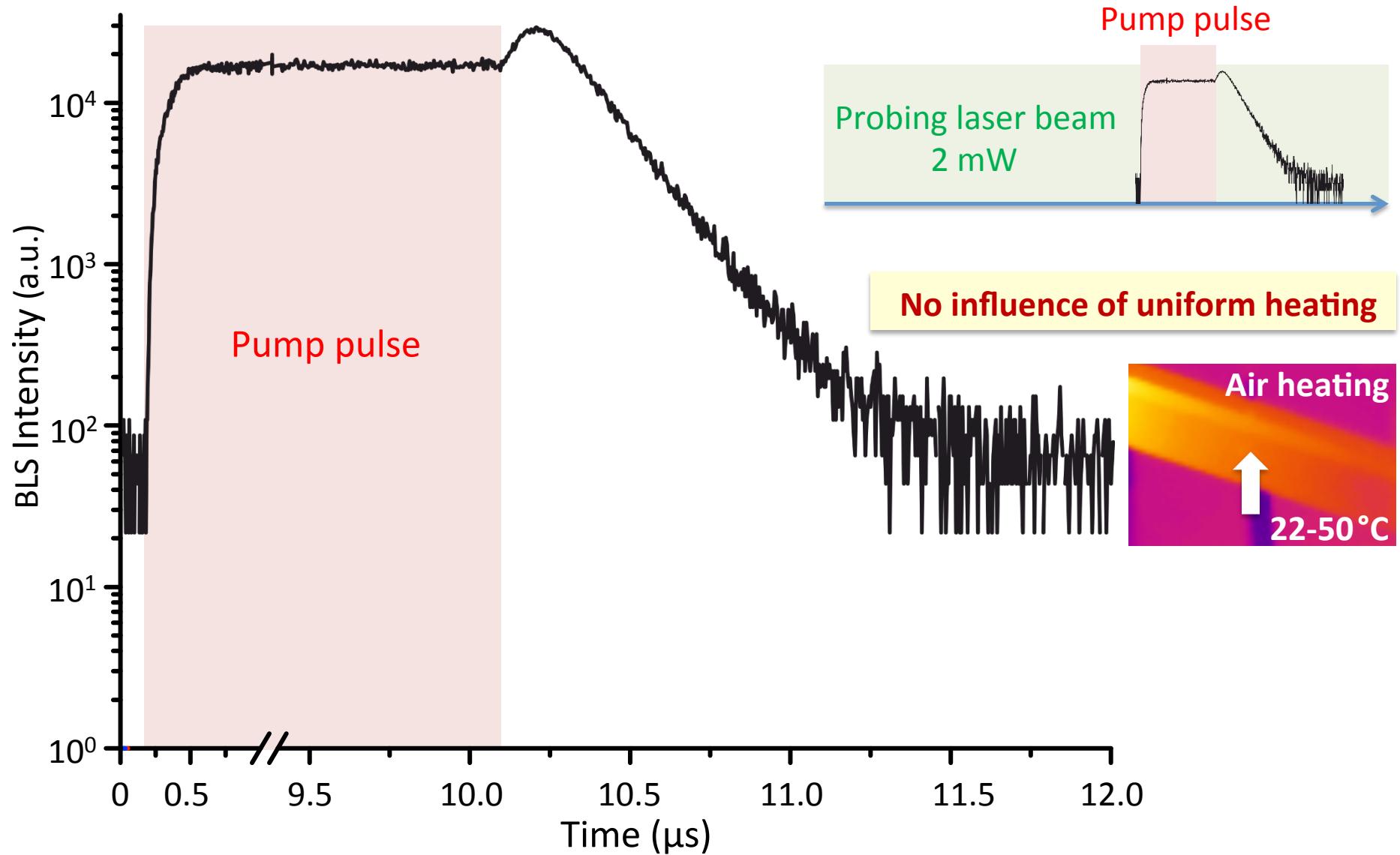


Bose-Einstein condensate

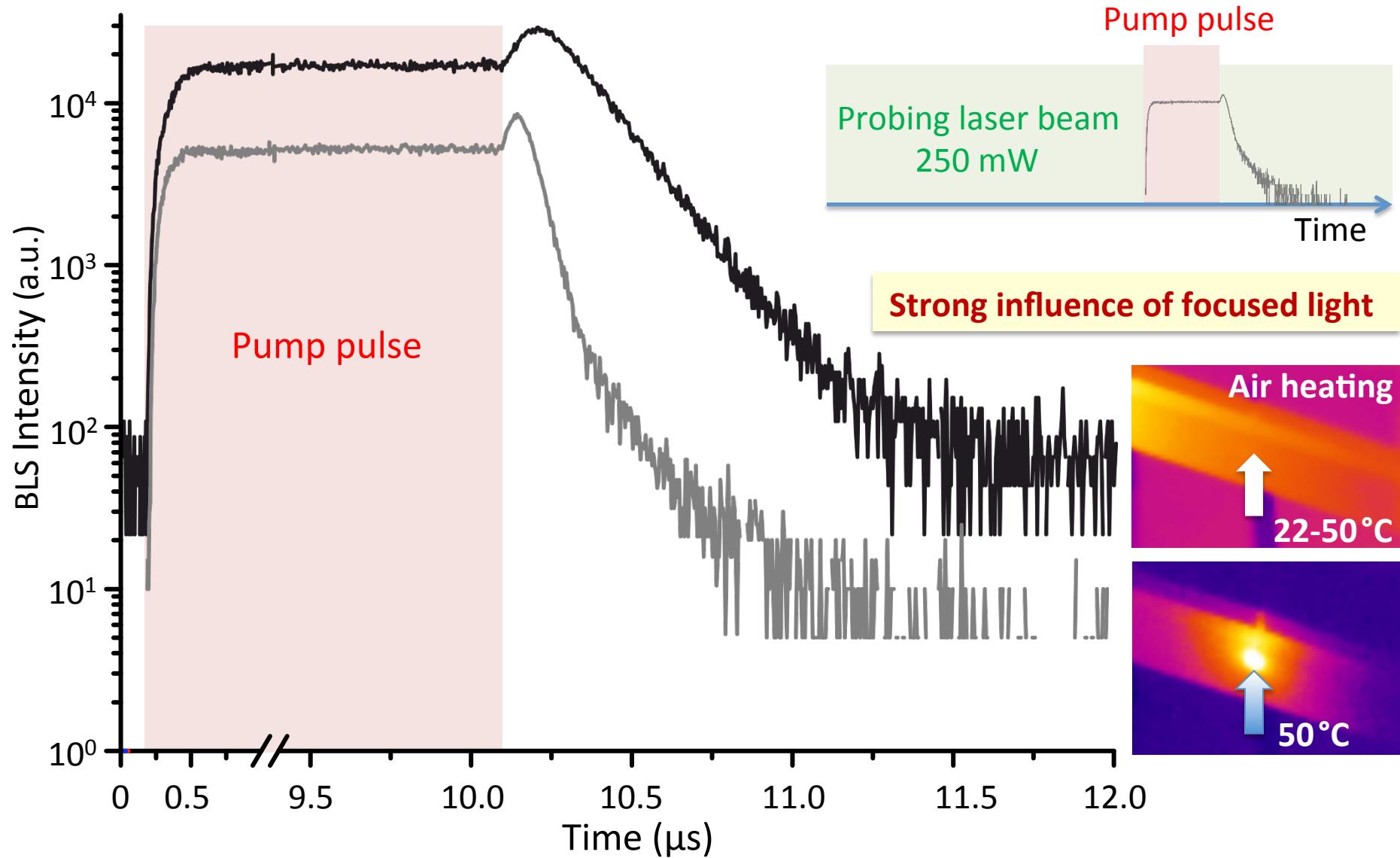
- Increasing laser power results in significant decrease of the BEC's lifetime and amplitude !
- Only weak frequency shift of BEC due to temperature change



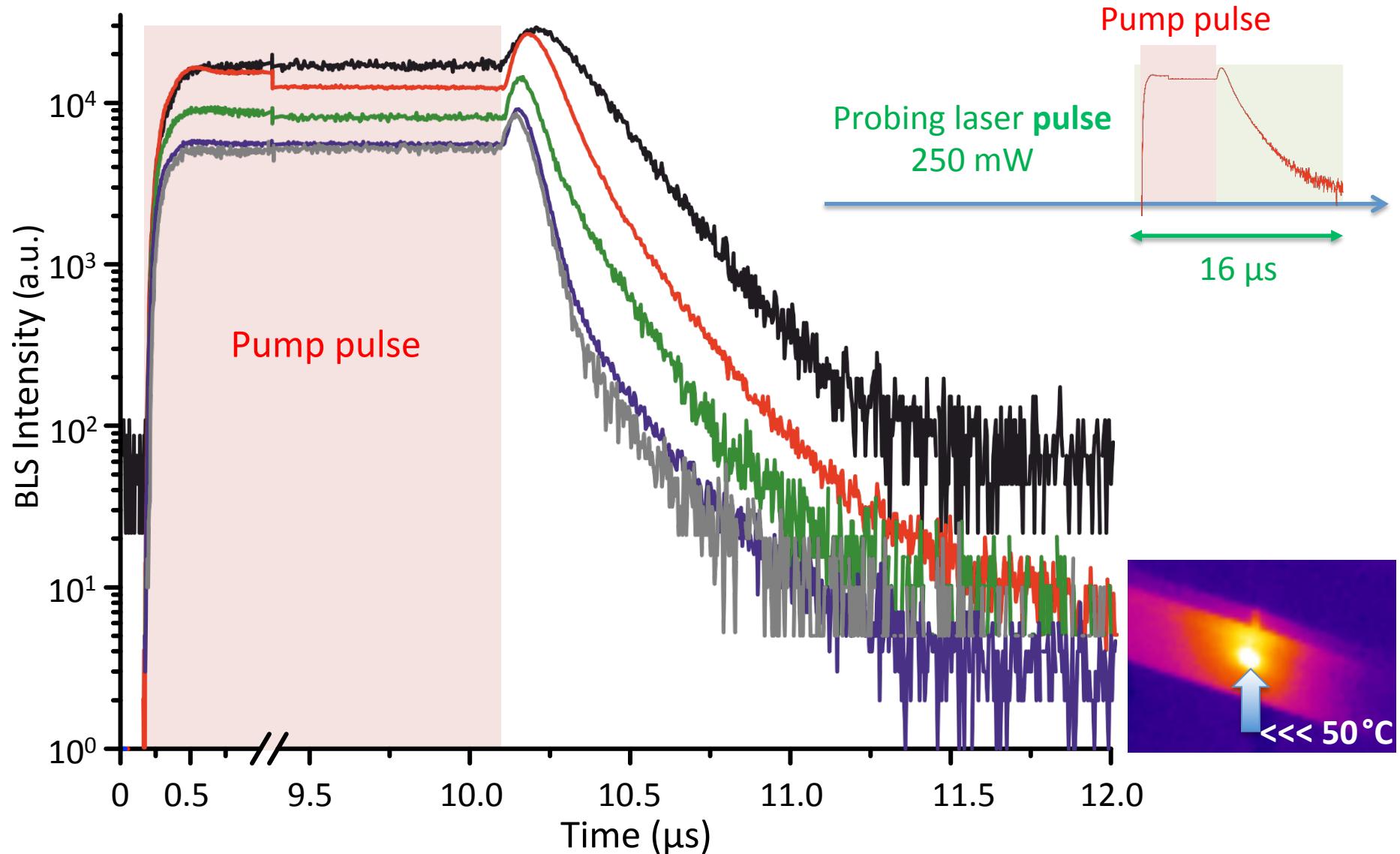
Magnon supercurrents – magnon BEC dynamics in thermal gradient



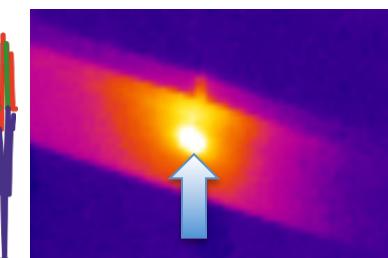
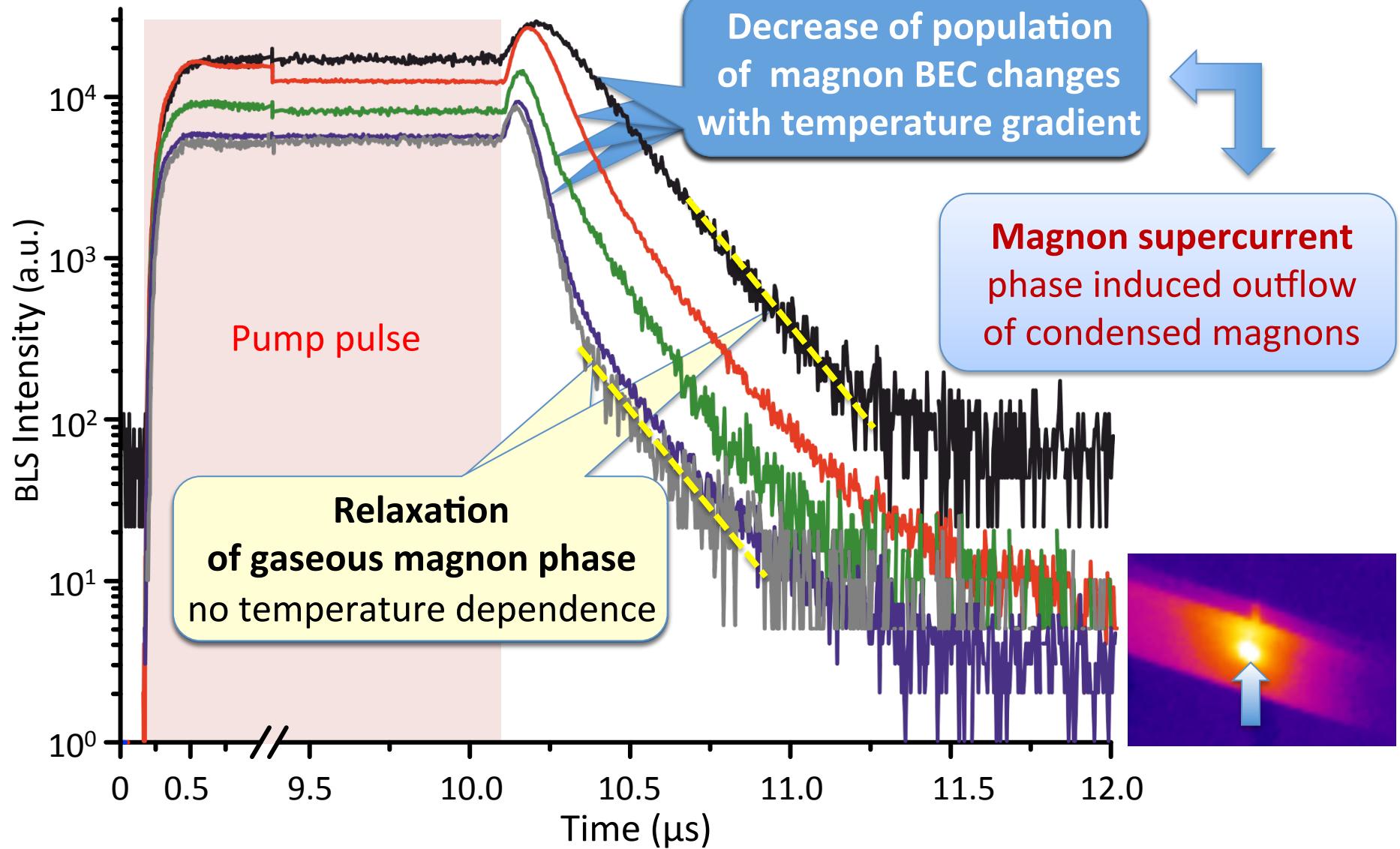
Magnon supercurrents – magnon BEC dynamics in thermal gradient



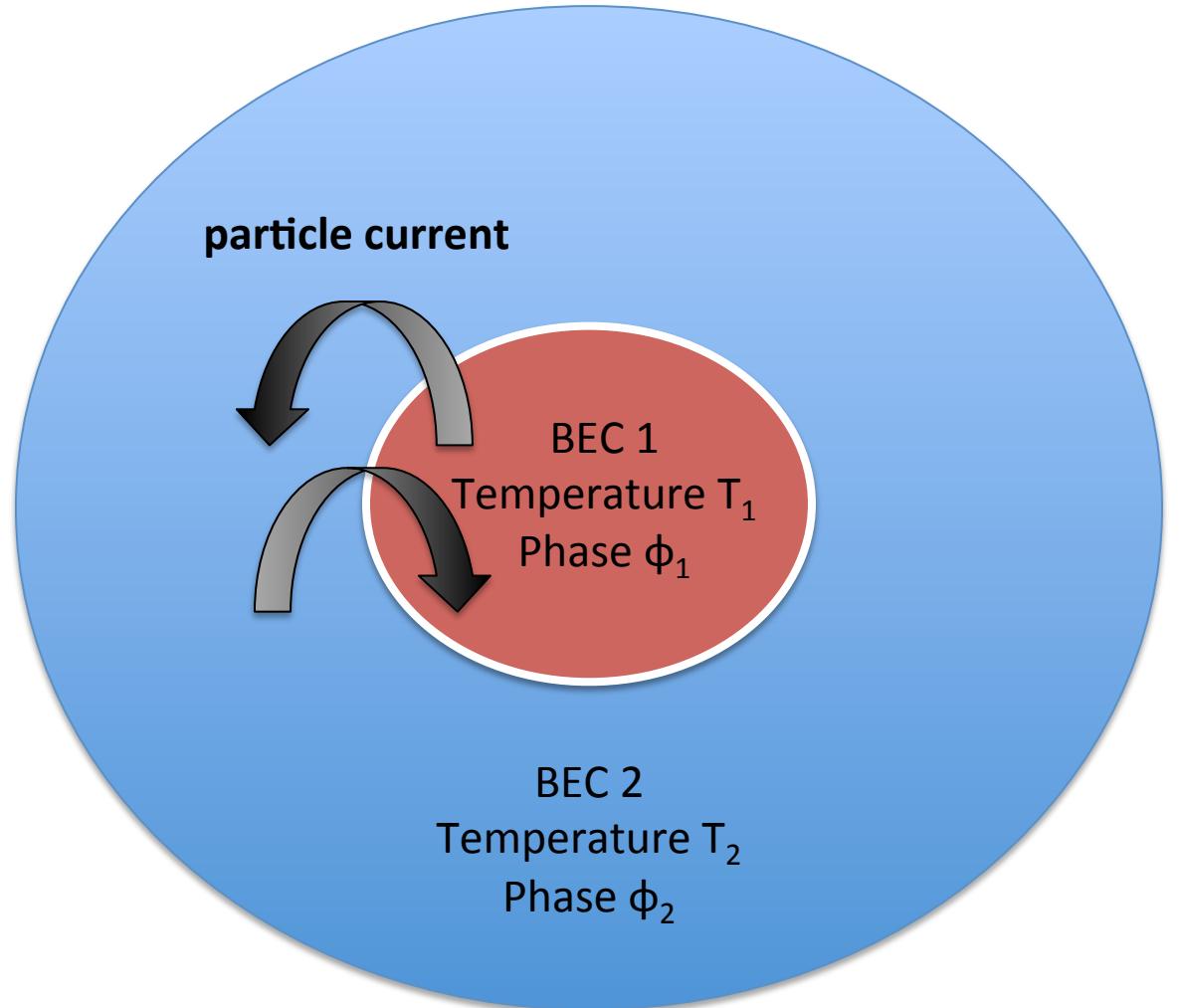
Magnon supercurrents – magnon BEC dynamics in thermal gradient



Magnon supercurrents – magnon BEC dynamics in thermal gradient

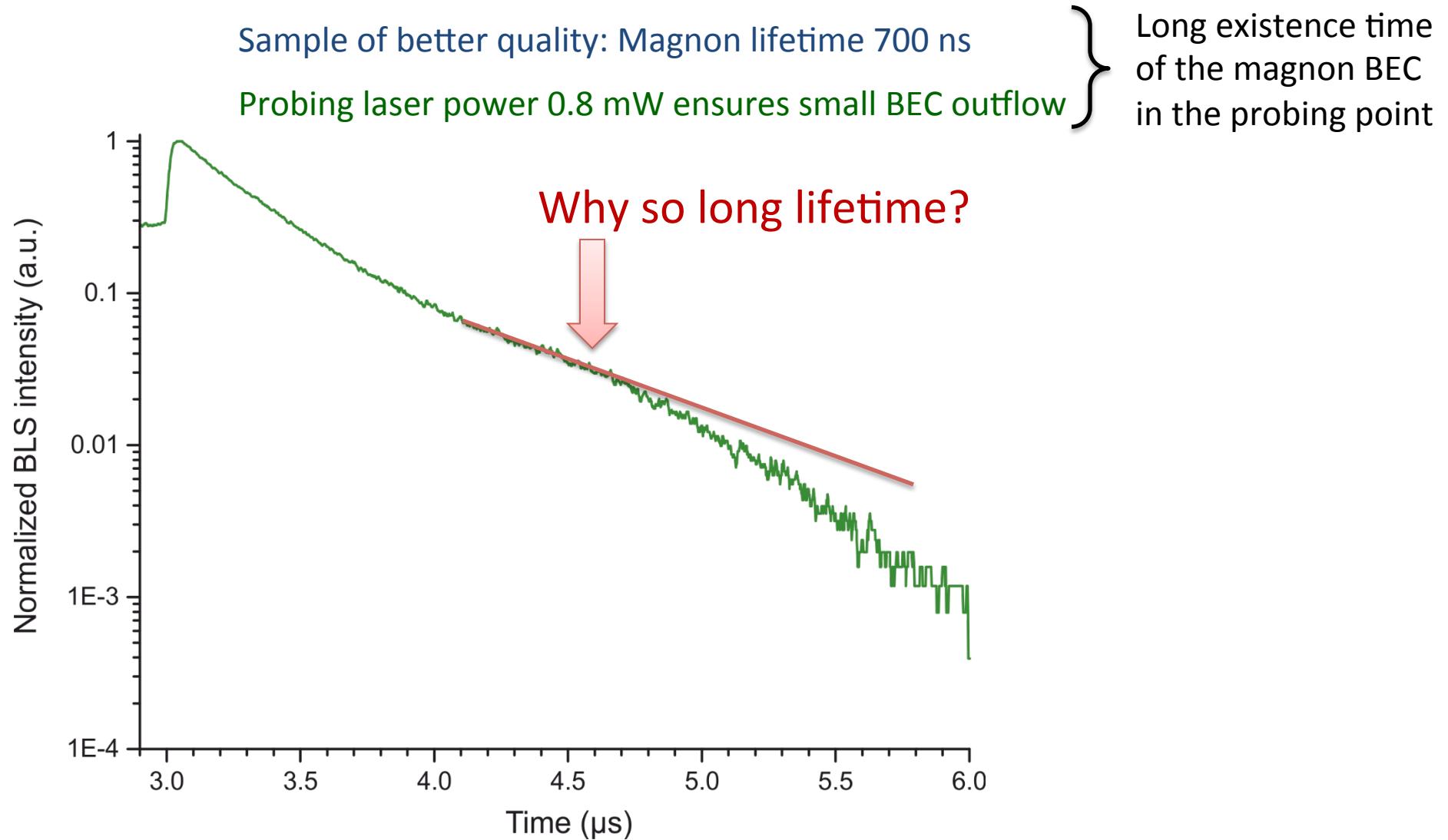


Magnonic supercurrents

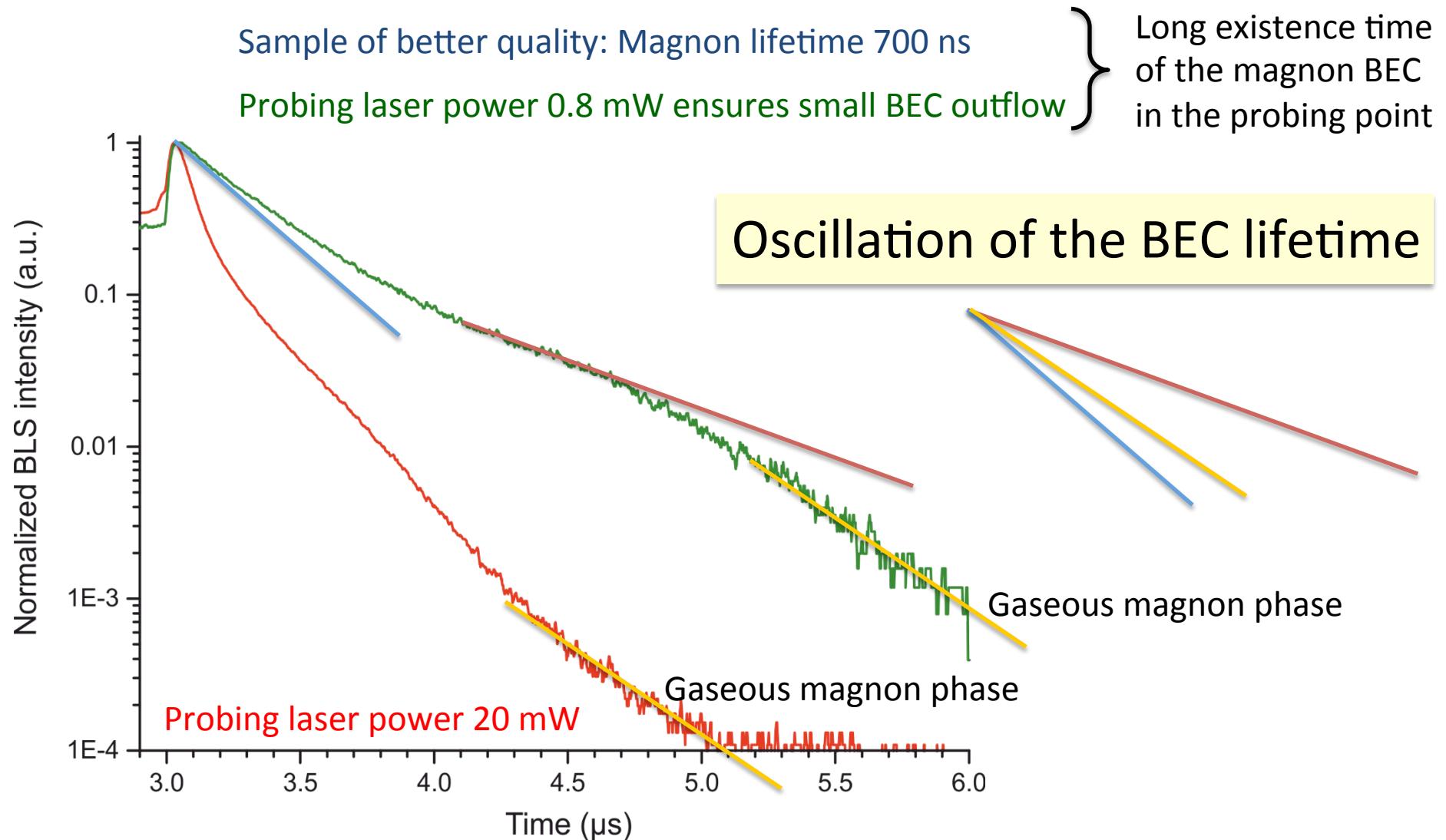


$$\begin{aligned} T_1 &> T_2 \\ \rightarrow \Phi_1 &\neq \Phi_2 \end{aligned}$$

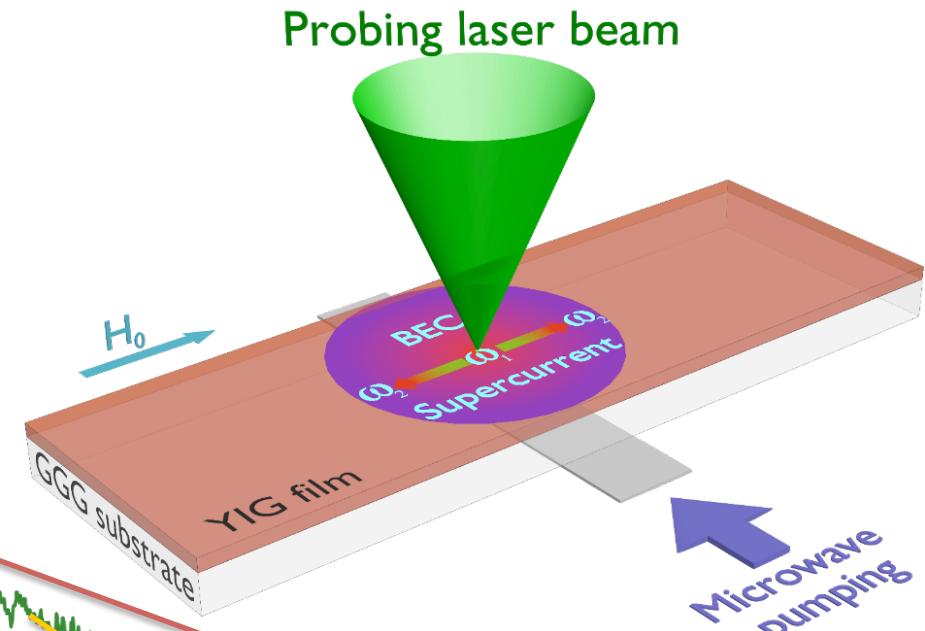
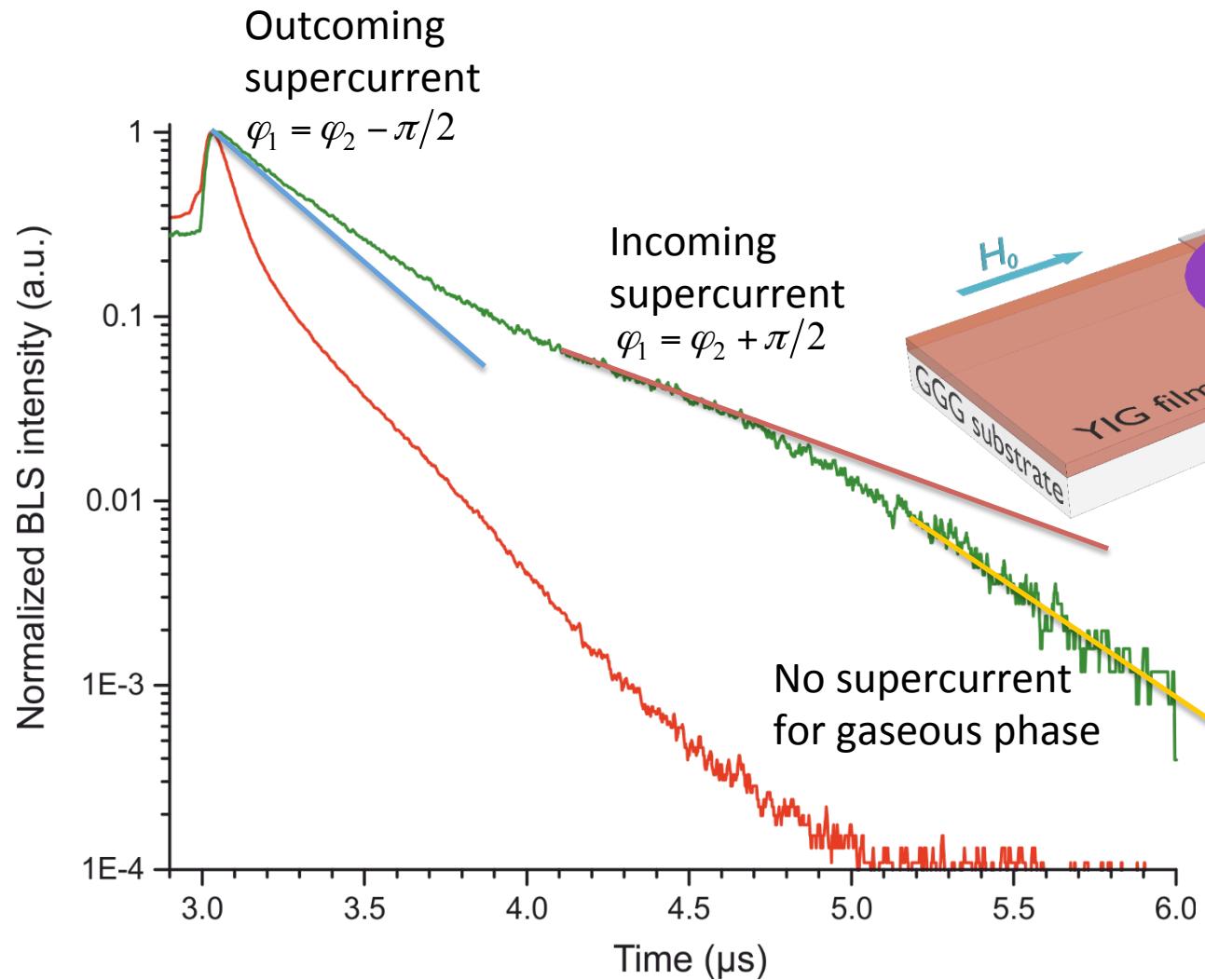
Long-lifetime BEC in thermal gradient



Oscillations of the BEC lifetime



BEC in gradient of saturation magnetization



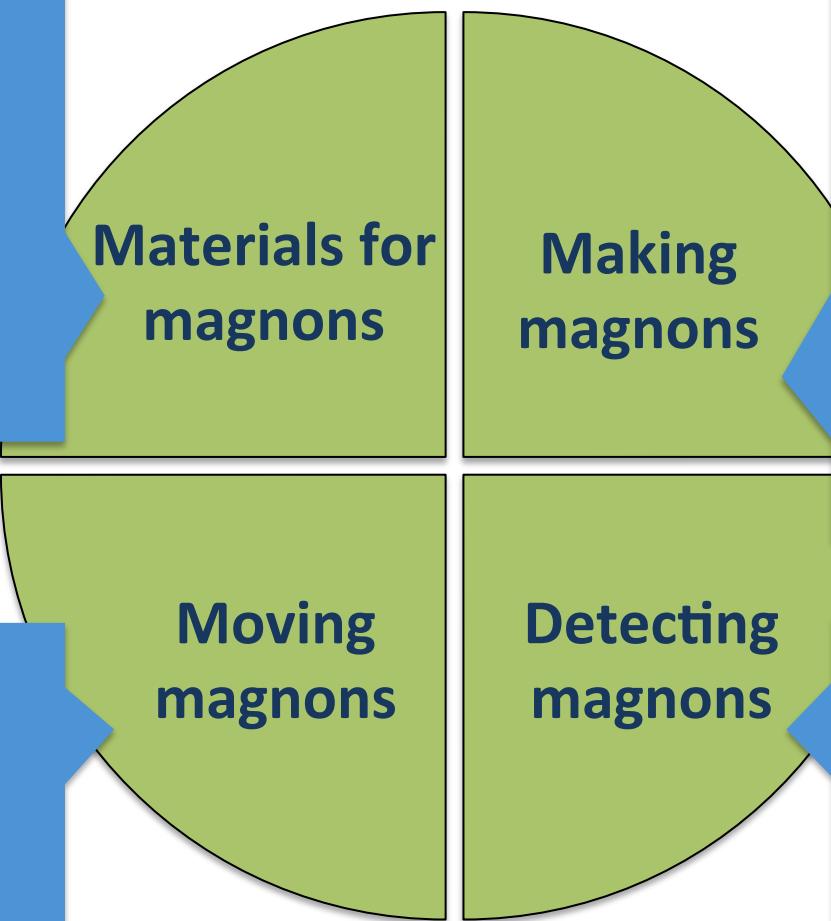
$$J \propto \nabla \varphi$$

$$\Delta\varphi = \Delta\omega t$$

$$J \propto \sin(\Delta\omega t)$$

Magnonics: Challenges

- New insulators
- Materials with controlled damping
- Novel artificial materials
- Materials with controlled spin-orbit coupling



Materials for magnons

Making magnons

Moving magnons

Detecting magnons

- Magnon Bose (Einstein) condensates
- Spin Transfer Torque effect in insulators
- Miniaturization
- THz magnonics
- Single magnon gener. and detection

- Magnon supercurrents
- Dzyaloshinskii-Moria interactions
- Electric field effects

- Enhancement of existing techniques
- TMR
- Quantum effects